



# Mode Mismatching Analysis

#beaminthehole

LIGO SURF 2022

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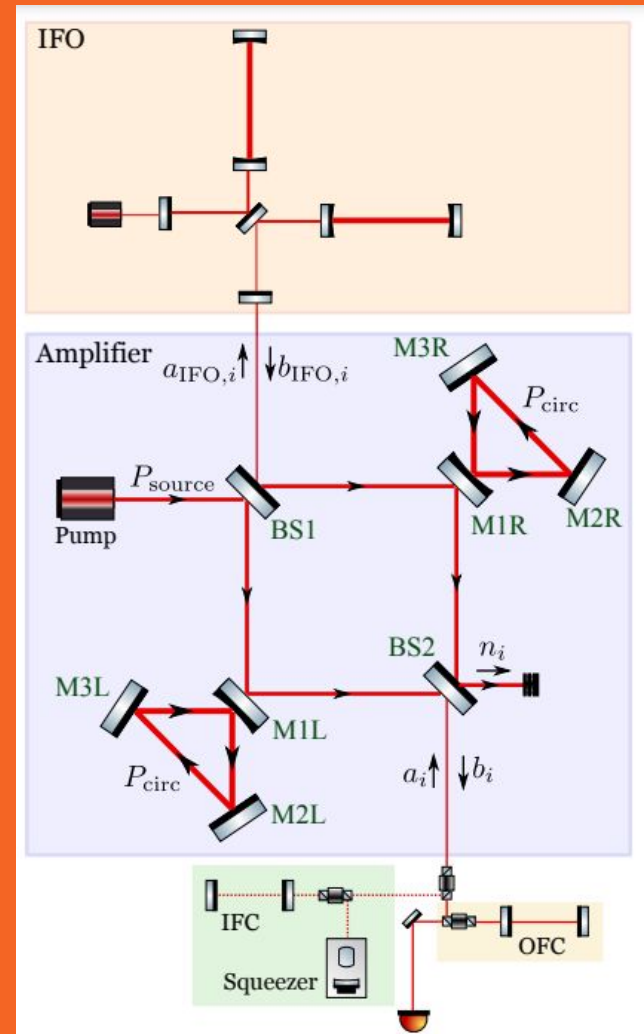
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## Project Overview:

- Original project motivations
- Background on mode mismatch
- Thick lens analysis and lens aberrations
- Modeling with Gaussian beam mode coupling
- Future plans

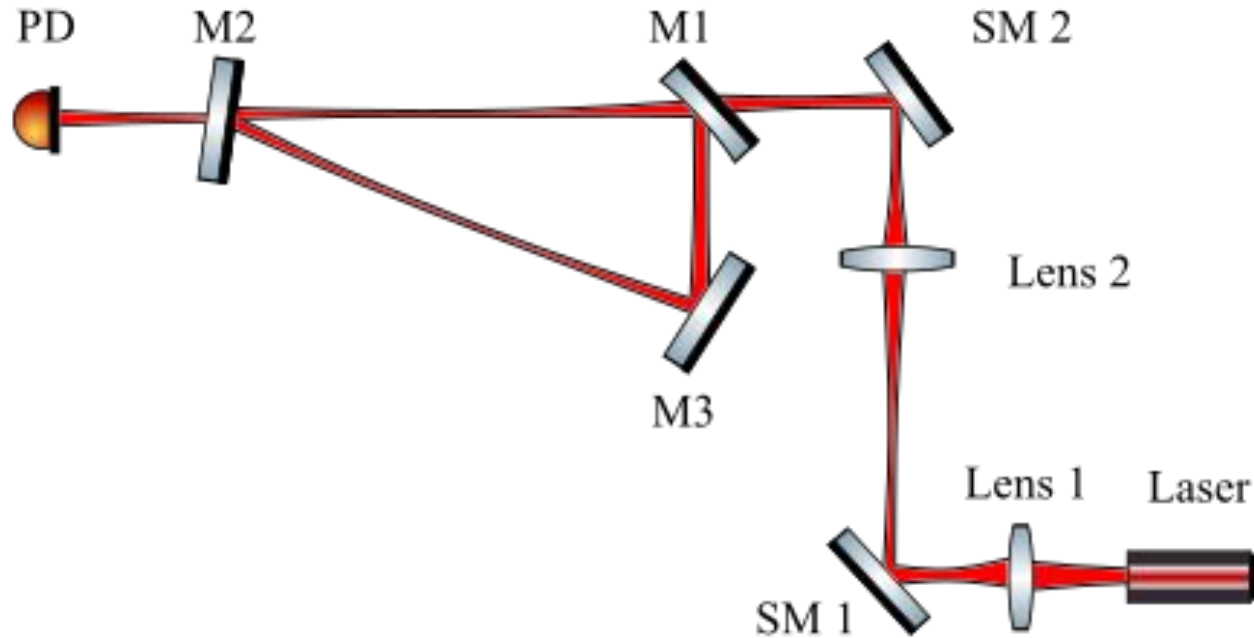
# Original Mission

- Phase Sensitive Optomechanical Amplifier (PSOMA) experiencing mode mismatching.
- Test thick lens analysis.
- Consider lens aberrations creating higher order modes

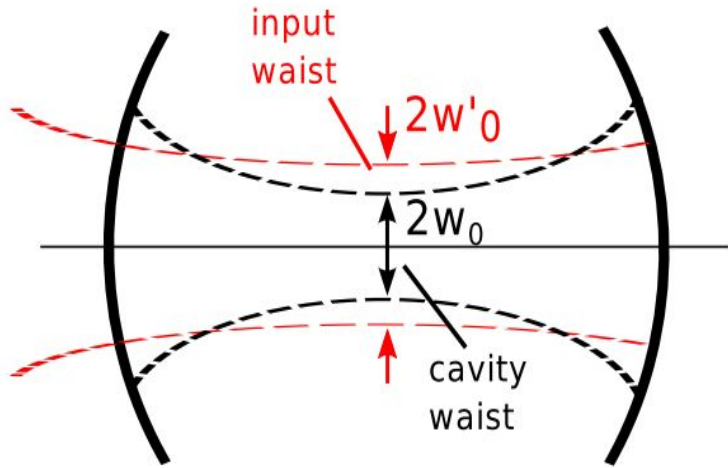


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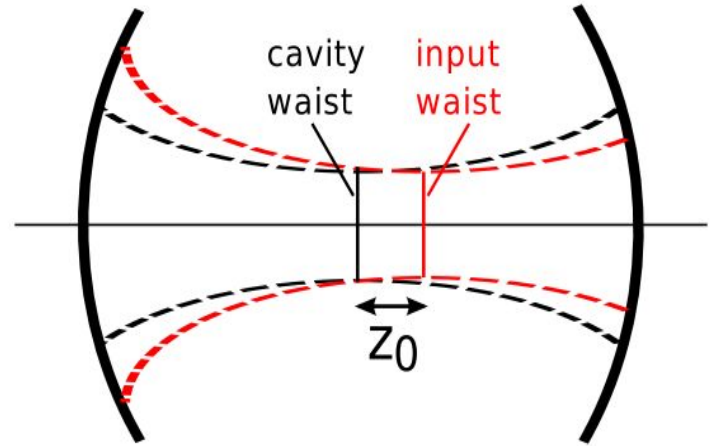
# Tabletop design



# #Beam in the hole



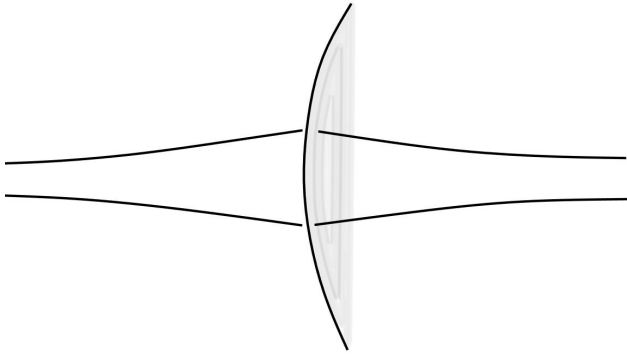
a)



b)

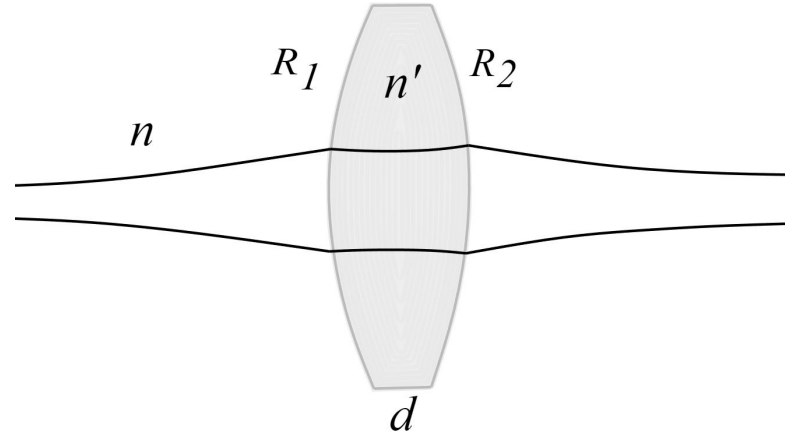
# Through Thin and Thick

Thin lens

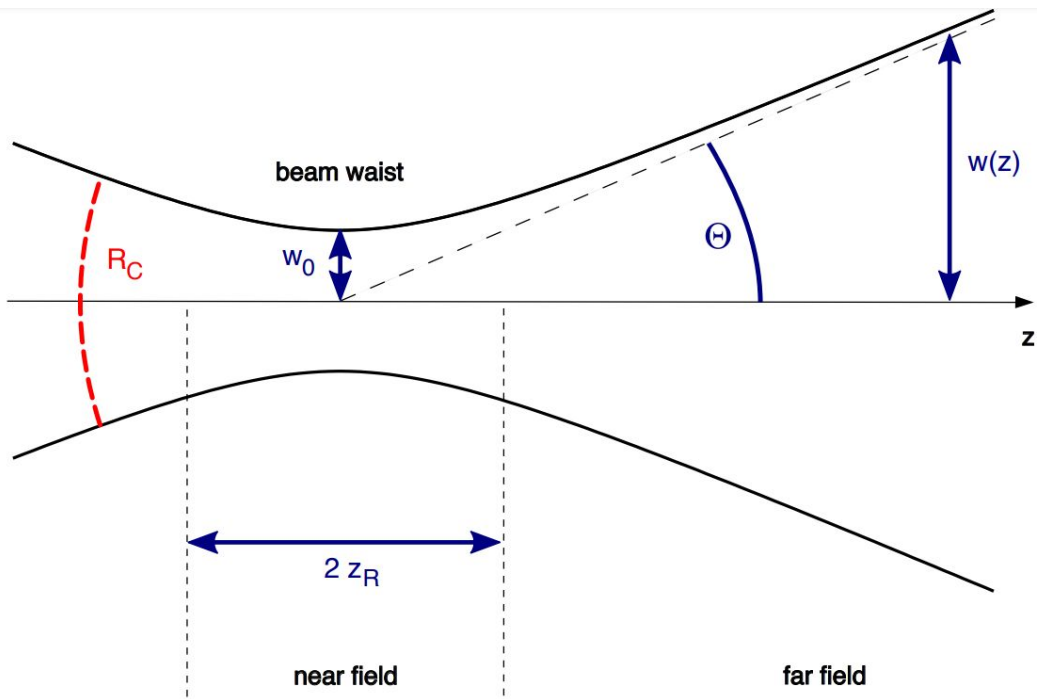


$$\begin{pmatrix} 1 & 0 \\ -\frac{1}{f} & 1 \end{pmatrix}$$

Thick lens



$$\begin{pmatrix} 1 & 0 \\ \frac{n' - n}{R_2} & 1 \end{pmatrix} \begin{pmatrix} 1 & \frac{d}{n'} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ \frac{n - n'}{R_1} & 1 \end{pmatrix}$$



$R_c$  = Beam front radius of curvature

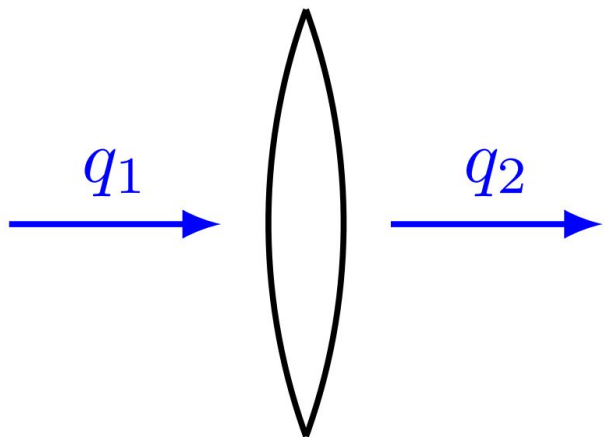
$w_0$  = Beam waist

$w(z)$  = Width of beam at some  $z$  position

$Z_R$  = Rayleigh range

$$q(z) = \left( \frac{1}{R_c} - i \frac{\lambda}{\pi w^2(z)} \right)^{-1}$$

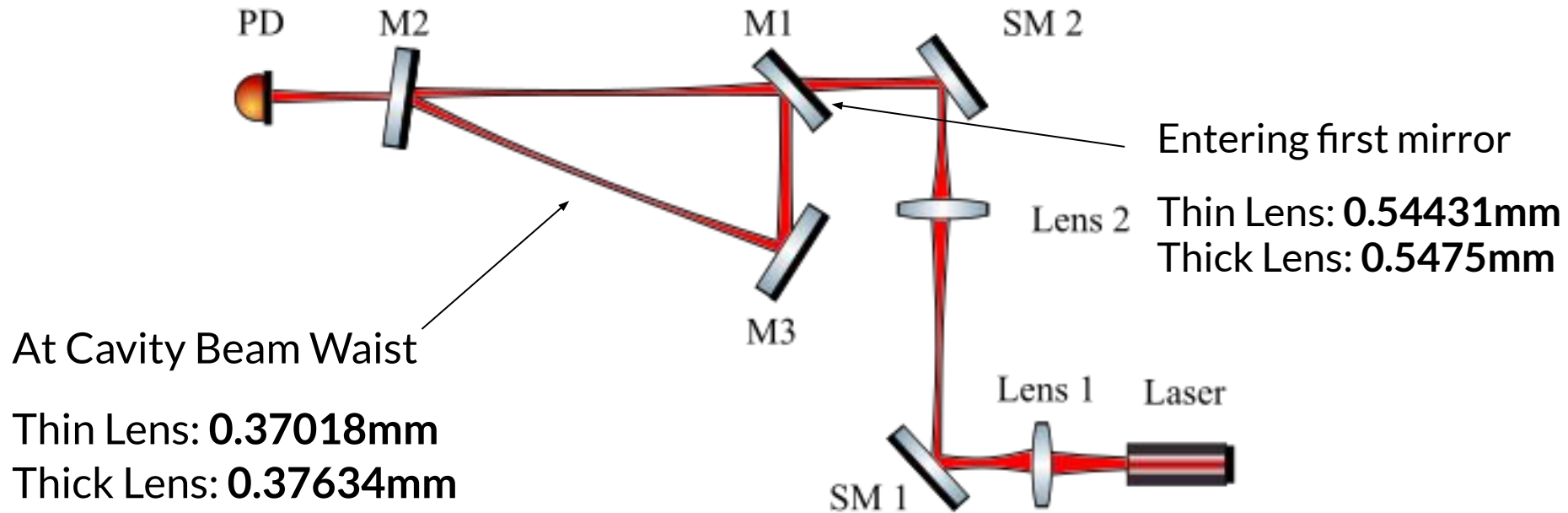
$$\begin{pmatrix} \cdot & \cdot \\ \cdot & \cdot \end{pmatrix} \begin{pmatrix} \cdot & \cdot \\ \cdot & \cdot \end{pmatrix} \begin{pmatrix} \cdot & \cdot \\ \cdot & \cdot \end{pmatrix} \begin{pmatrix} \cdot & \cdot \\ \cdot & \cdot \end{pmatrix} \cdots = \begin{pmatrix} A & B \\ C & D \end{pmatrix}$$



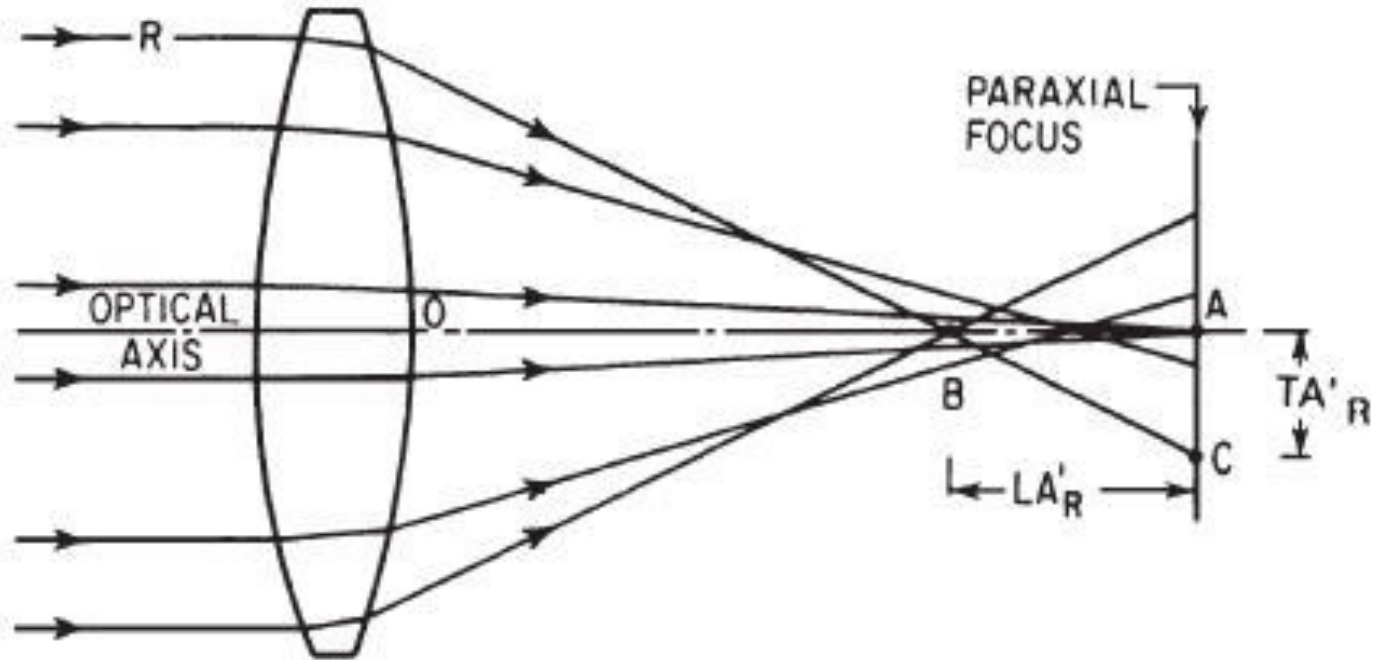
$$\frac{Aq_1 + B}{Cq_1 + D} = q_2$$



## Results from Thick Lens Calculation



# Lens Aberrations with Ray Tracing



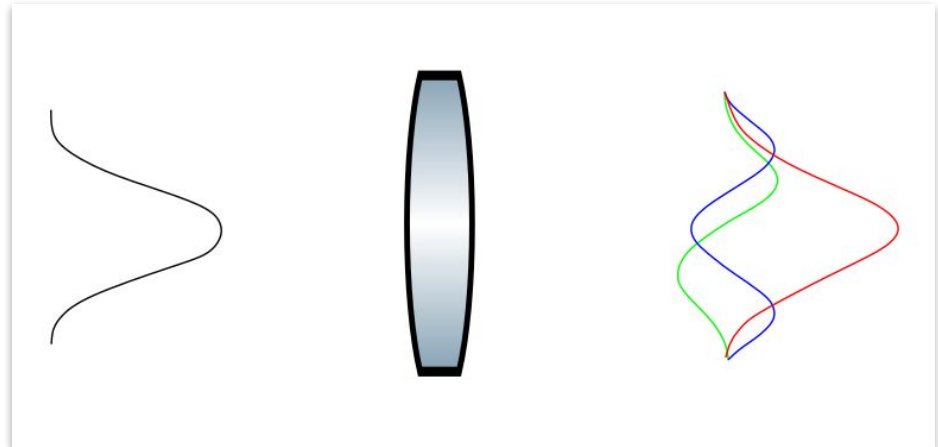
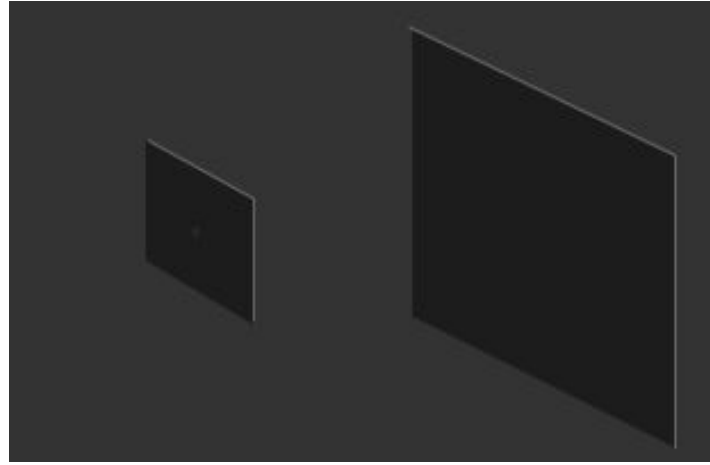
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# Ray Tracing to Gaussian Beams

A purely Gaussian laser can  
scatter into Higher Order  
Modes.

This is the Gaussian version  
of lens aberrations.

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## Quick Math Behind Gaussian Beams

$$u(x, y, z) = \sqrt{\frac{2}{\pi}} \frac{1}{\omega(z)} e^{\left(-ik \frac{x^2+y^2}{2R(z)} - \frac{x^2+y^2}{\omega(z)}\right)} e^{i\psi(z)}$$

Equations to get Hermite Gauss patterns

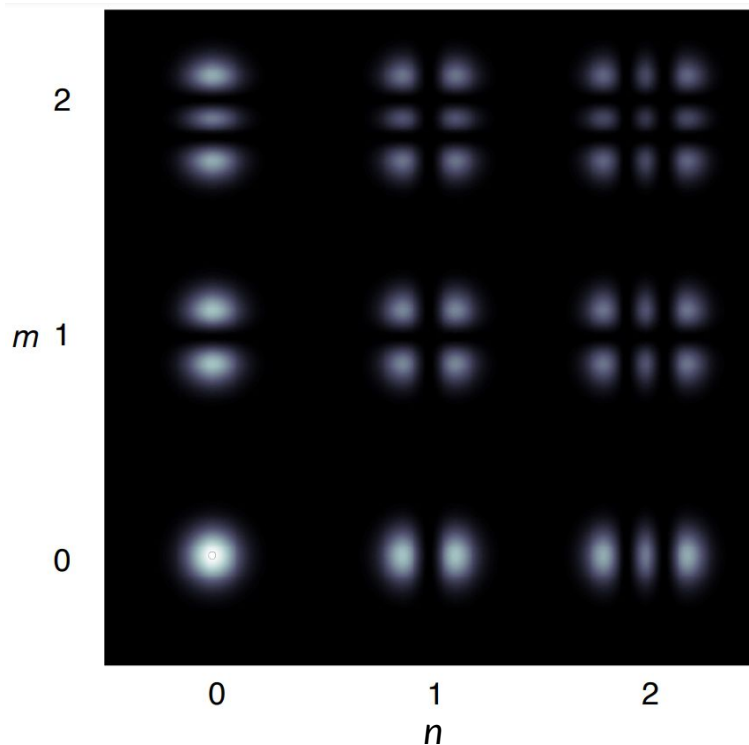
$$u_{nm} = \left(2^{n+m-1} n! m! \pi\right)^{-1} \frac{1}{\omega(z)} e^{i(n+m+1)\psi(z)} H_n \frac{\sqrt{2}x}{\omega(z)} H_m \frac{\sqrt{2}y}{\omega(z)} e^{\left(-ik \frac{x^2+y^2}{2R(z)} - \frac{x^2+y^2}{\omega(z)}\right)}$$

Equations to get Laguerre Gauss patterns

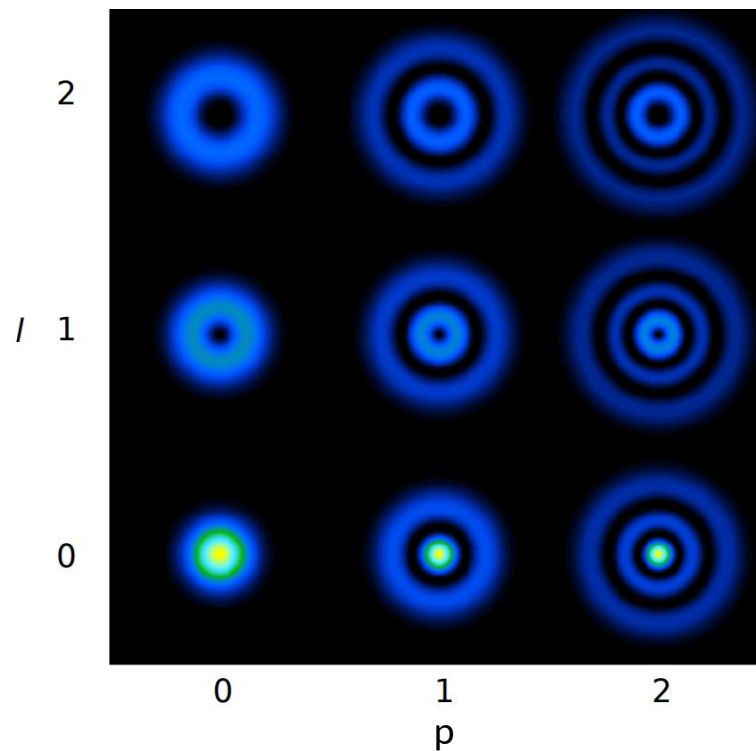
$$u_{pl} = \frac{1}{\omega(z)} \sqrt{\frac{2p!}{\pi(|l|+p)!}} e^{i(2p+|l|+1)\psi(z)} \left(\frac{\sqrt{2}r}{\omega(z)}\right)^{|l|} L_p^{|l|} \left(\frac{2r^2}{\omega(z)^2}\right) e^{\left(-ik \frac{r^2}{2q(z)} + il\phi\right)}$$

# Meet the Families of Modes

## Hermite Gauss Modes

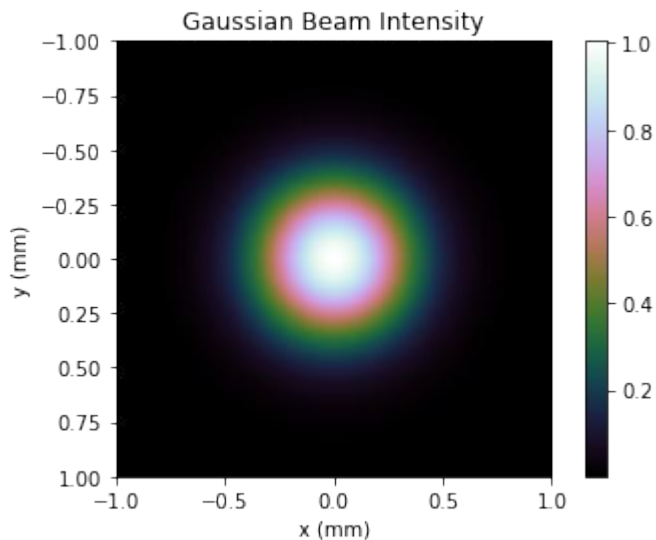


## Laguerre Gauss Modes



—  
Just add them  
together  
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$$C_{nmn'm'} = \iint_{-\infty}^{\infty} u_{nm} u_{n'm'}^* e^{(2ikZ(x,y))} dx dy$$

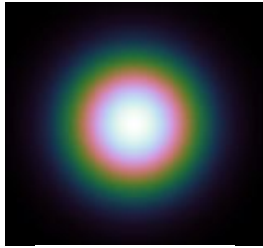


$Z(x, y)$



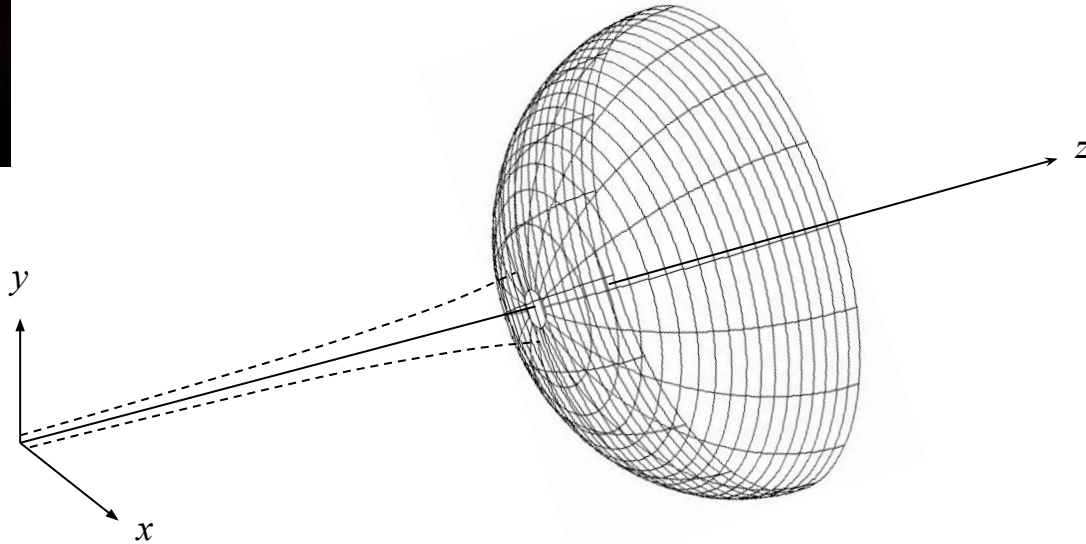
$$Z(x, y) = \sqrt{R^2 - x^2 - y^2}$$

We start with a pure Gaussian beam.  
 $p=l=0$



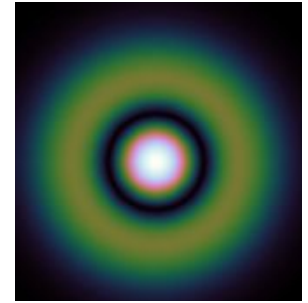
$u_{00}$

$$Z(x, y) = \sqrt{R^2 - x^2 - y^2}$$



$$|c_{0,0}|^2, |c_{1,0}|^2, |c_{2,0}|^2, \dots, |c_{n,m}|^2 \leq 1$$

The first higher order mode.  
 $p'=1, l'=0$

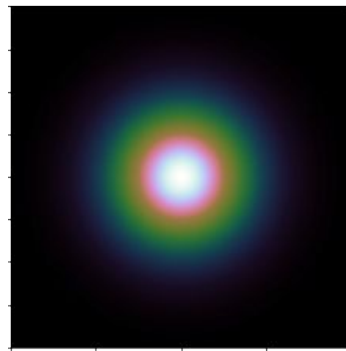
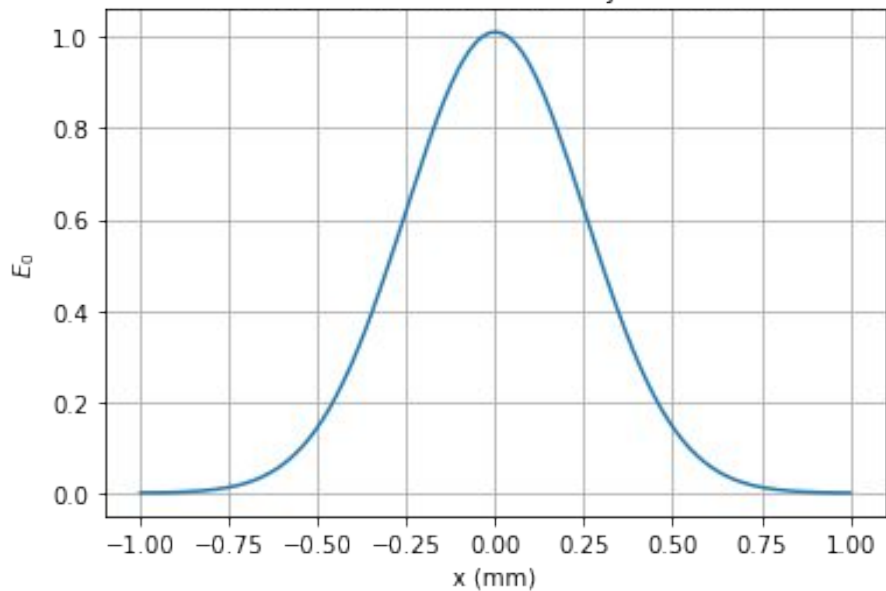


$$|c_{10}|^2 u_{10}$$

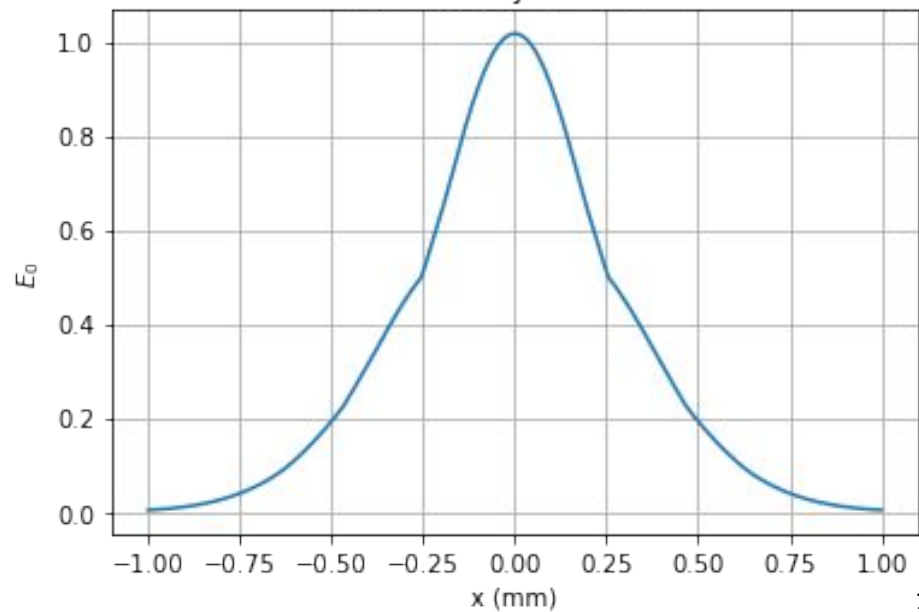
A coefficient  
 between 0 and 1



Pure Gaussian Beam Intensity Distribution

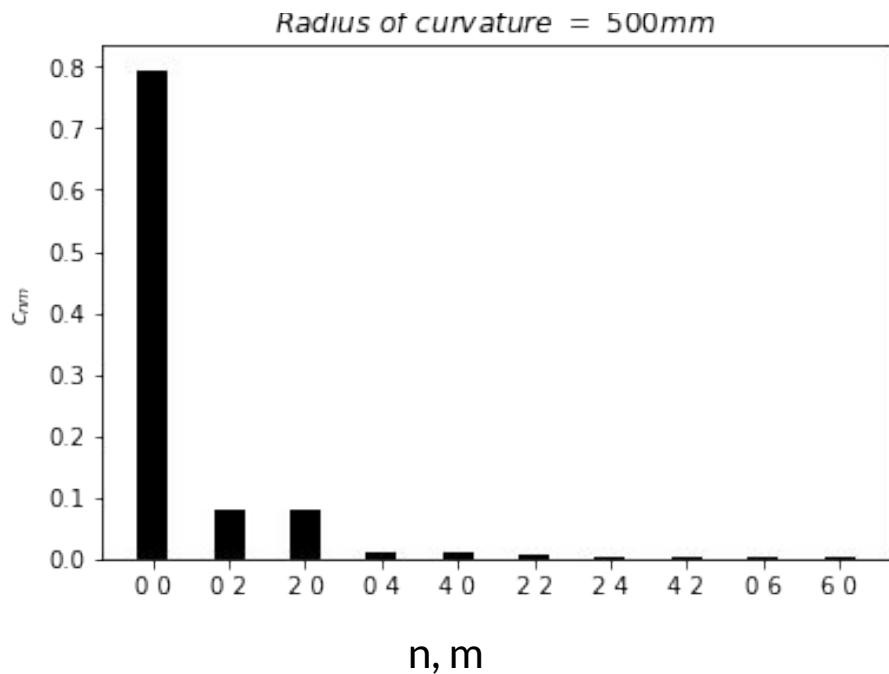


HOM Intensity Distribution

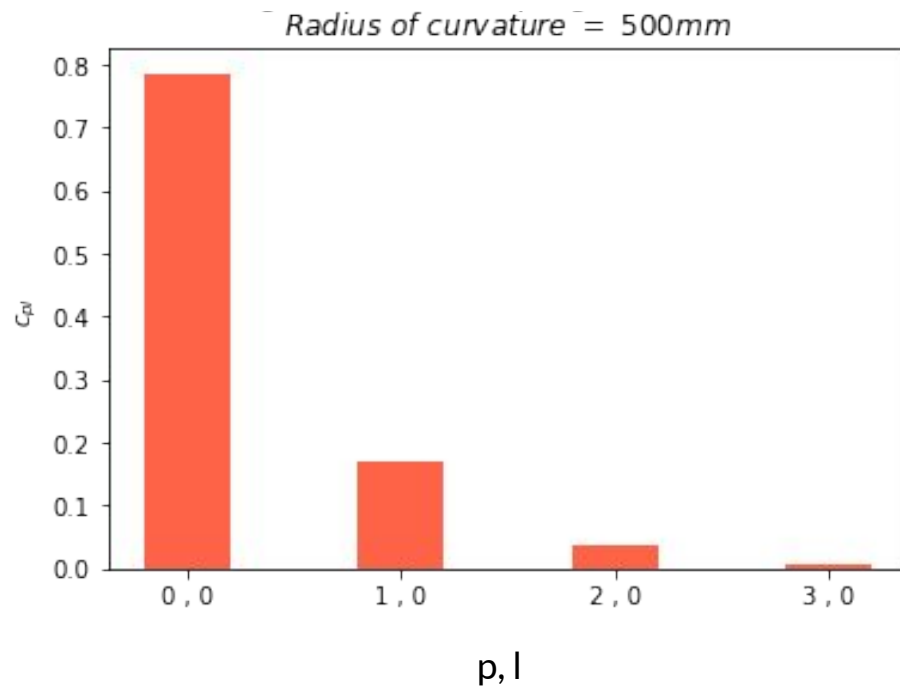




# Hermite Gauss Coupling Coefficients

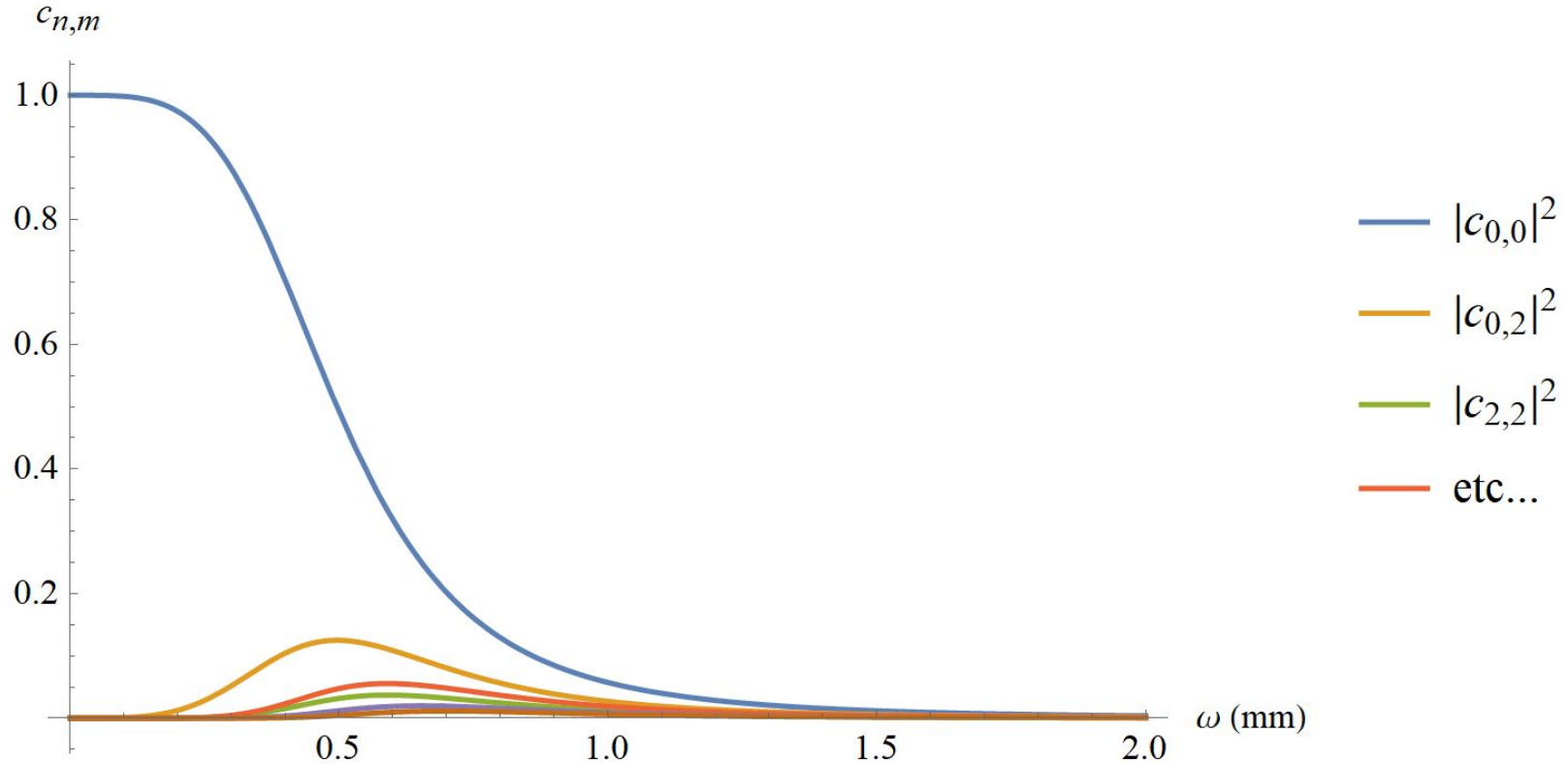


# Laguerre Gauss Coupling Coefficients

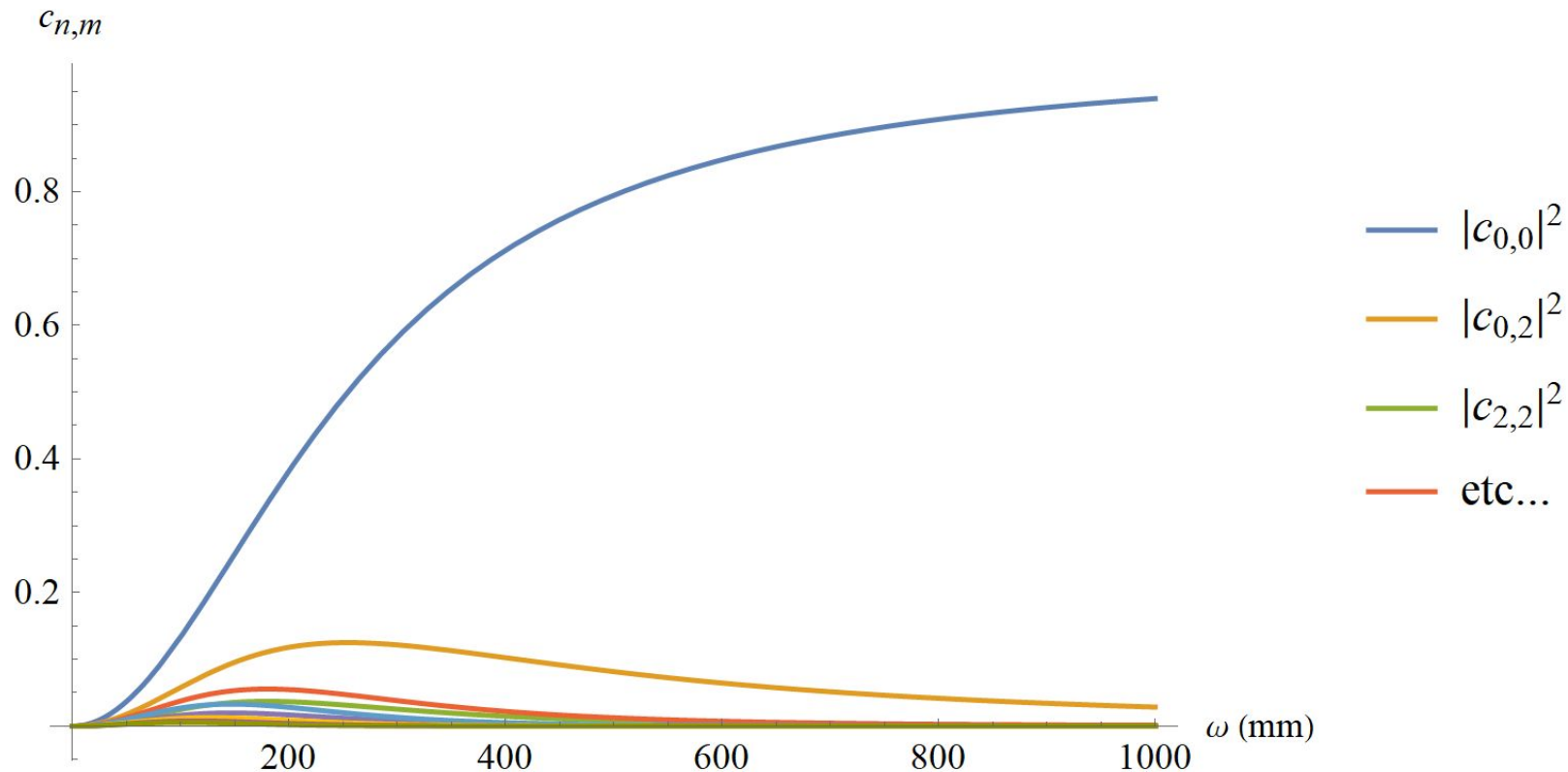


# Coupling Coefficients vs. Beam Width

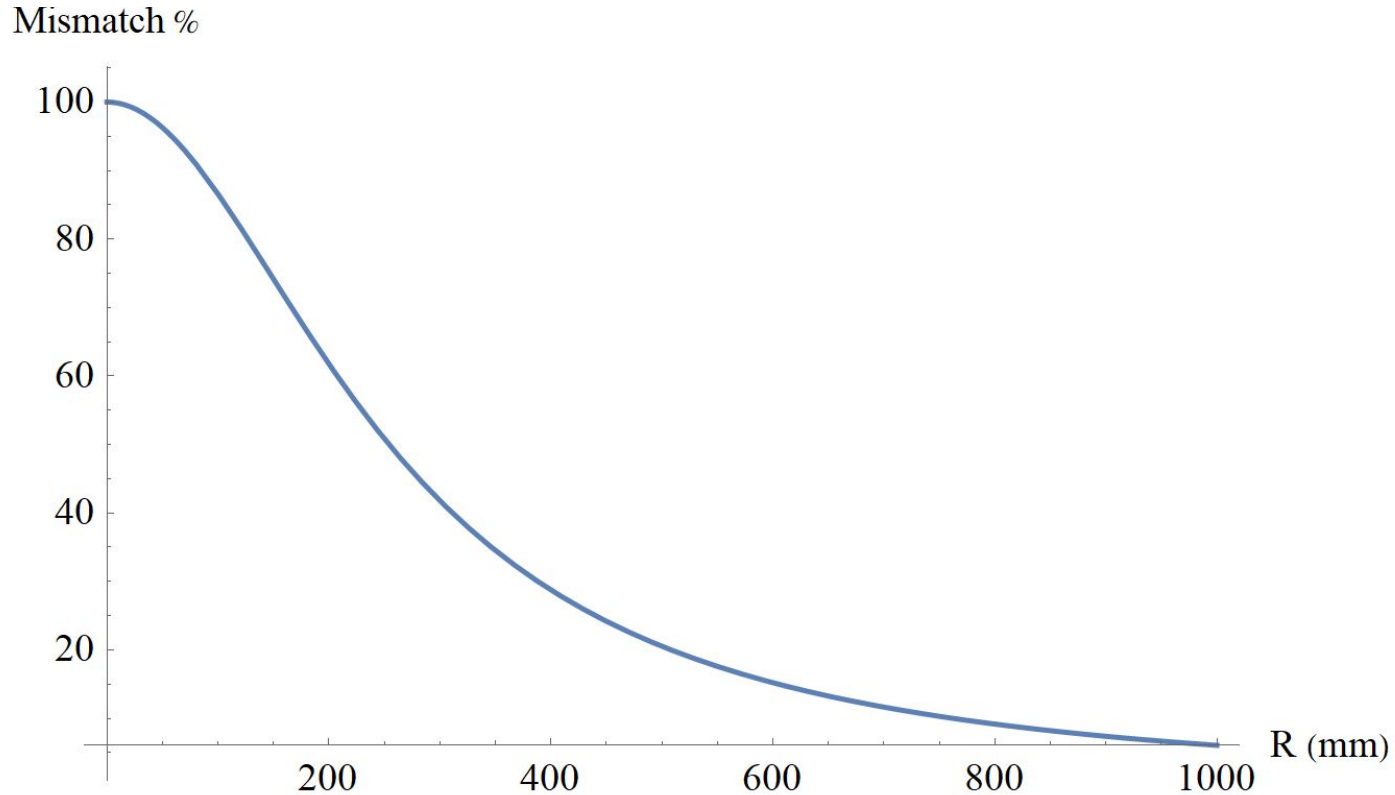
*Fixed Radius of curvature R at 500mm*



# Coupling Coefficients vs. $Z(x,y)$ Surface Radius of Curvature



# Mode Mismatch vs. Z(x,y) Surface Radius of Curvature



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# The Work Continues

- Further research on how mode mismatching affects quantum squeezing loss.
- Test different  $Z(x,y)$  functions to see how different surfaces affect beams.
- Simulate some radius  $R$ , find the  $q$  parameter and compare with ABCD matrix approach.
- Simulate a two lens system with integration method for mode matching losses.
- Find best parameters for mode matching with numerical integration.
- If work yields good results, try modeling for LIGO optics

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# References

- [1] Yuntao Bai Gautam Venugopalan Kevin Kuns Christopher Wipf Aaron Markowitz Andrew R. Wade Yanbei Chen and Rana X Adhikari, A phase-sensitive optomechanical amplifier for quantum noise reduction in laser interferometers. *Physical Review*, (2020).
- [2] Eric D. Hall, An introduction to Pound Drever Hall frequency stabilization *Physical Review*, (2001)
- [3] Xinqian Guo Linbo Zhang Jun Liu Long Chen Le Fan Guanjun Xu Tao Liu Ruifang Dong Shougang Zhang, An automatic frequency stabilized laser with herz-level linewidth *Physical Review*, (2022)
- [4] Kenneth Strain Andreas Freise Charlotte Bond Daniel Brown, Interferometer Techniques for Gravitational-Wave Detection *Physical Review*, (2015)
- [5] Denton Wu, Automated Laser Frequency Re-Stabilization *Physical Review* 2017,(2018)
- [6] K Huang H Le Jeannic J Ruaudel O Morin J Laurat, Microcontroller based Locking in optics *Physical Review*, (2014)
- [7] <https://www.liquidinstruments.com/products/integrated-instruments/laser-lock-box-mokulab/>
- [8] Jenkins A Francis, Harvey E. White, *Fundamentals of Optics* (chp. 8-9). McGraw-Hill Inc. 1957
- [9] Alex Abramovici, Chapsky Jake. *Feedback Control Systems: A Fast-Track Guide for Scientists and Engineers*, Kluwer Academic Publishers, 2000
- [10] Peter Beyersdorf, “Thick Lenses and ABCD Formalism”, Lecture notes, 2006.

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*Mathematica*



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