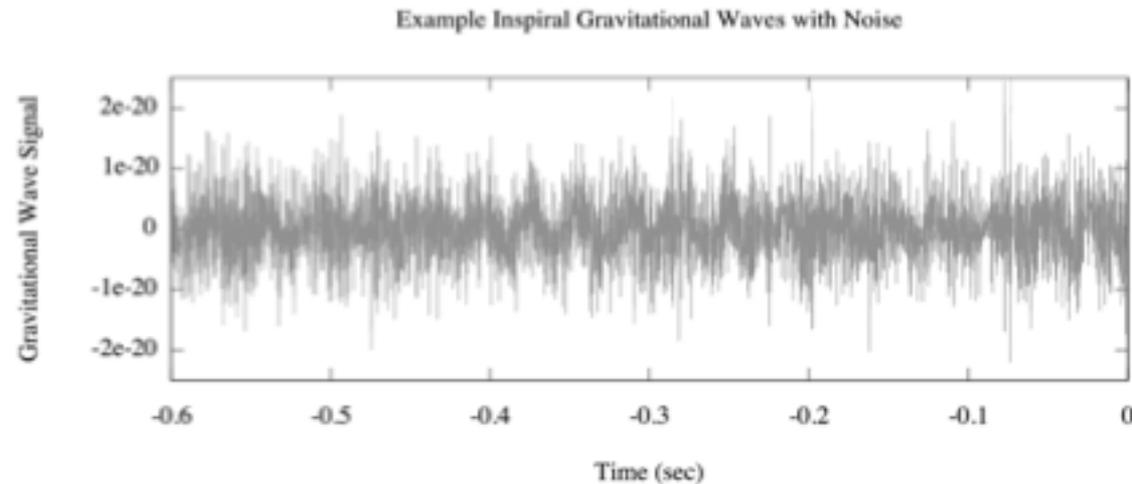
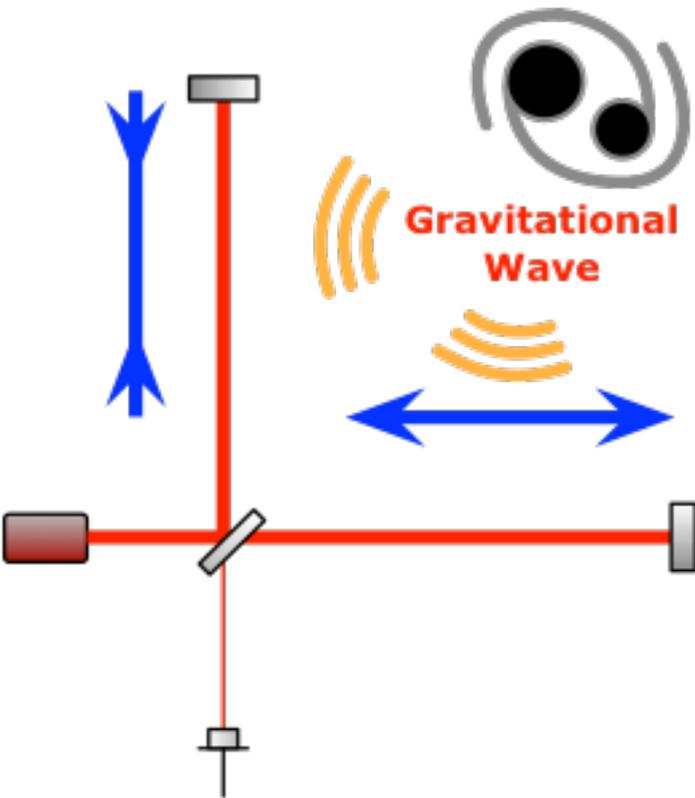


# Noises in gravitational wave detectors

Koji Arai – LIGO Laboratory / Caltech

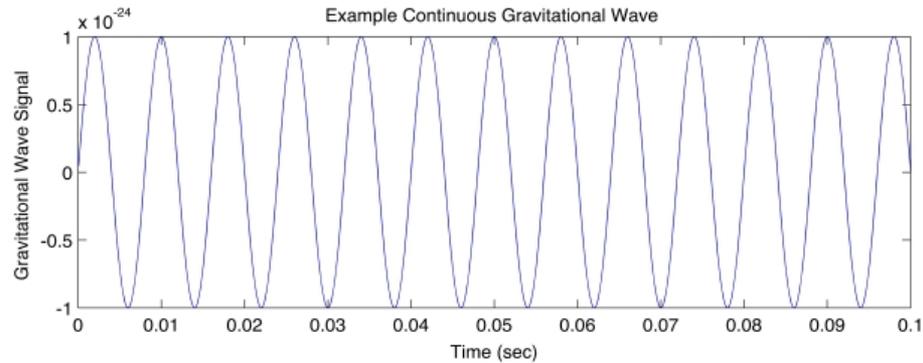
# Detector output signal

- $h(t)$  = differential arm strain signal
- Single time series data

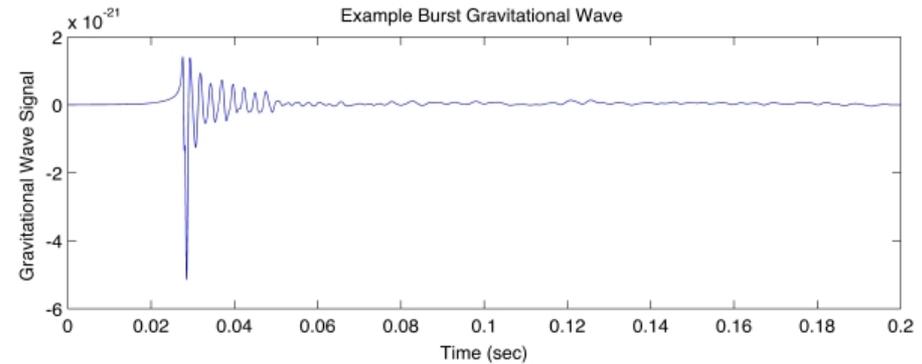


# Signal vs Noise

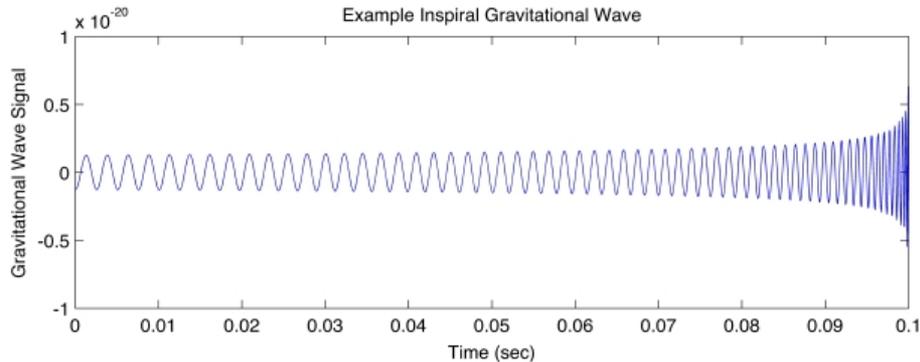
- **Waveform of GW signals**
- **Continuous**



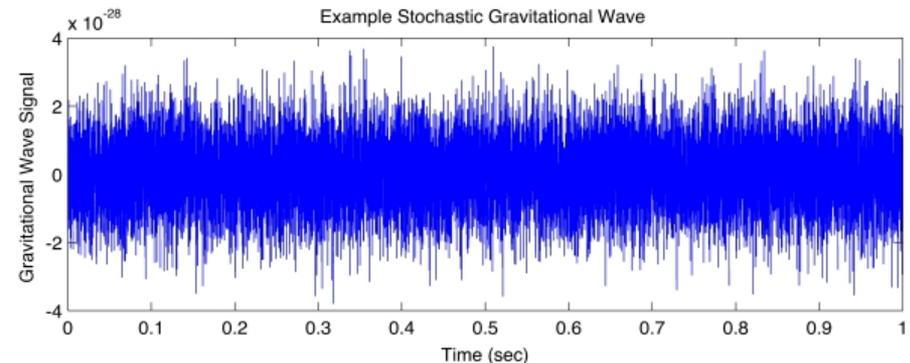
- **Burst** (unpredictable w.f.)



- **Compact Binary Coalescence**

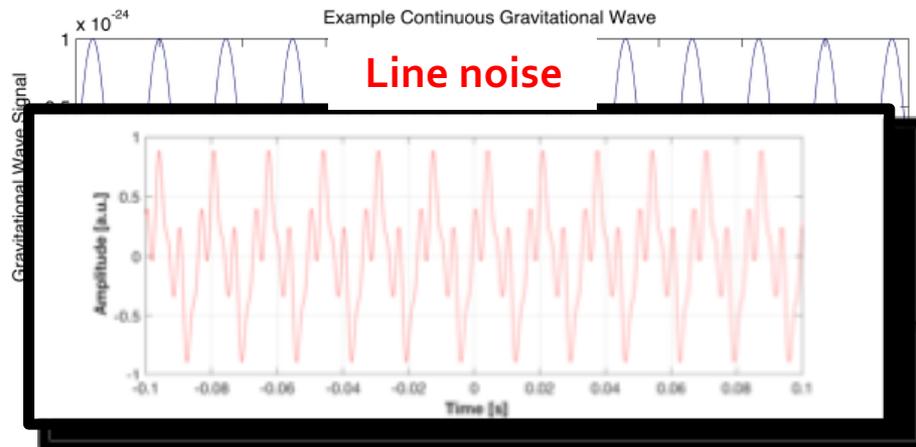


- **Stochastic**

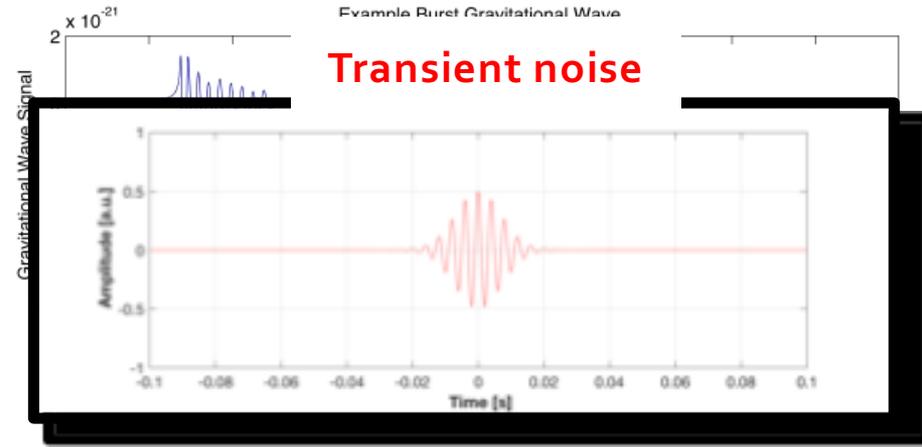


# Signal vs Noise

- Waveform of GW signals vs noises
- Continuous

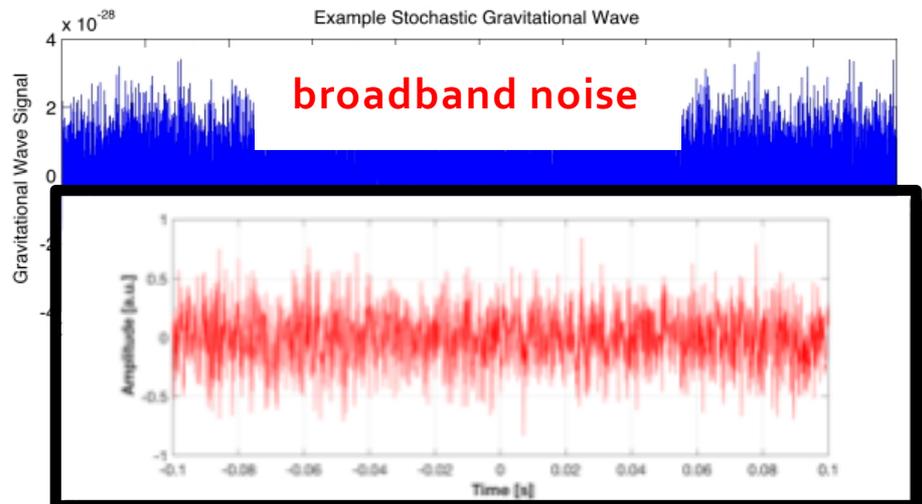
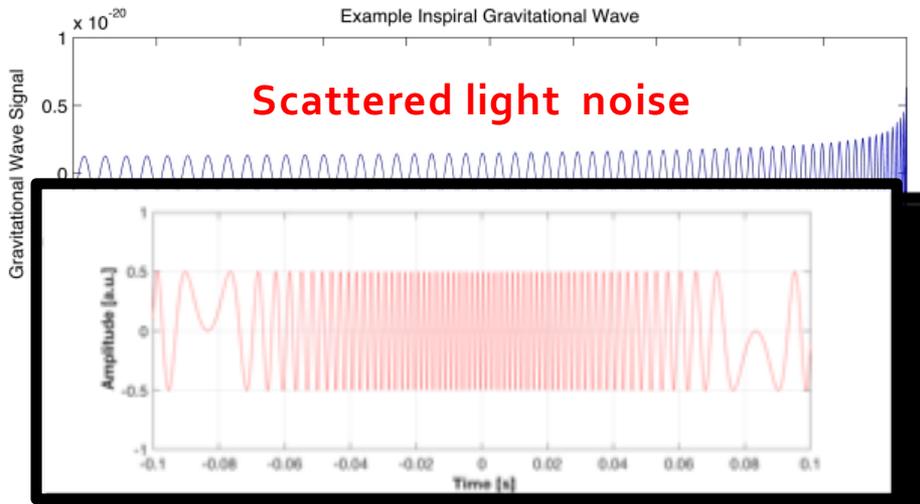


- Burst (unpredictable w.f.)



- Compact Binary Coalescence

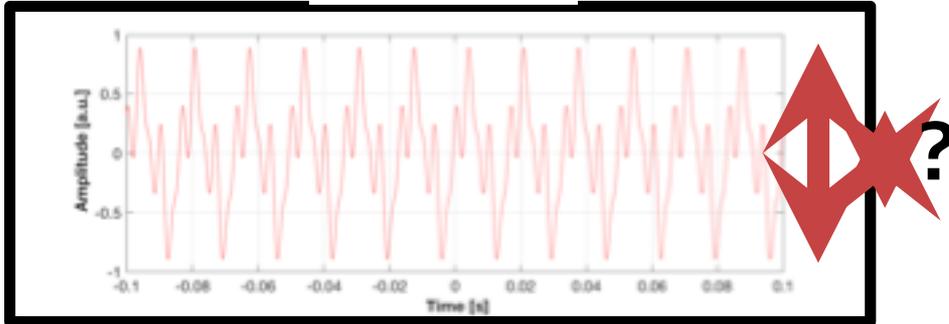
- Stochastic



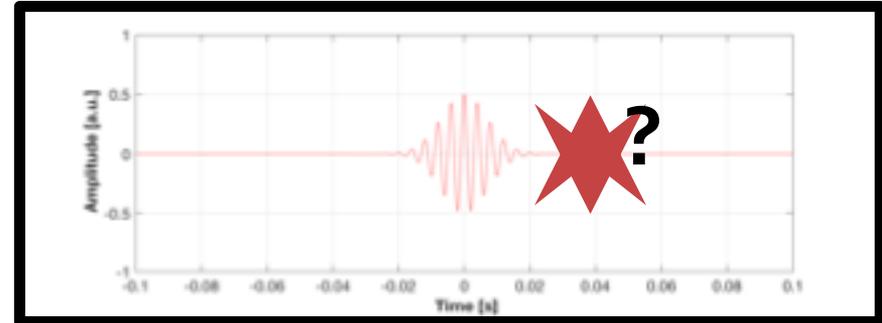
# Characterization of time-series data

## ■ Amplitude

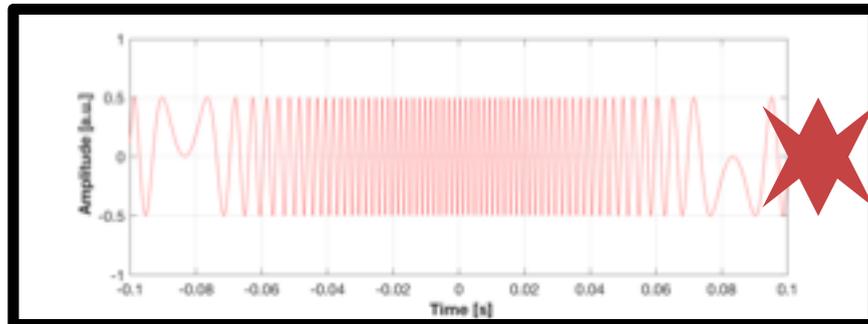
Line noise



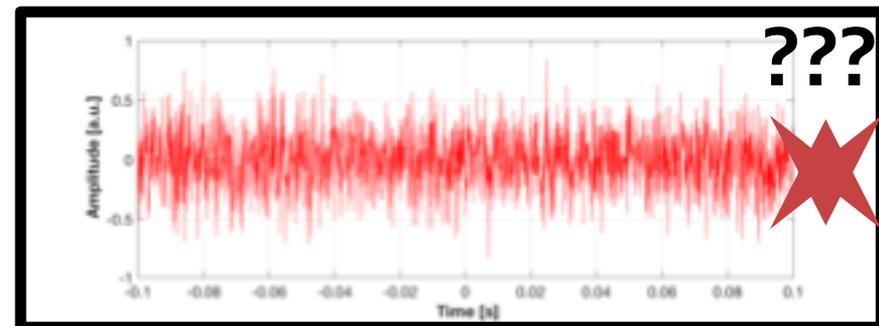
Transient noise



Scattered light noise



broadband noise

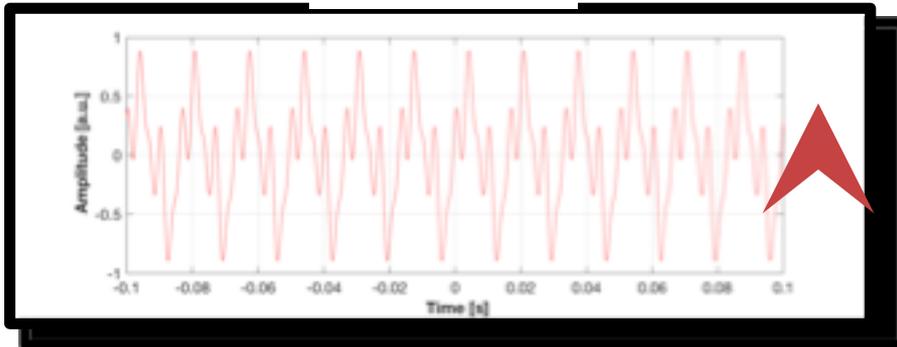


# Characterization of time-series data

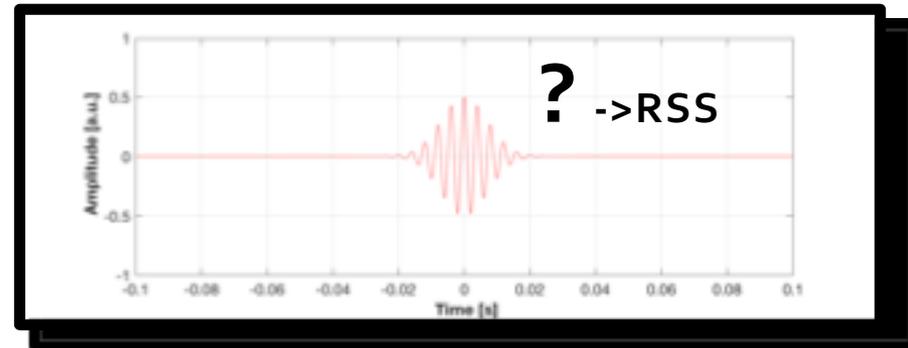
- RMS (Root Mean Square)

$$x_{\text{RMS}} = \sqrt{\frac{1}{T} \int_{T_1}^{T_2} [x(t)]^2 dt}$$

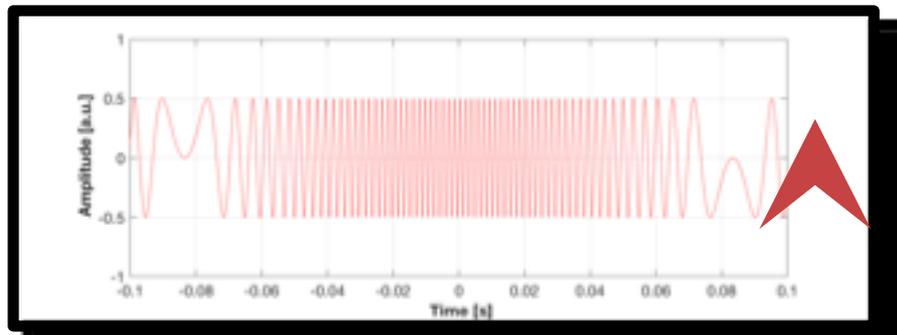
Line noise



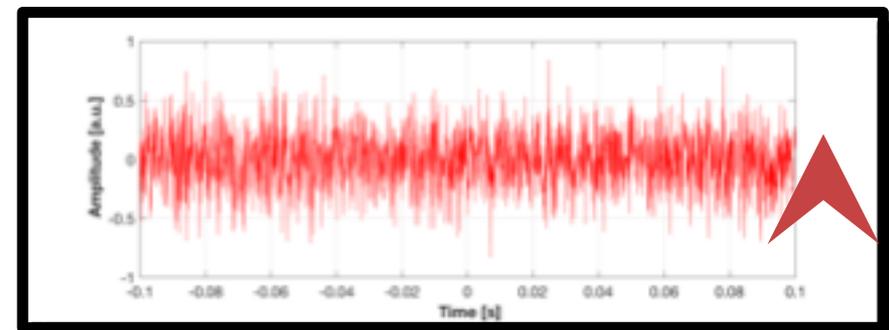
Transient noise



Scattered light noise



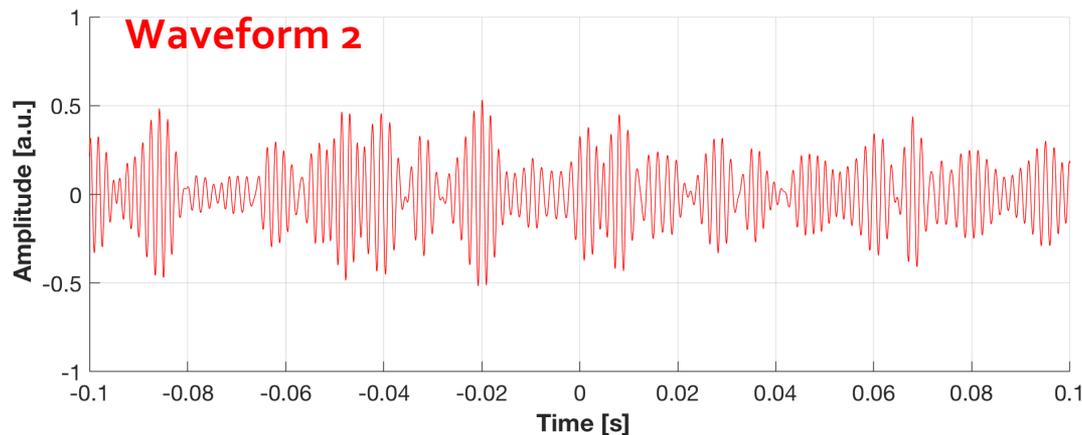
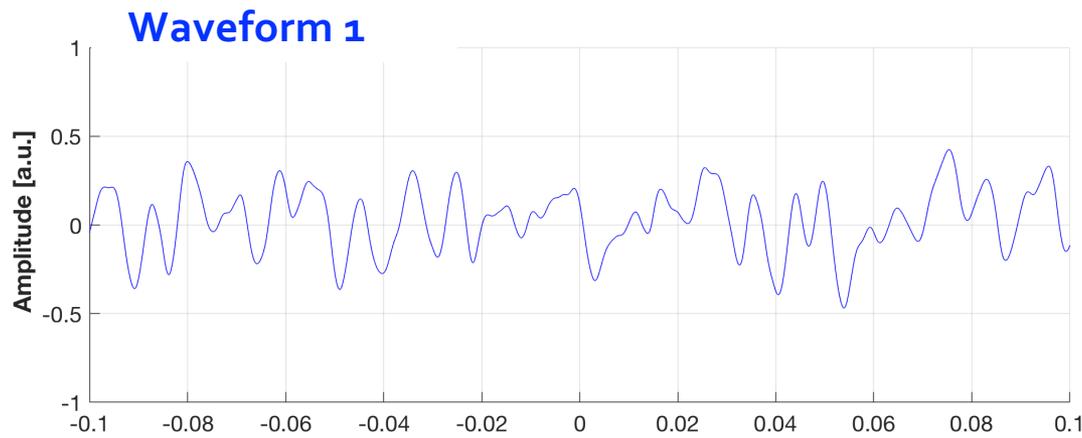
broadband noise



# Characterization of time-series data

- Same RMS ( $\sim 0.2$ ), but they look different

$$x_{\text{RMS}} = \sqrt{\frac{1}{T} \int_{T_1}^{T_2} [x(t)]^2 dt}$$



# Freq domain: Power Spectral Density

- Double sided PSD (-Infinity < f < Infinity)

$$S_{\text{DS}}(f) = \lim_{T \rightarrow \infty} \frac{1}{T} \left| \int_{-T/2}^{T/2} x(t) e^{-2\pi i f t} dt \right|^2$$

- Single sided PSD (0 <= f < Infinity)

$$S_x(f) = 2S_{\text{DS}}(f) \quad [\text{x}_{\text{unit}}^2 / \text{Hz}]$$

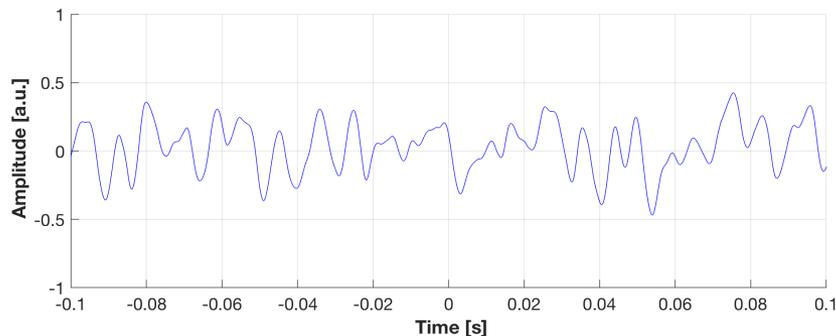
- Linearized PSD:

$$G_x(f) = \sqrt{S_x(f)} \quad [\text{x}_{\text{unit}} / \text{sqrtHz}]$$

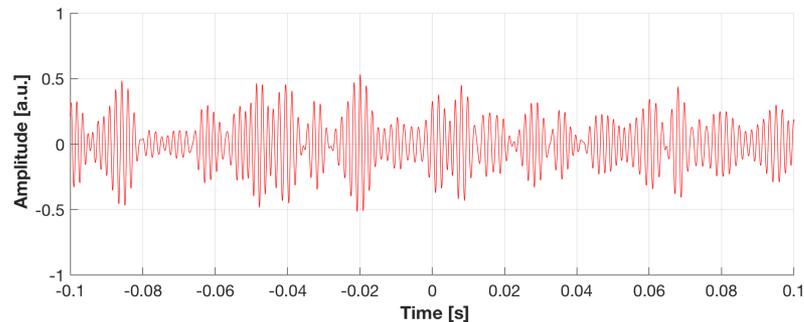
# Characterization of time-series data

- Same RMS ( $\sim 0.2$ ), but they look different

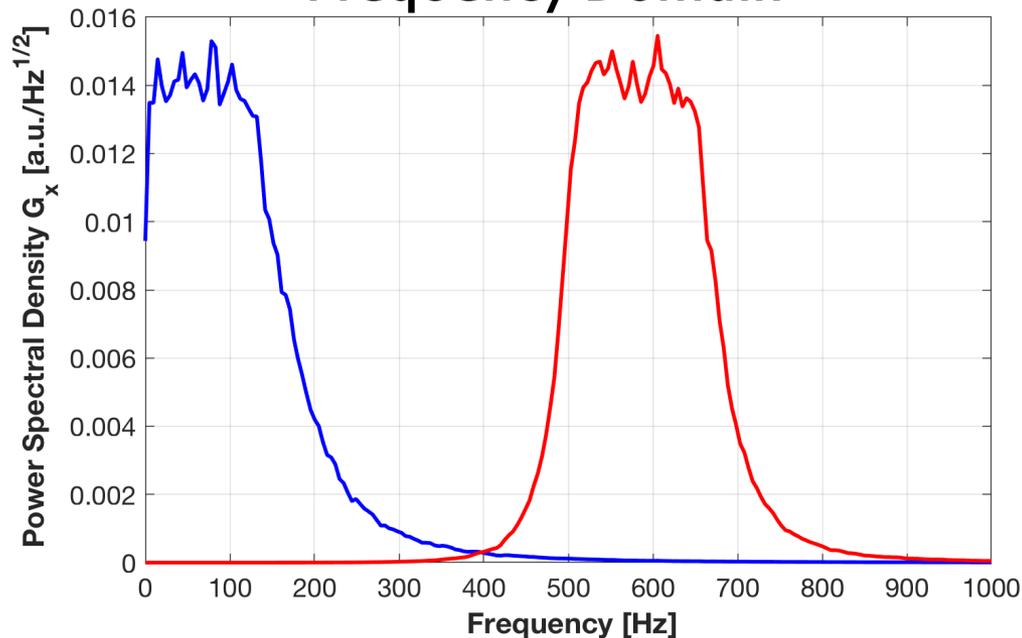
Waveform 1



Waveform 2



## Frequency Domain



# PSD $\leftrightarrow$ RMS

- Parseval's Theorem for signal RMS and PSD

$$\begin{aligned}\overline{x^2(t)} &= \int_0^{\infty} S_x(f) df \\ &\equiv x_{\text{RMS}}^2\end{aligned}$$

Root Mean of  $x(t)$ :

average signal power density (per sec)  
(cf. variance, std deviation)

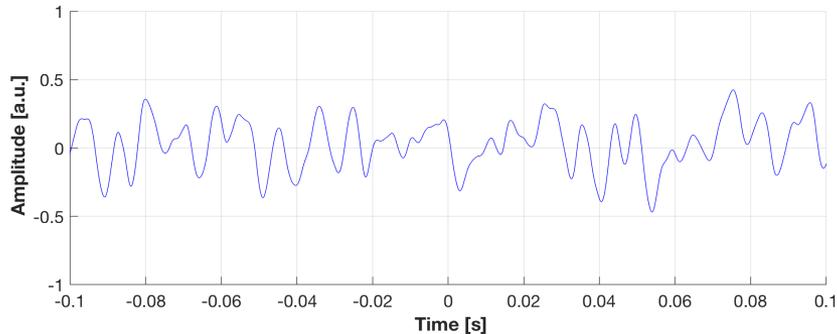
PSD  $S_x(f)$ :

power density per frequency (per sec)

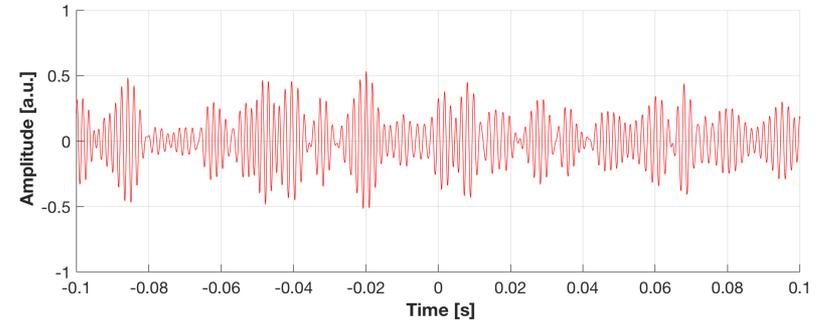
# PSD $\leftrightarrow$ RMS

- Same RMS ( $\sim 0.2$ ), but they look different

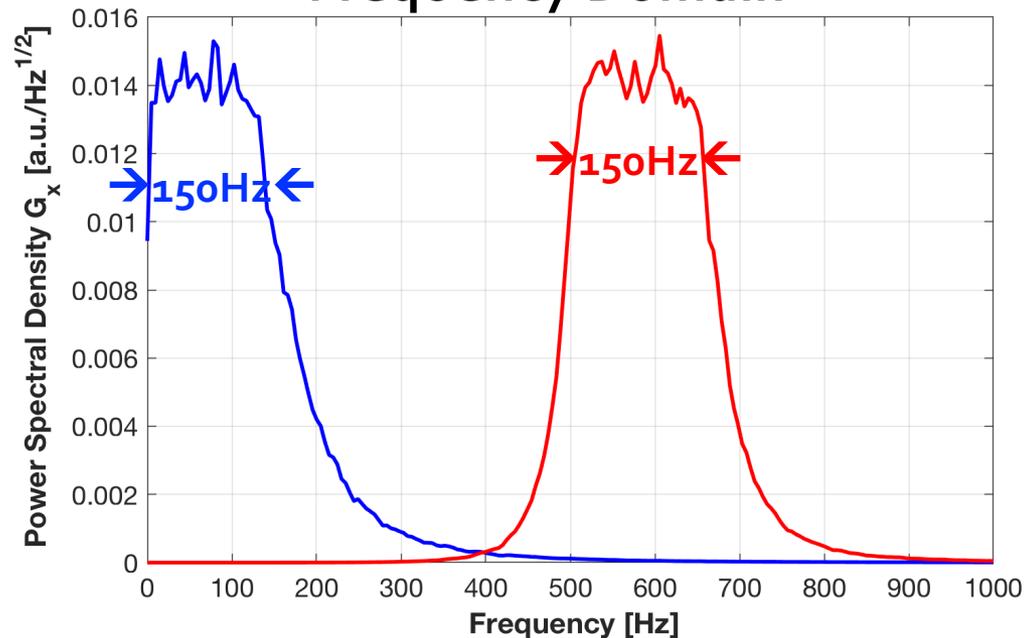
Waveform 1: RMS  $\sim 0.18$



Waveform 2: RMS  $\sim 0.18$



## Frequency Domain



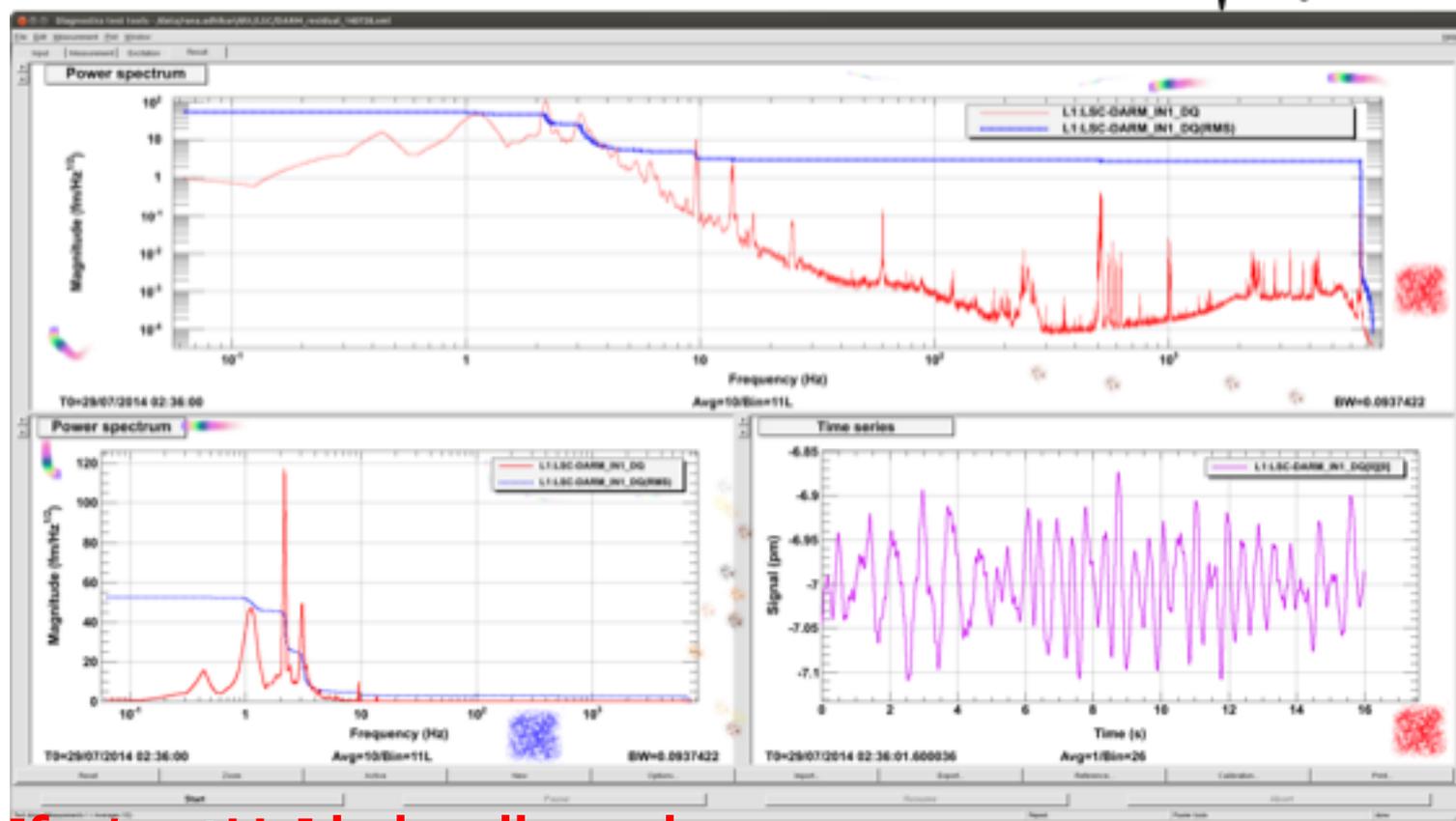
$$\text{Sqrt}(0.014^2 * 150) \\ = \sim 0.17$$

# Time series, PSD, and RMS

## Example

PSD [fm/sqrtHz] in log-log scale,

$$x_{RMS}(f) = \sqrt{\int_f^\infty S_x(f')df'}$$



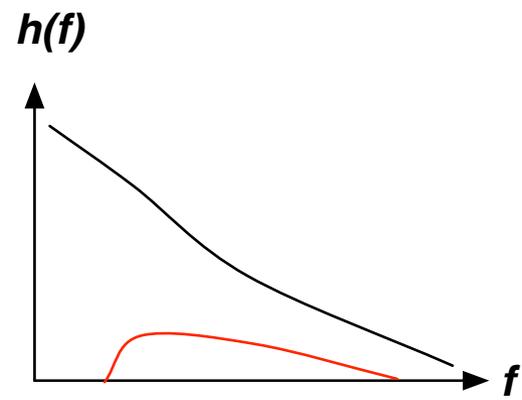
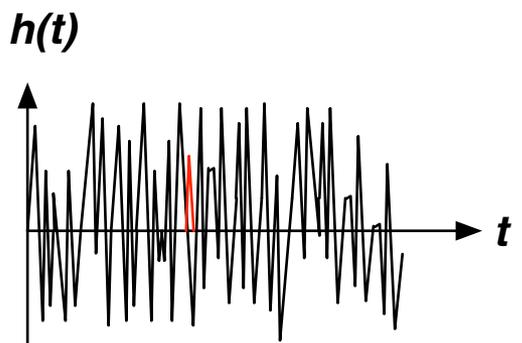
PSD [fm/sqrtHz] in log-lin scale

RMS [fm] ~ 50fm = 0.05pm

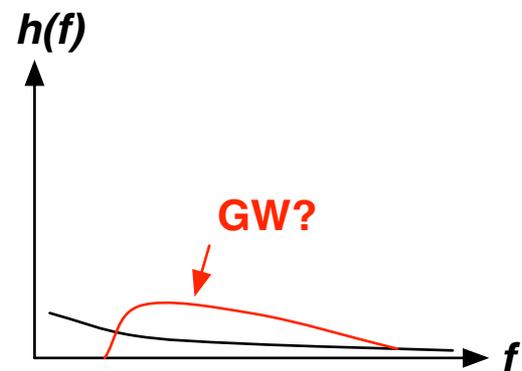
Time series [pm]

# Time domain vs frequency domain

- Time domain vs frequency domain



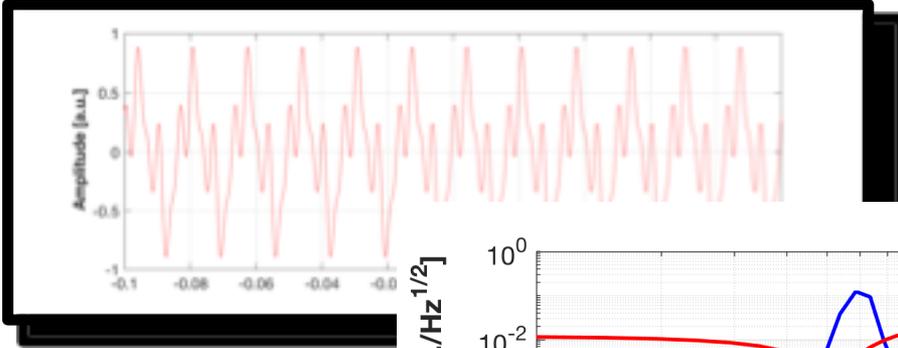
Noise Reduction



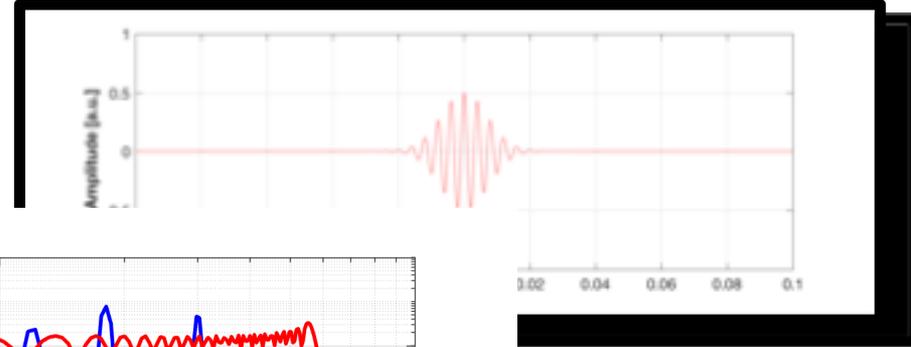
- Time domain: transient noises
- Frequency domain: stationary noises

# Time domain vs frequency domain

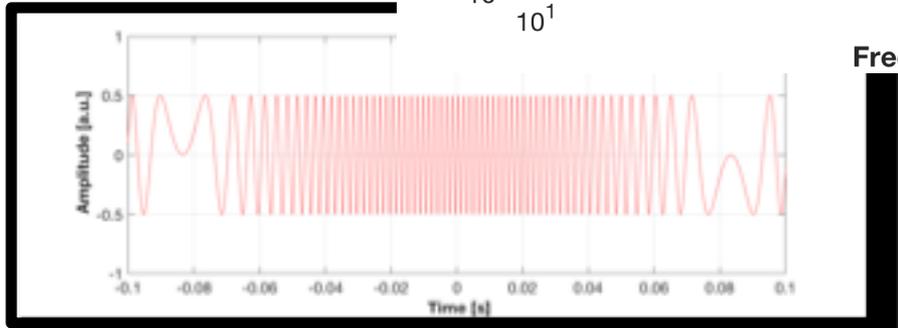
Line noise \*



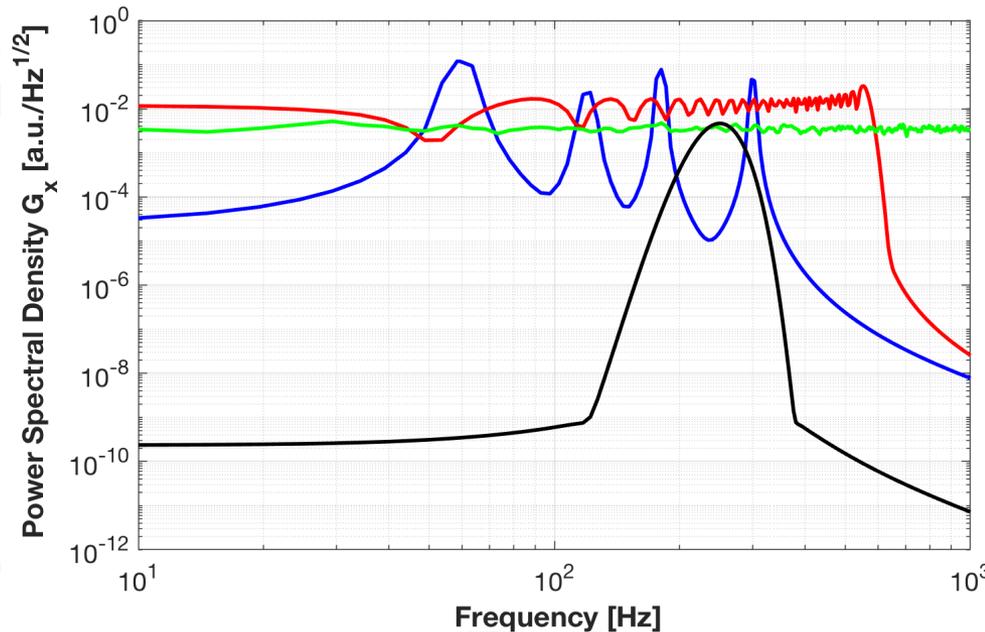
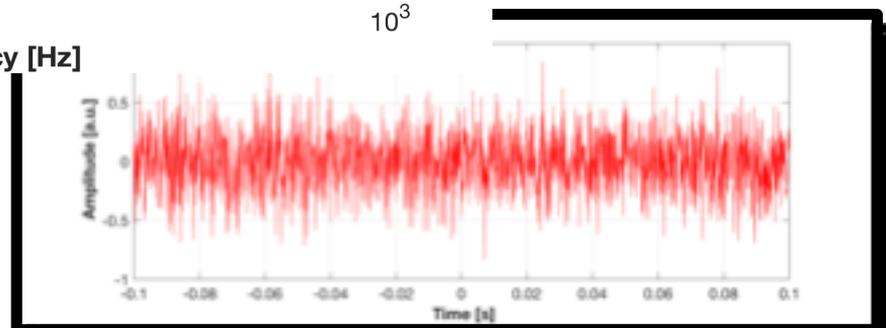
Transient noise \*



Scattered light noise \*

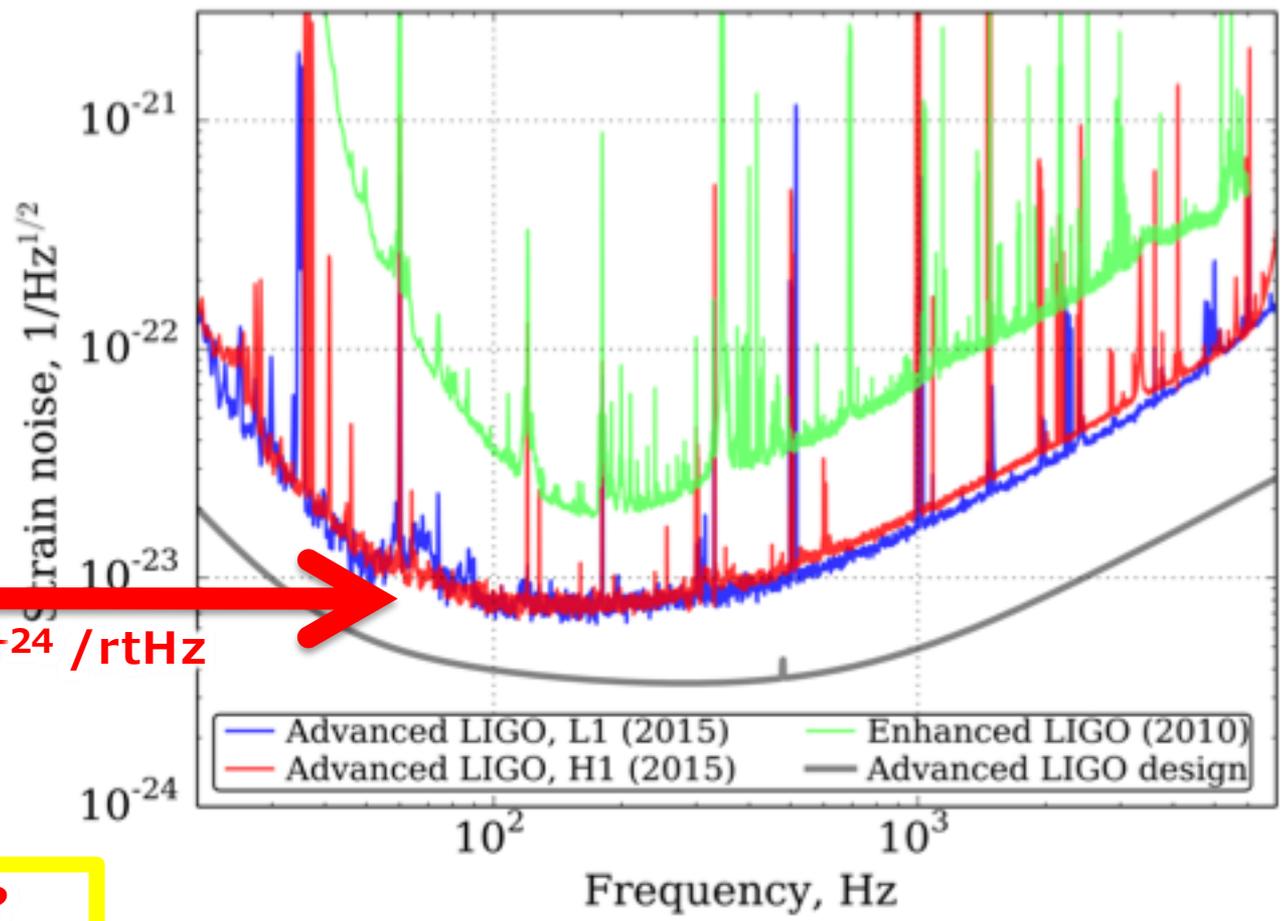


Broadband noise



# Sensitivity and noise

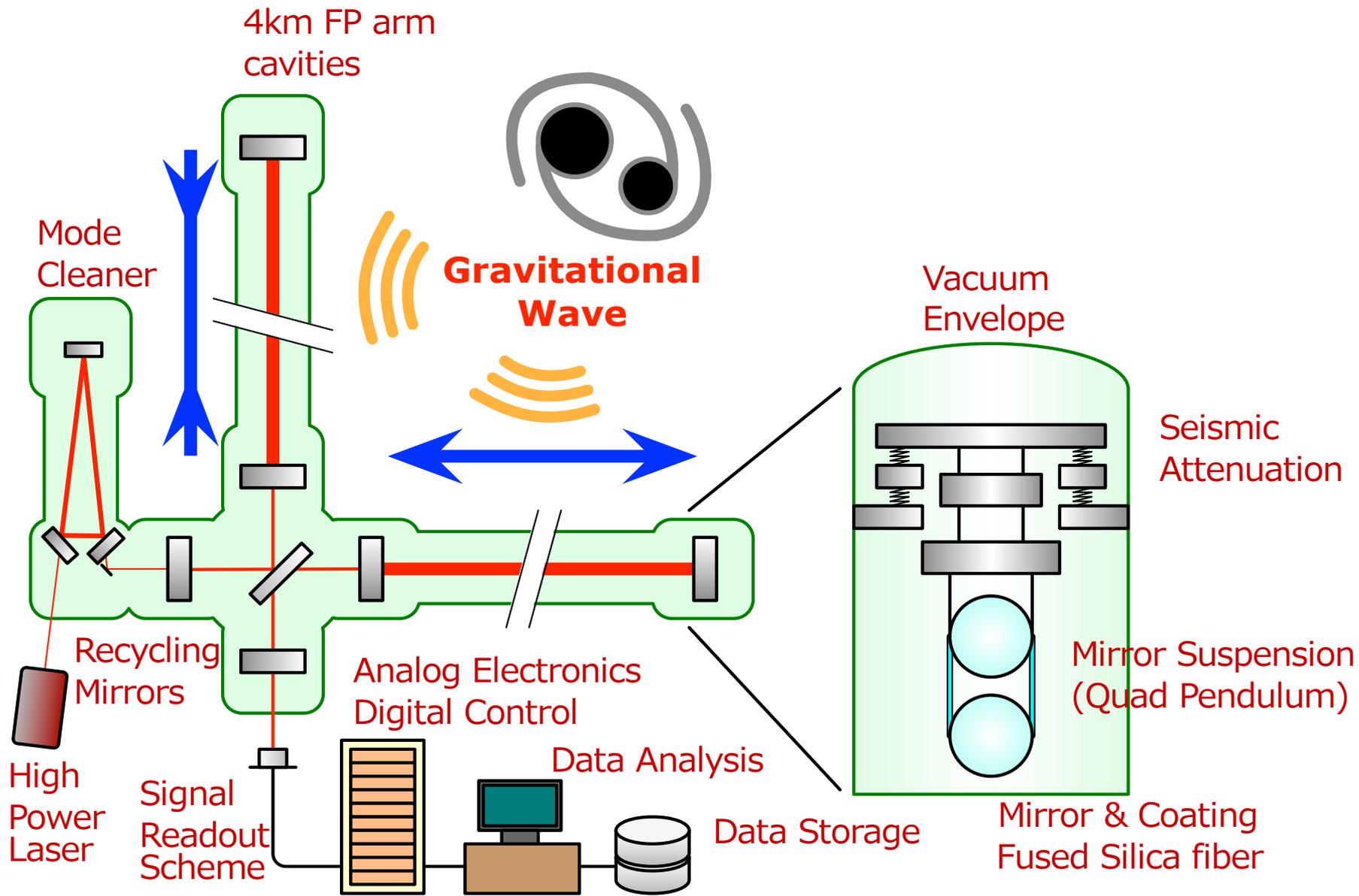
- Sensitivity (=noise level) of Advanced LIGO
- Current sensitivity



$h = 8 \times 10^{-24} / \text{rtHz}$

**RMS?**

# Actual GW detector

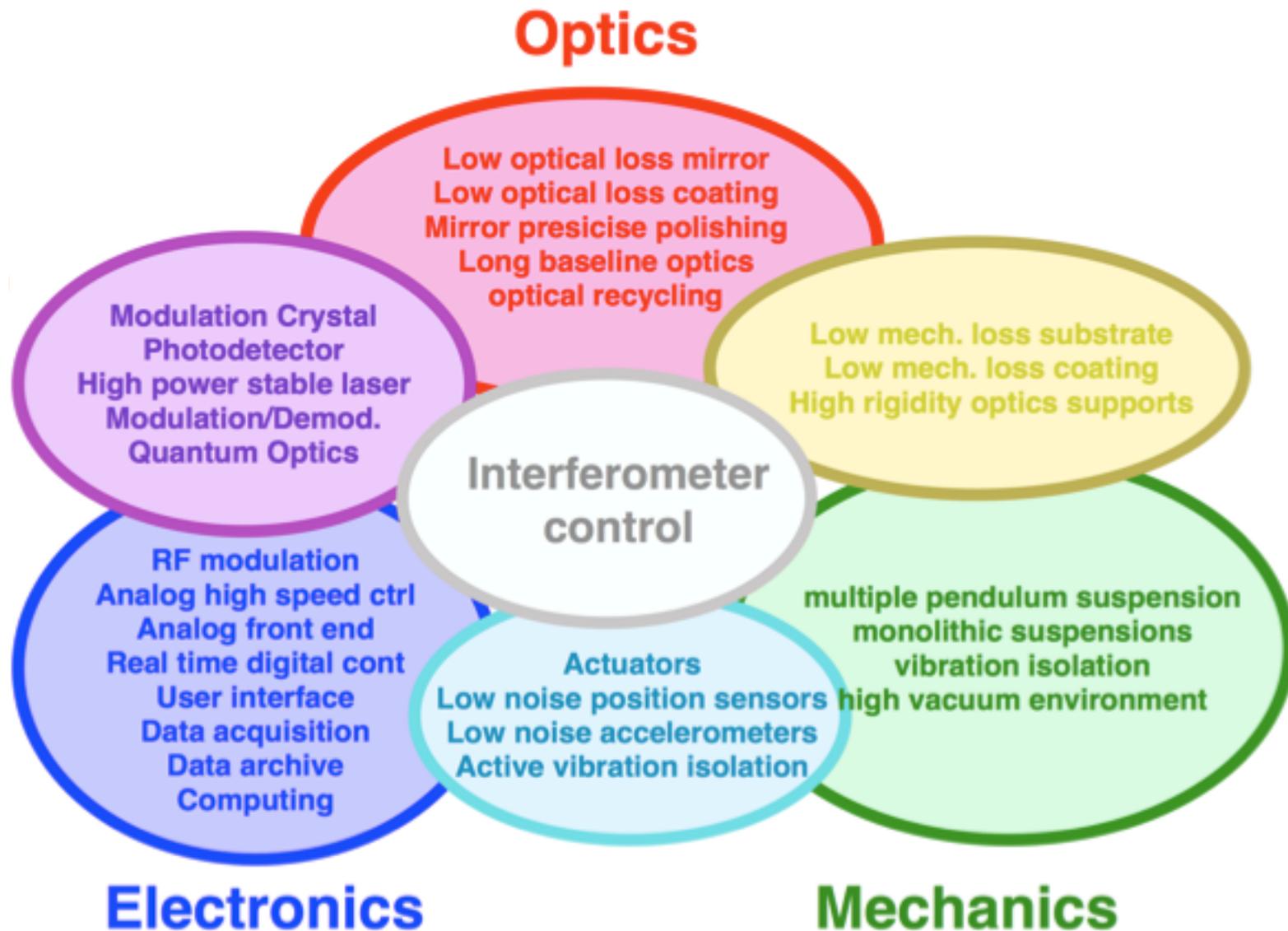


# Components of the interferometer

- 3 elements of a GW detector
  - Mechanics
  - Optics
  - Electronics

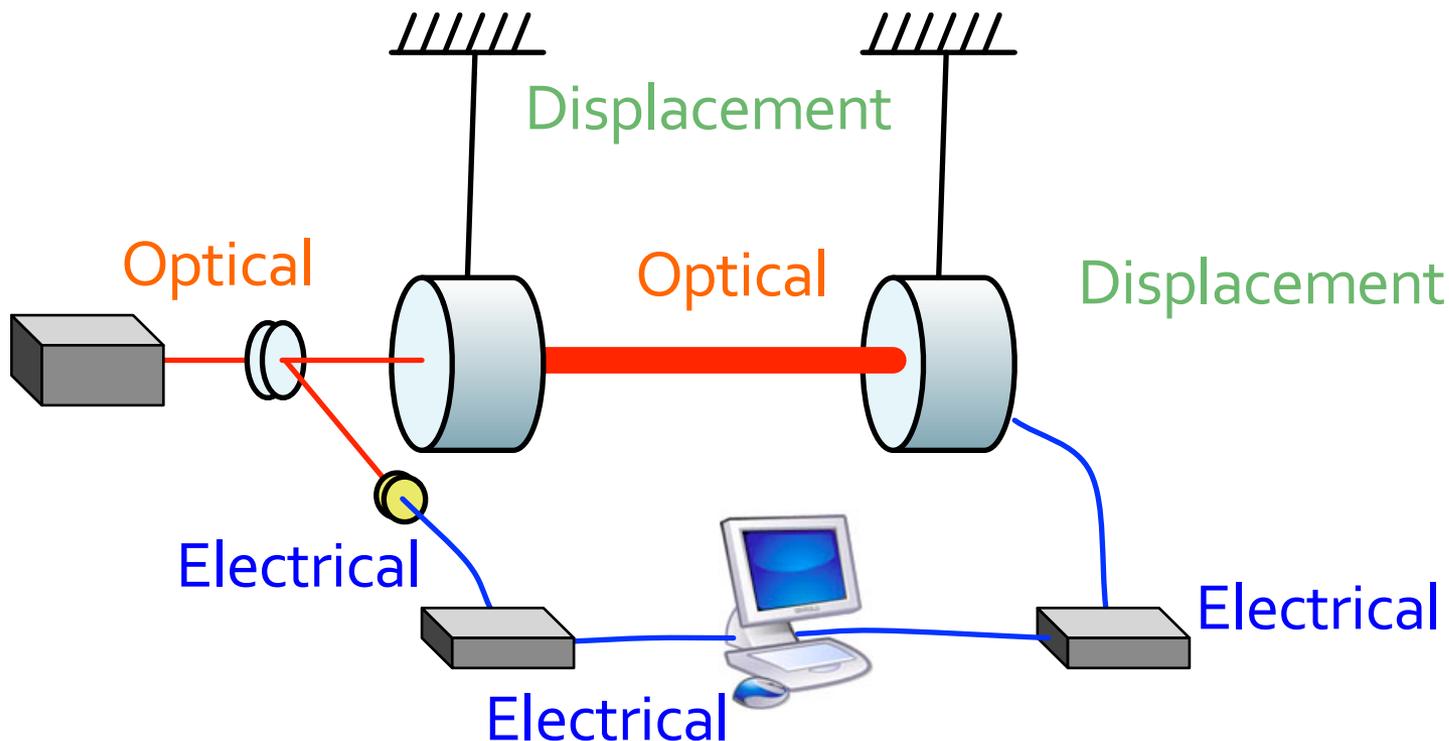
# Components of the interferometer

■ 3



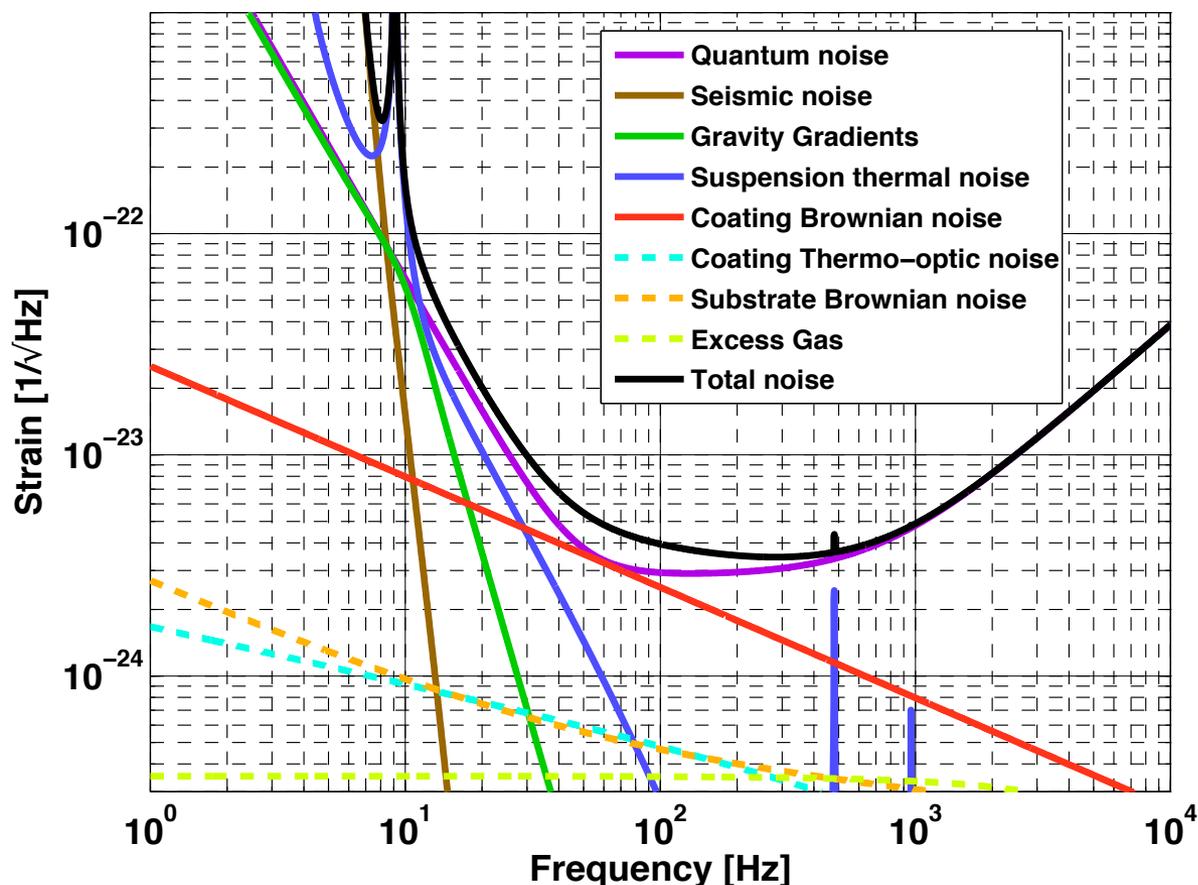
# Noise categories

- 3 fundamentals of the GW detector
- **Mechanics** -> **Displacement noises**
- **Optics** -> **Optical noises**
- **Electronics** -> **Electrical noises**



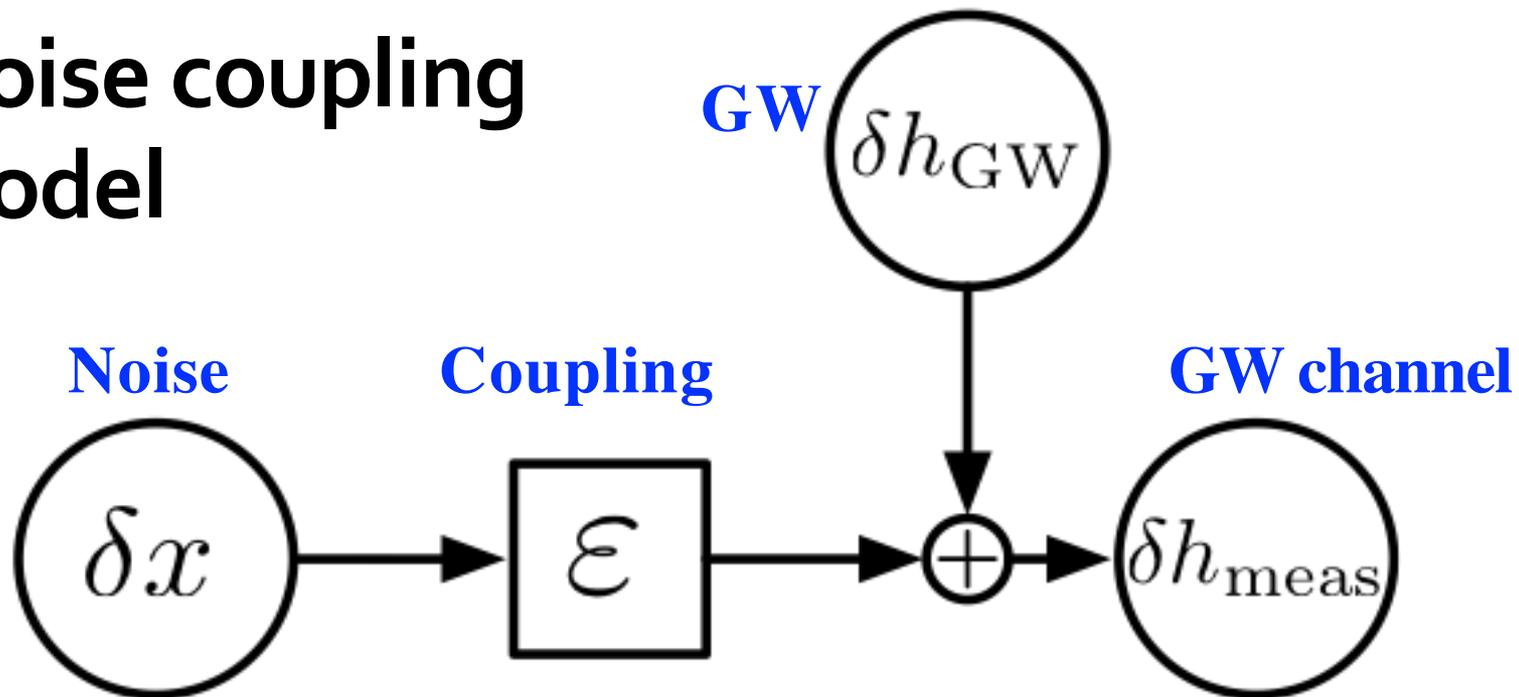
# Sensitivity and noise

- Sensitivity (=noise level) of Advanced LIGO
- Design



# Noise coupling to the GW channel

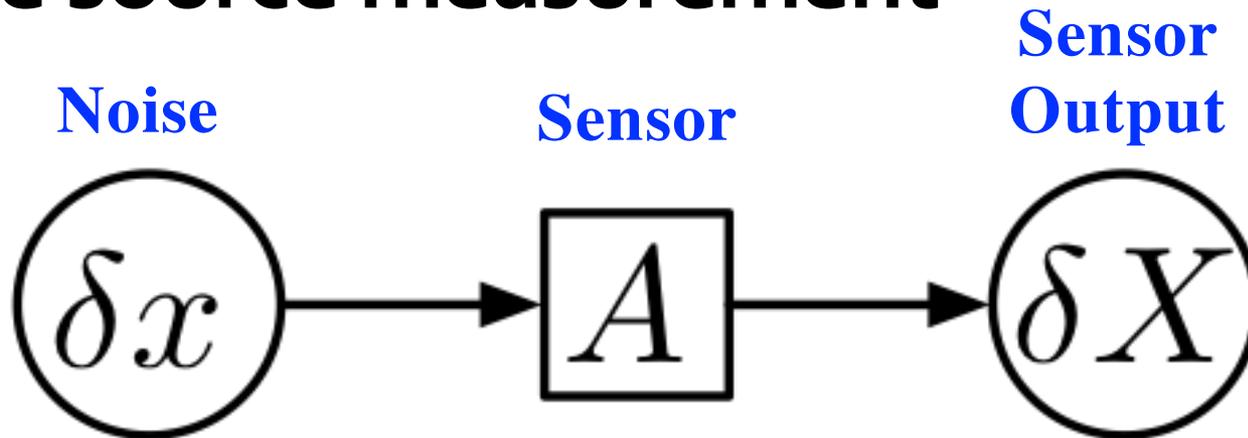
- Noise coupling model



- Make the noise ( $dx$ ) small
- Make the coupling ( $\varepsilon$ ) small
- It is simple... isn't it?

# Noise coupling to the GW channel

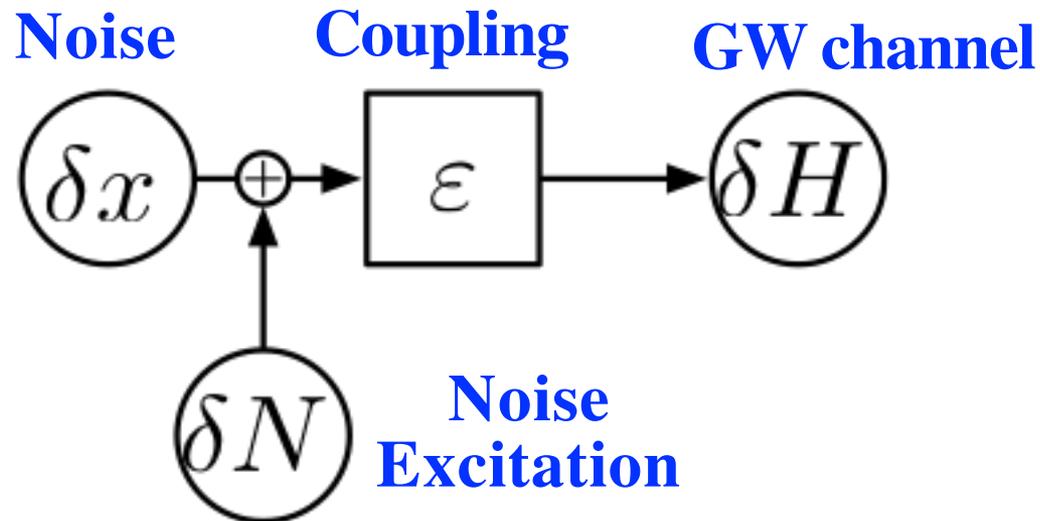
- Noise source measurement



- Use a sensor that is specifically sensitive to  $dx$ .
- This measurement itself is often difficult.

# Noise coupling to the GW channel

## ■ Coupling measurement

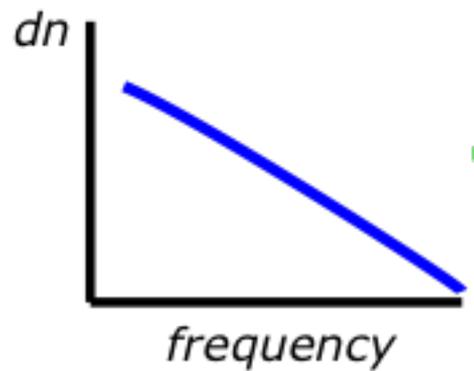


- $dH = \epsilon dN$ .
- The system should be designed to allow the excitation.

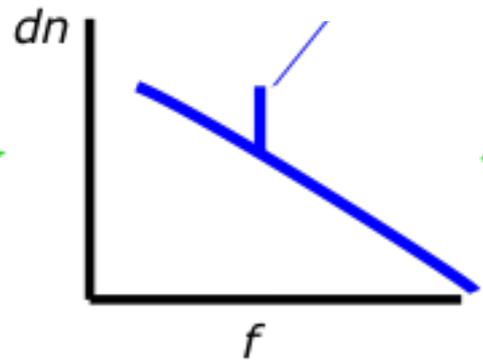
# Noise coupling to the GW channel

## ■ Noise budget

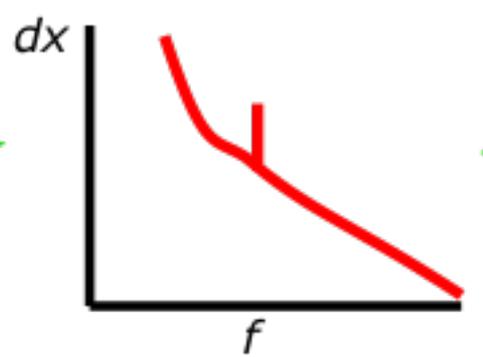
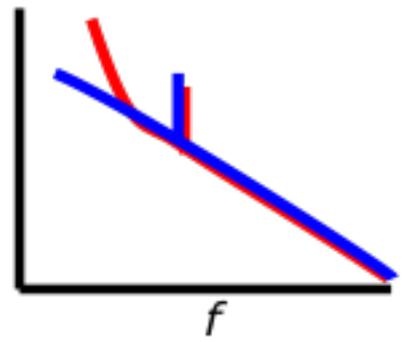
Noise measurement



Artificial Excitation



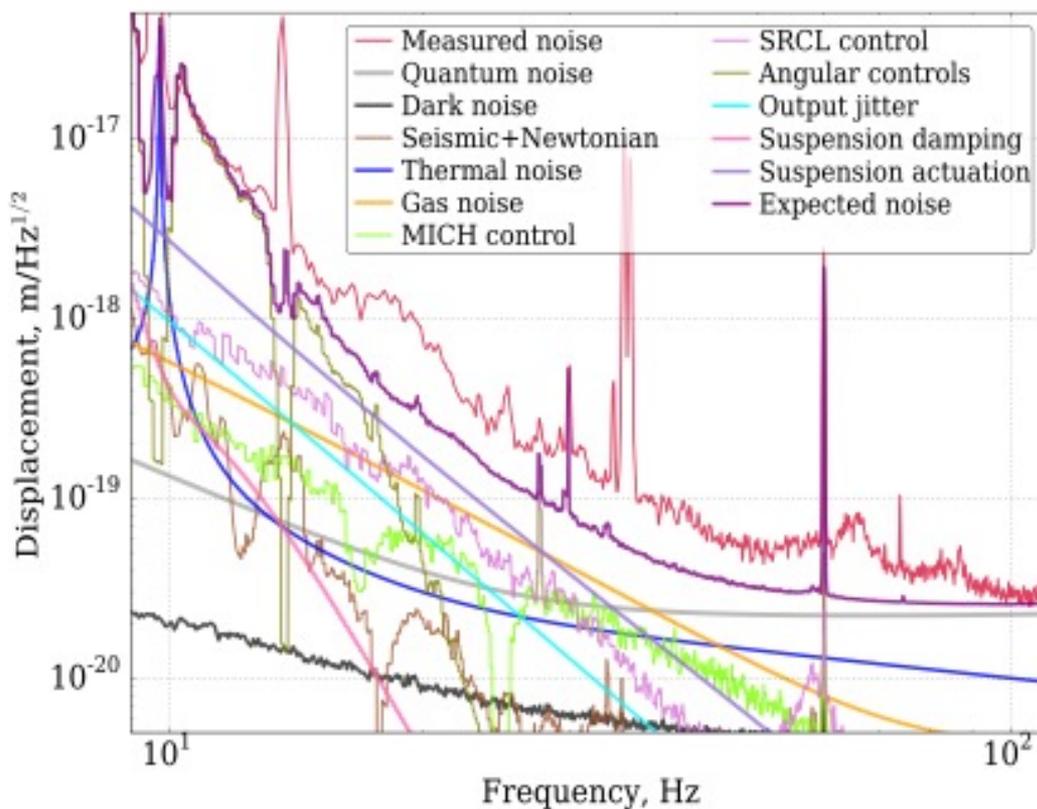
$dx$



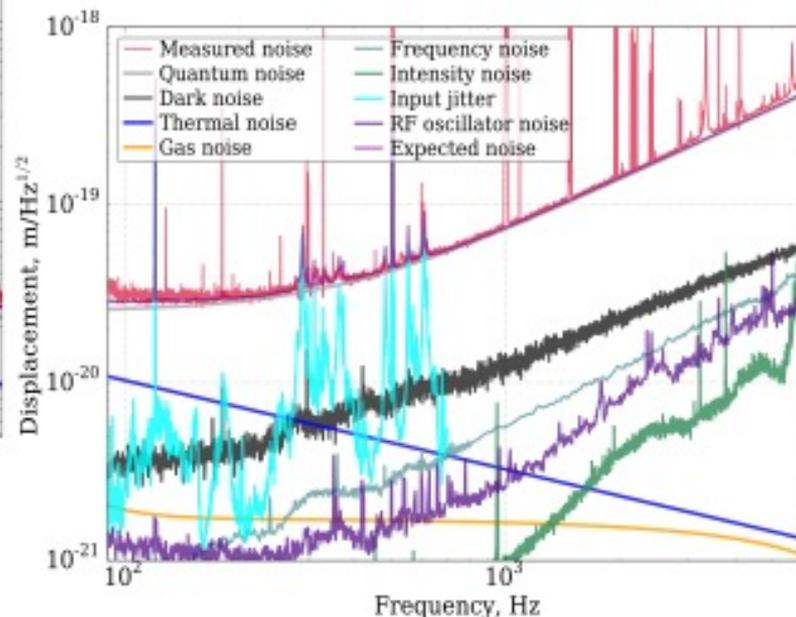
GW channel

# Sensitivity and noise

- Sensitivity (=noise level) of Advanced LIGO
- Noise budget



(a) LIGO Livingston Observatory



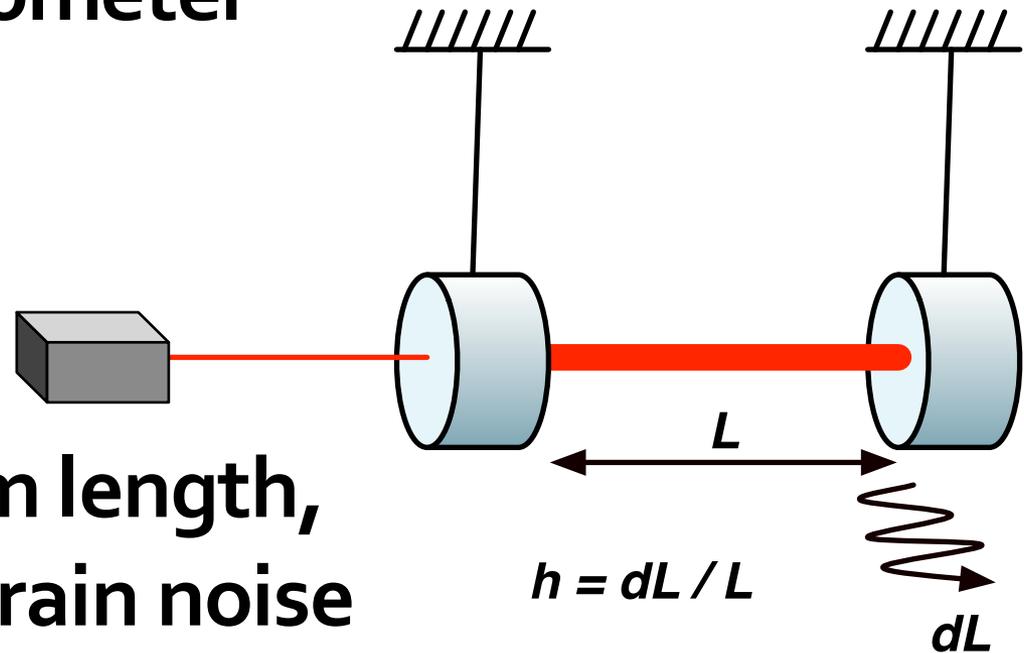
(b) LIGO Hanford Observatory

# Displacement noises

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# Displacement noise

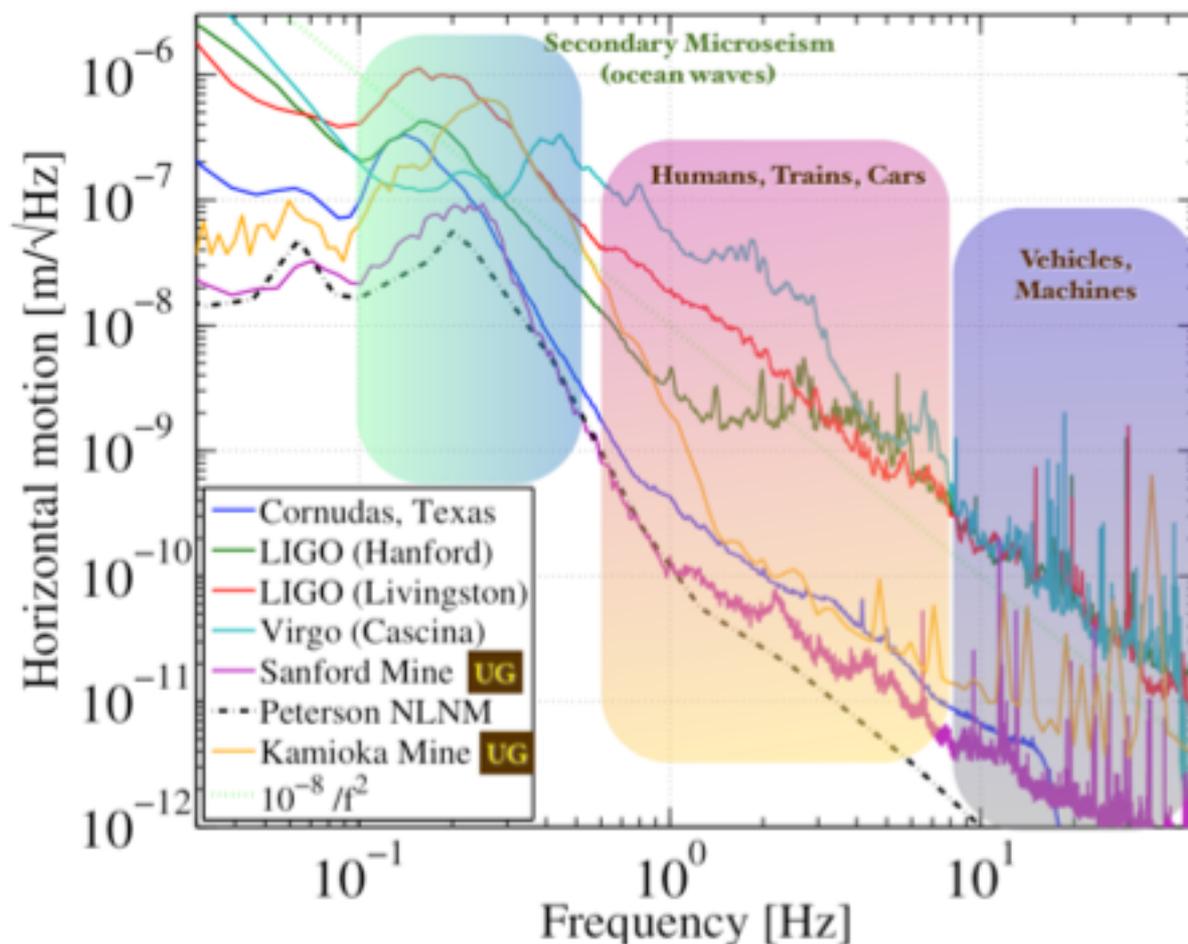
- Mechanical displacement sensed by a laser interferometer



- The longer the arm length, the smaller the strain noise
  - Seismic noise
  - Thermal noise
  - Newtonian Gravity noise

# Displacement noise

- Seismic noise
  - Even when there is no noticeable earth quake...

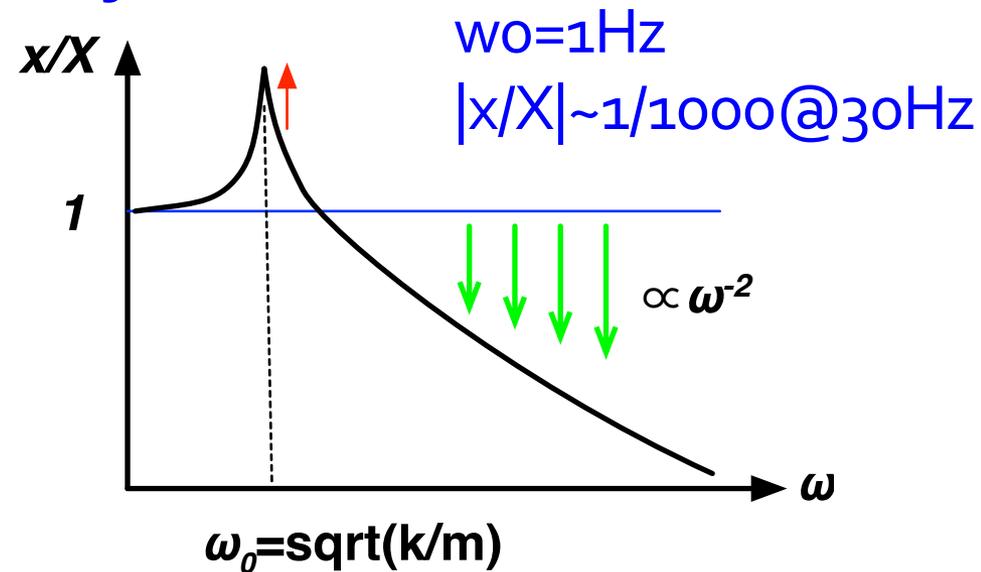
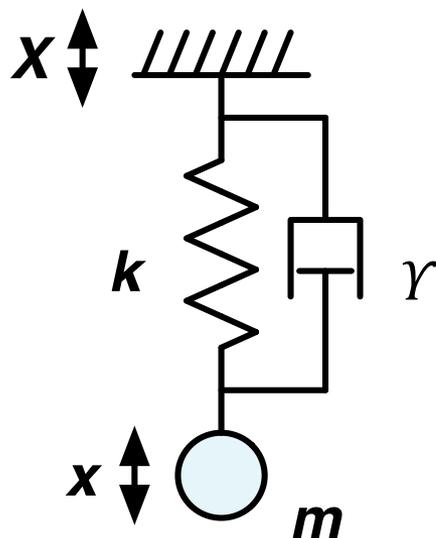


Target  
disp. noise  
 $10^{-20}$  m/rHz



# Displacement noise

- Vibration isolation ~ utilize a harmonic oscillator
  - A harmonic oscillator provides vibration isolation above its resonant frequency



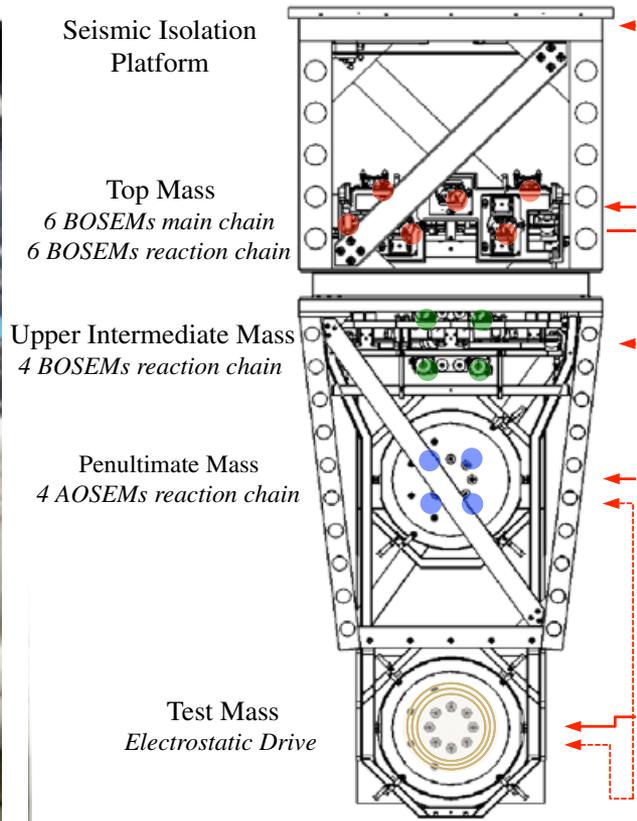
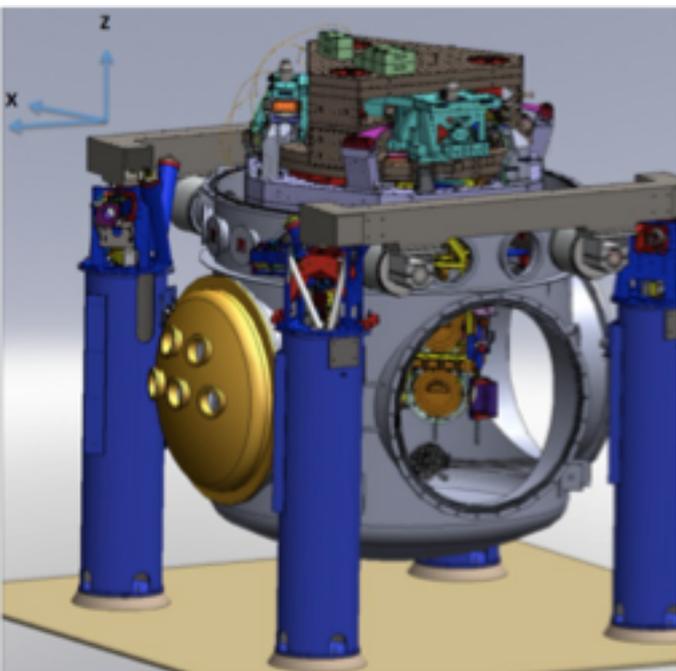
$$m\ddot{x} = -k(x - X) - \gamma(\dot{x} - \dot{X})$$

$$\left(\omega_0^2 + i\frac{\gamma}{m}\omega - \omega^2\right) \tilde{x} = \left(\omega_0^2 + i\frac{\gamma}{m}\omega\right) \tilde{X}$$

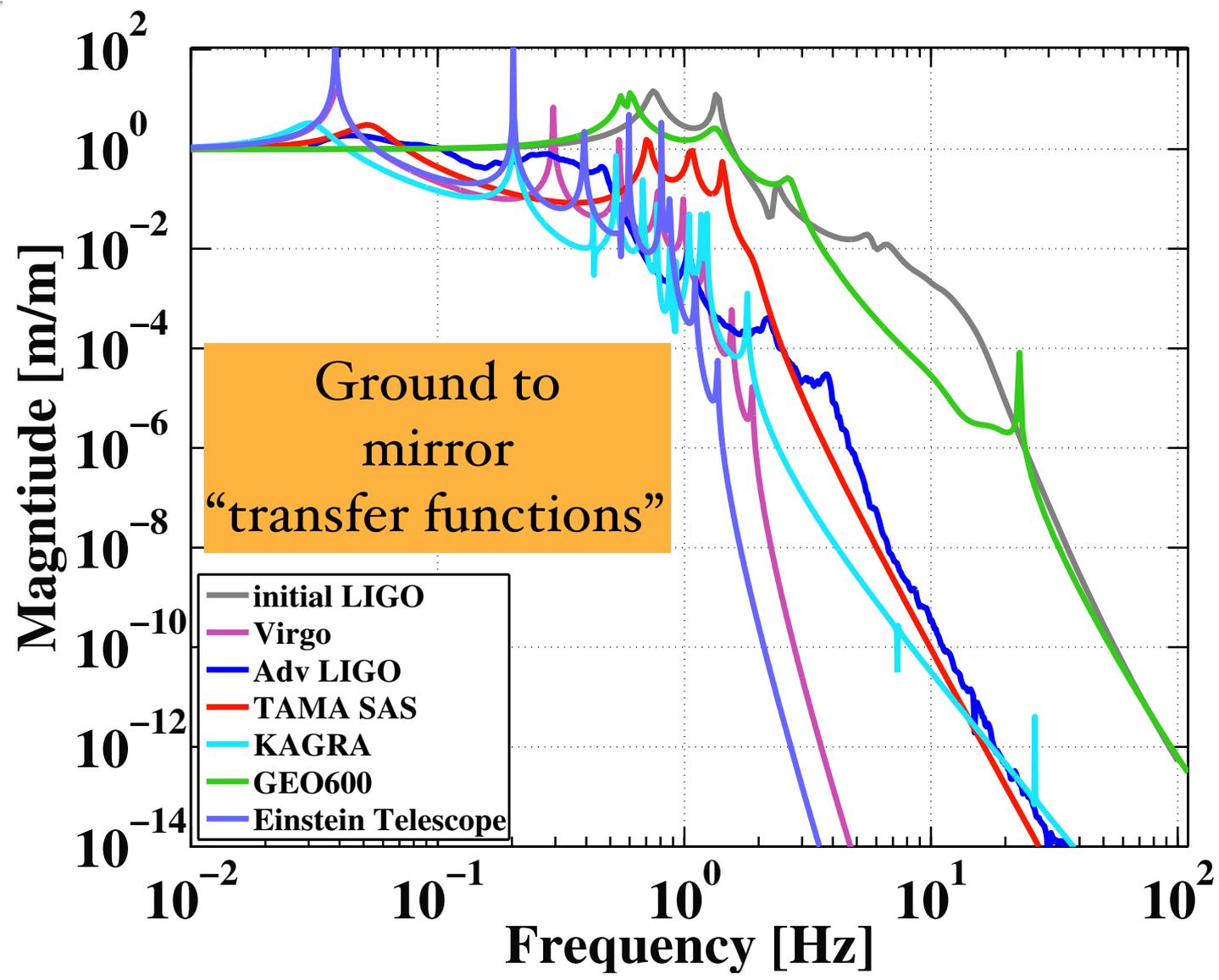
$$\frac{\tilde{x}}{\tilde{X}} = \frac{\omega_0^2 + i\frac{\gamma}{m}\omega}{\omega_0^2 + i\frac{\gamma}{m}\omega - \omega^2}$$

# Displacement noise

- aLIGO vibration isolation
- Hydraulic active isolation / Invacuum Active Isolation Platforms / Multiple Pendulum

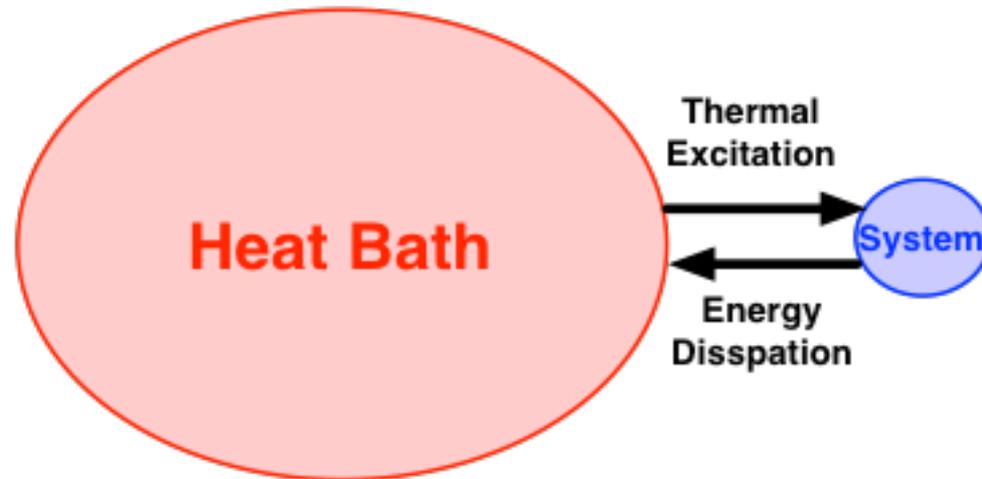


# Displacement noise



# Displacement noise

- Thermal noise:
- System in thermal equilibrium
  - the system can dissipate its energy to the heat bath
  - the system is thermally excited by thermal fluctuation



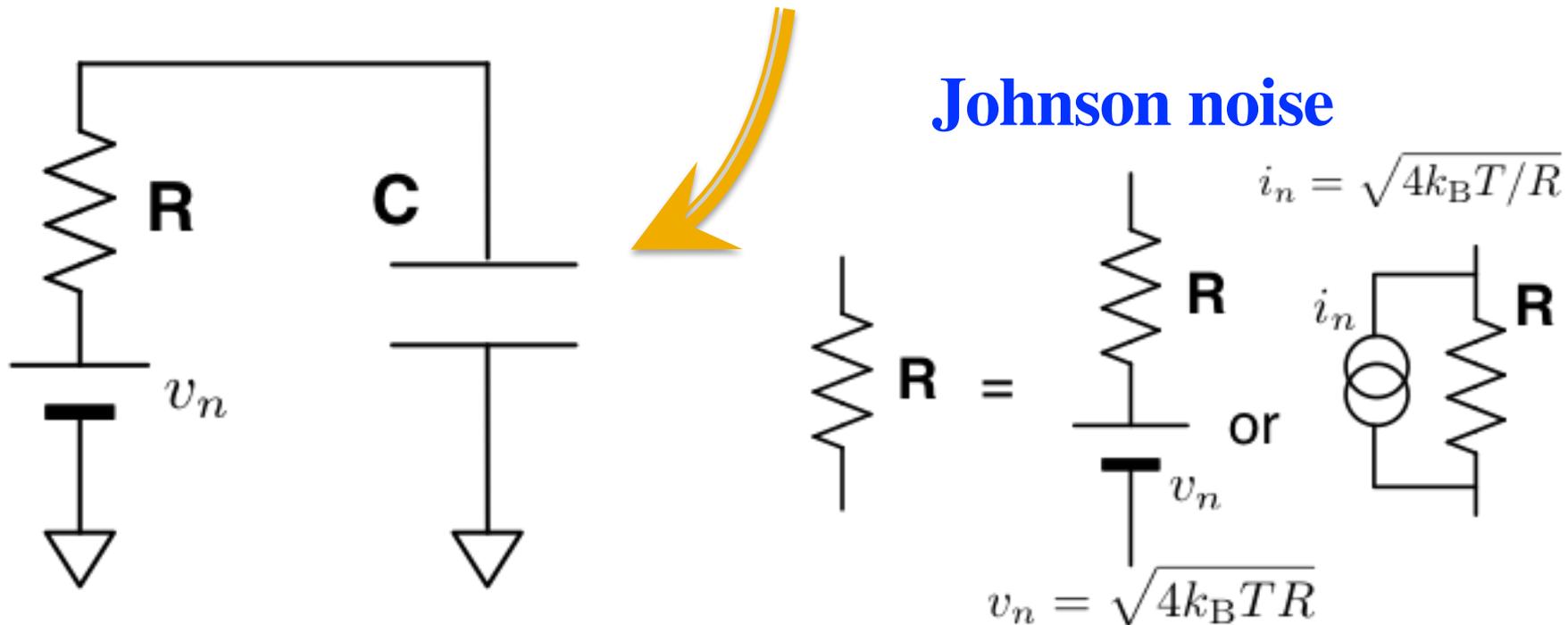
- The level of thermal excitation can be evaluated using the energy dissipation of the system (e.g. friction, resistance)  
"Fluctuation Dissipation Theorem"

# Displacement noise

- Equipartition Theorem

In thermal equilibrium, each d.o.f. has the energy of

$$\langle E \rangle = \frac{1}{2} k_B T$$



# Displacement noise

- Mechanical resonance
- Quality factor  $Q$

$$Q \stackrel{\text{def}}{=} 2\pi \times \frac{\text{energy stored}}{\text{energy dissipated per cycle}}$$

- High  $Q$   
=> most of the  $kBT/2$  energy contained in the resonant freq

=> Lower thermal motion at off-resonance

Low  $Q$



High  $Q$



# Displacement noise

- **Mirror substrate thermal noise**

- **Brownian motion**

Mechanical loss associated with the internal friction

↔ **Thermally excited body modes**

Optical coating (higher mech. loss)  
**will be limiting noise source in aLIGO**

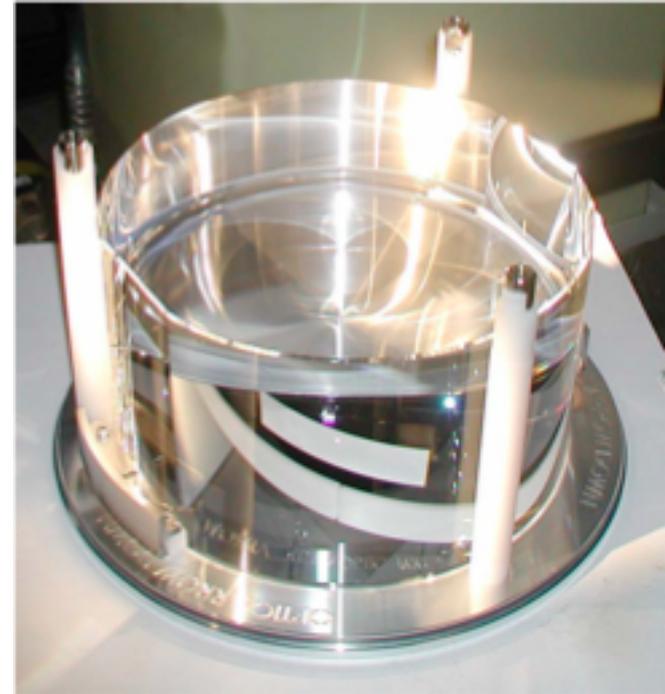
- **Thermo elastic noise**

Elastic strain & thermal expansion coefficient  
=> cause heat distribution & flow in the substrate

↔ **Temperature fluctuation causes mirror displacement**

- **Thermo-refractive noise**

↔ **Temp. fluctuation causes fluctuation of refractive index**



# Displacement noise

- **Suspension thermal noise**

- **Brownian motion**

Mechanical loss of the suspension fiber

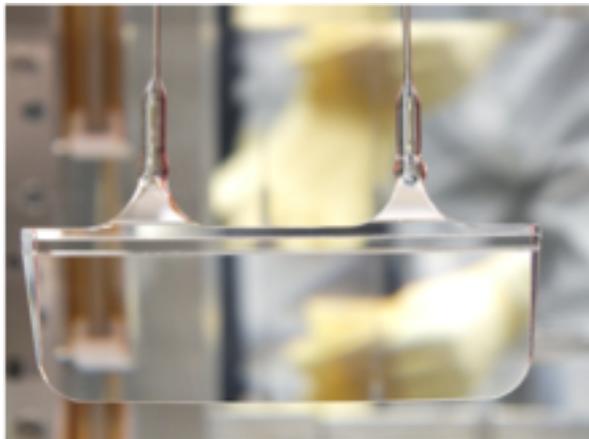
↔ **Thermally excited pendulum modes**

- **Thermo elastic noise**

Elastic strain of the fiber & thermal expansion coefficient

=> cause heat distribution & flow in the fiber

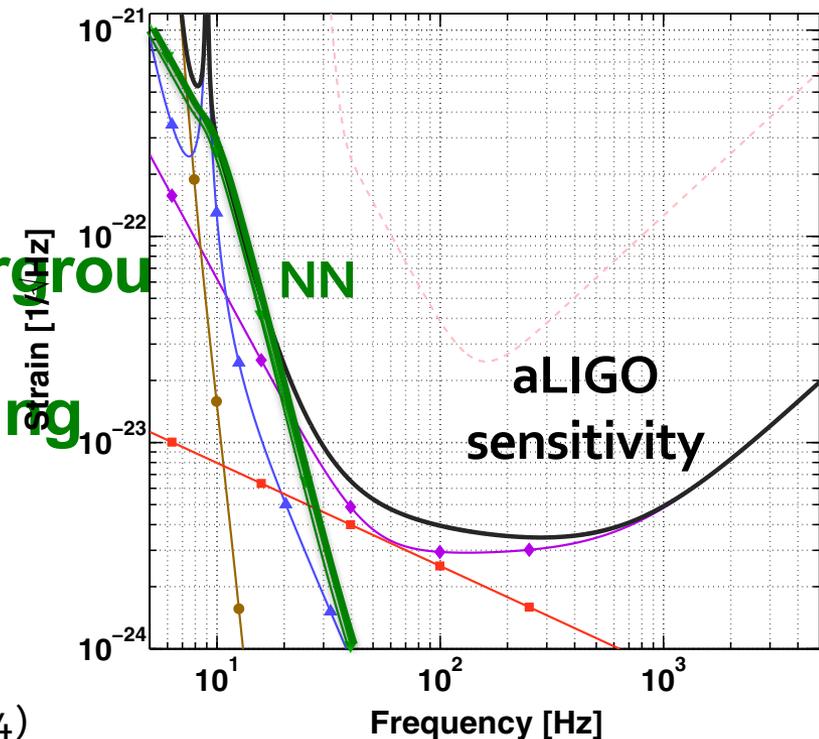
↔ **Temperature fluctuation causes mirror motion**



**<- Monolithic suspension  
for high pendulum Q**

# Displacement noise

- **Newtonian Gravity noise**
  - Mass density fluctuations around the test masses  
=> **test mass motion via gravitational coupling**
  - Dominant source of Newtonian noise  
= **Seismic surface wave**
- **Mitigation**
  - 1) Going to quiet place (underground)
  - 2) Feedforward subtraction
  - 3) Passive reduction by shaping local topography



J Driggers, et al, PRD 86, 102001 (2012)

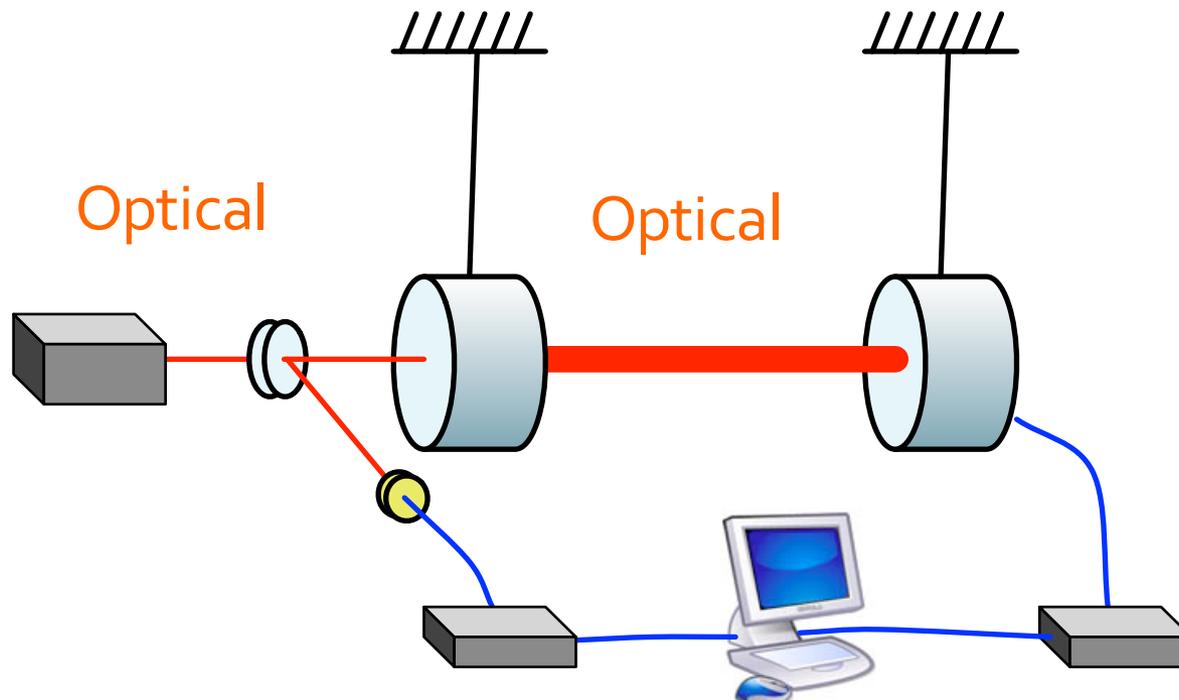
J Harms, et al, Class. Quantum Grav. 31 185011 (2014)

# Optical noises

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# Optical noises

- **Noises that contaminate the readout signal**
  - **Quantum noises (shot noise, radiation pressure noise)**
  - **Laser technical noises (frequency/intensity noise)**
  - **Modulation noises**
  - **Scattered light noise**



# Optical noises

- **Quantum noises: Shot noise**
  - Noise due to photon counting statistics
  - N detected photon => standard deviation  $\sqrt{N}$
  - Increasing the incident power  $P_{in}$ ,
    - => The shot noise is increased by  $\sqrt{P_{in}}$
    - => The signal amplitude is increased by  $P_{in}$
  - In total, the signal-to-noise ratio is improved by

$$\text{SNR} \propto \sqrt{P_{in}}$$

# Optical noises

- **Quantum noises ~ Radiation pressure noise**

- **Photon number fluctuation in the arm cavity**

**=> Fluctuation of the back action force**

- **Quantum noise of the input laser**

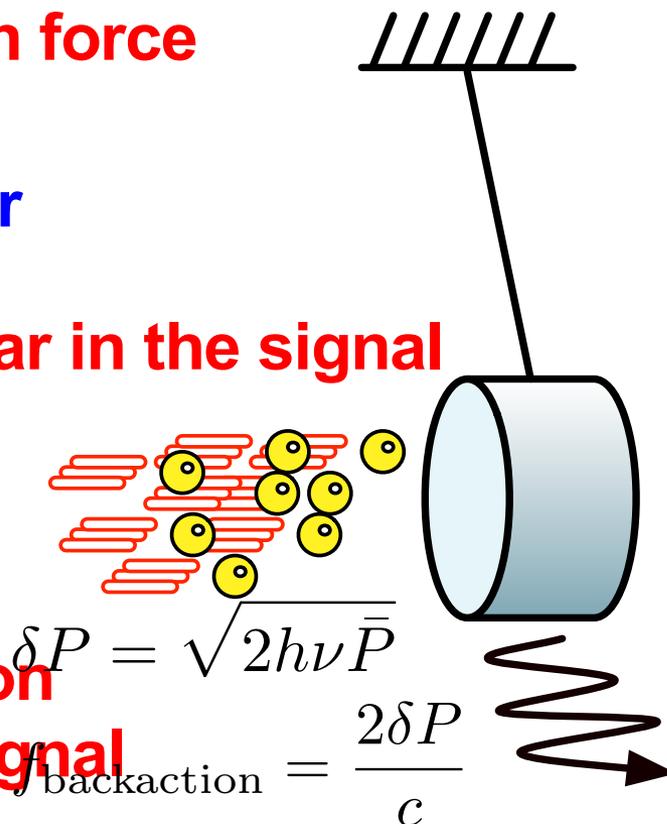
**=> Common noise for two arms**

**=> cancelled and does not appear in the signal**

- **Vacuum fluctuation injected from the dark port**

**=> Differentially power fluctuation**

**=> Cause the noise in the GW signal**



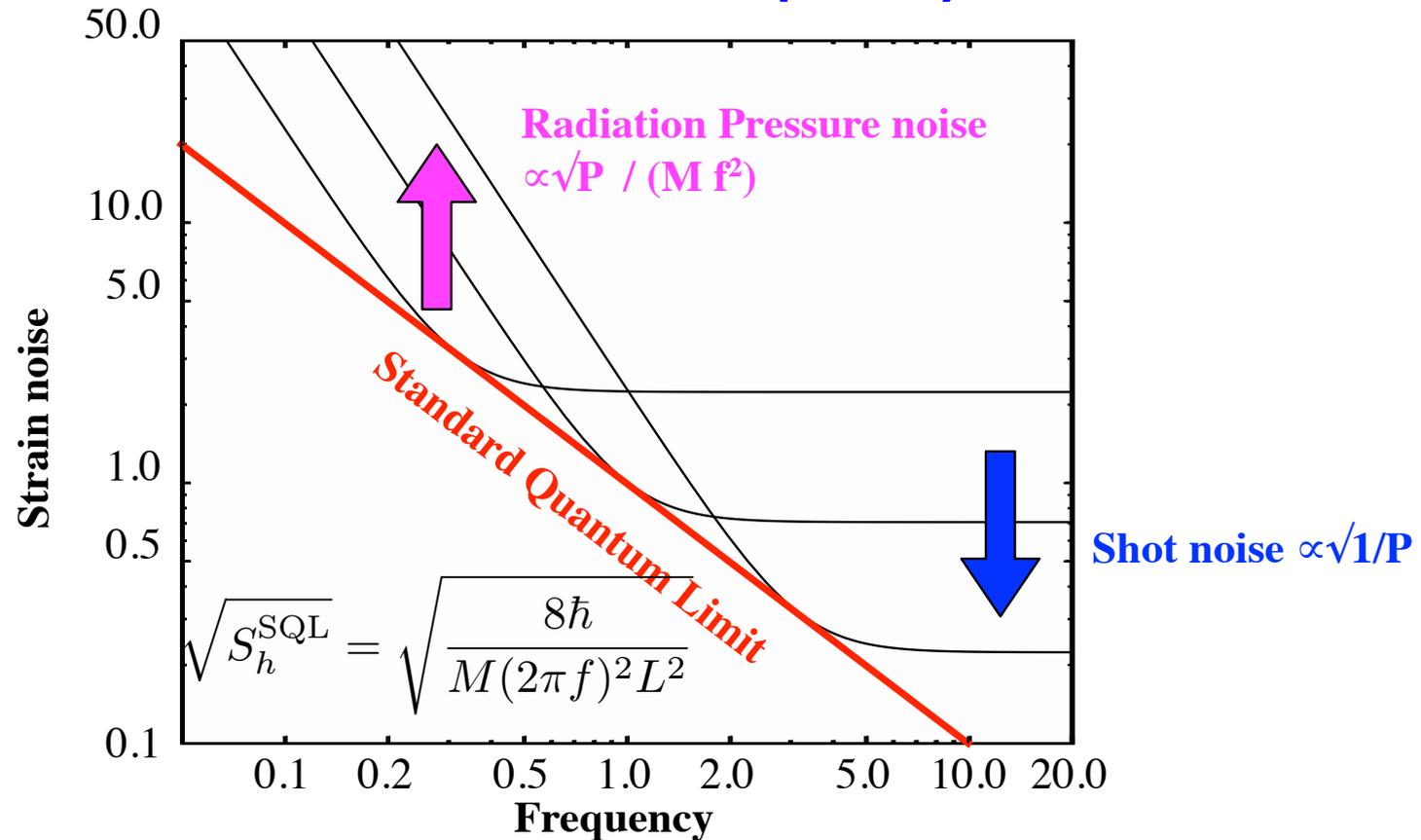
$$\delta P = \sqrt{2h\nu\bar{P}}$$

$$f_{\text{backaction}} = \frac{2\delta P}{c}$$

$$\tilde{x} = \frac{f_{\text{backaction}}}{M\omega^2}$$

# Optical noises

- Quantum noises
  - Standard Quantum Limit (SQL)



- Trade-off Between Shot Noise and Radiation-Pressure Noise
- Uncertainty of the test mass position due to observation

# Optical noises

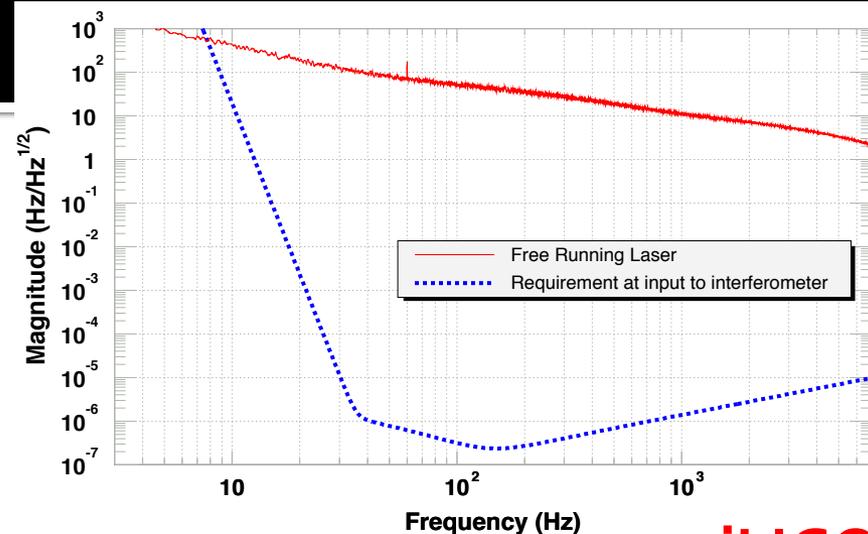
- **Laser frequency noise**
  - **Laser wavelength ( $\lambda = c / \nu$ )**  
= reference for the displacement measurement
  - **Optical phase  $\phi = 2 \pi \nu L / c$**   
 $d\phi = 2 \pi / c (L d\nu + \nu dL) \ll \text{indistinguishable}$

$$\frac{dL}{L} = \frac{d\nu}{\nu}$$

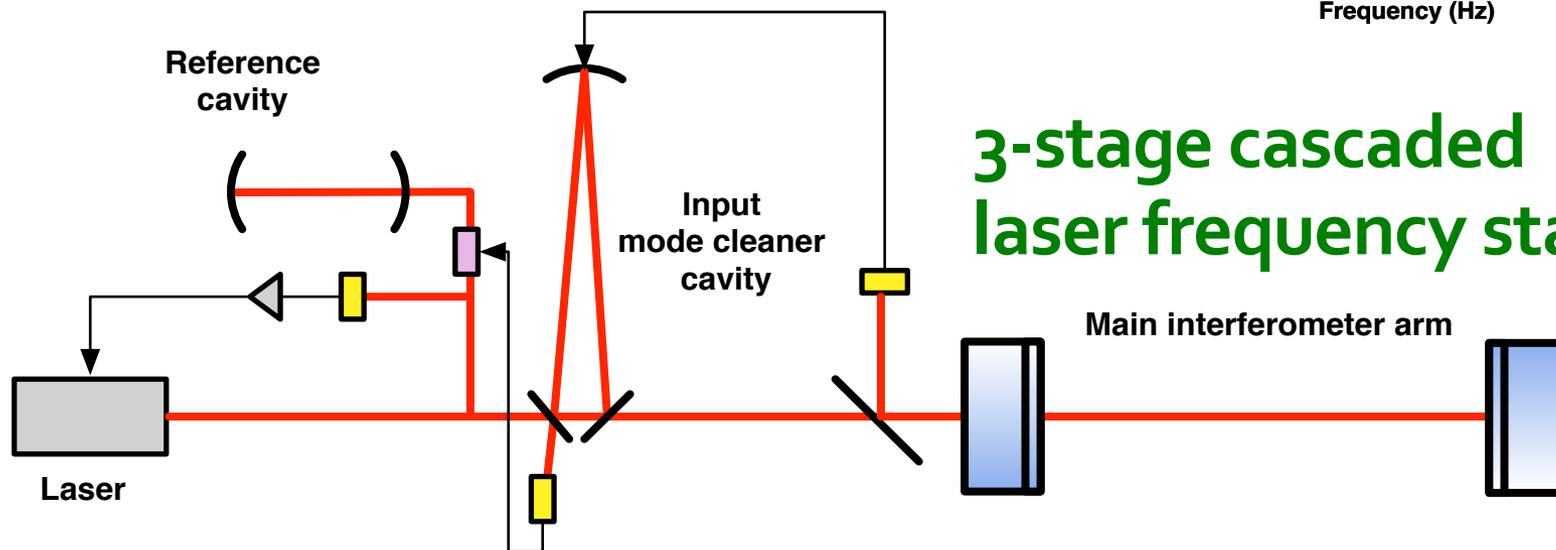
- **$dL/L$  target  $10^{-24}$**   
 =>  **$d\nu = 10^{-24} \times 300 \text{ THz}$  (1064nmYAG laser)**  
        **$= 3 \times 10^{-10} \text{ Hz/rtHz}$**

# Optical noises

- **Laser frequency noise**
  - **Target:  $dv_{\text{eff}} = 3 \times 10^{-10}$  Hz/rHz**
  - **Laser stability**  
 **$dv = 10 \sim 100$  Hz/rHz @100Hz**



iLIGO



Michelson's differential sensitivity provides  
Frequency noise cancellation of  $1/100 \sim 1/1000$   
"Common Mode Rejection"

# Optical noises

- **Laser intensity noise**
  - **Relative Intensity Noise (RIN):  $dP/P$**
  - **Sensor output  $V = P \times x$**   
 $\Rightarrow dV = P dx + x dP$        $\Leftarrow$  **indistinguishable**

$$\frac{dx}{x_{\text{offset}}} = \frac{dP}{P}$$

- **Requirement:  $RIN = 10^{-9} \text{ 1}/\sqrt{\text{Hz}}$**   
 $x_{\text{ofs}} = 10\text{e-}12$  (DC Readout)  
 $\Rightarrow dx = 1\text{e-}20 \text{ m}/\sqrt{\text{Hz}}$

# Optical noises

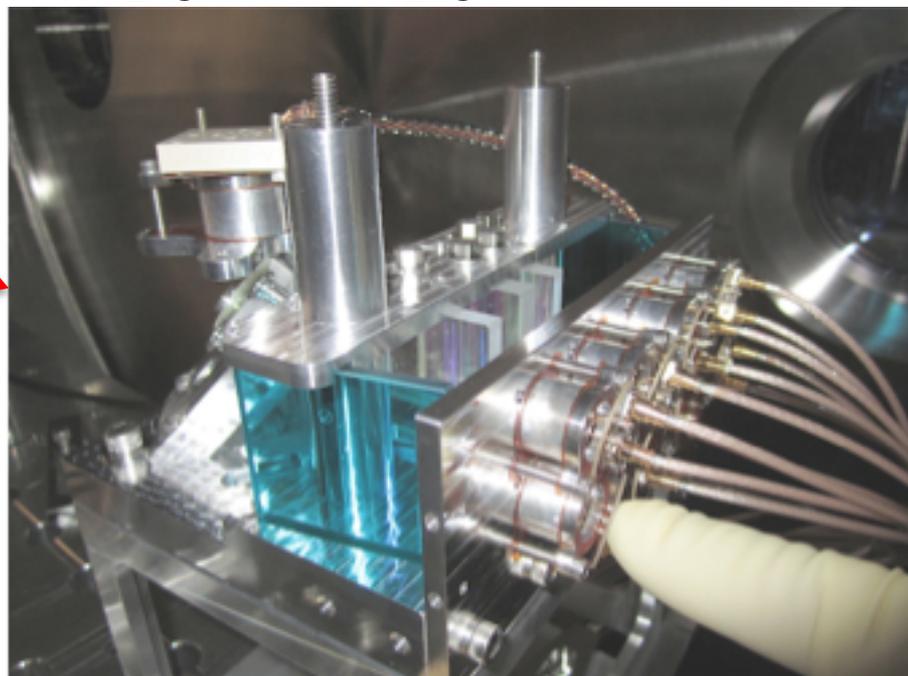
- **Laser intensity noise ~ intensity stabilization**

- Requirement:  $RIN = 10^{-9} 1/\sqrt{\text{Hz}}$
- 2-stage cascaded intensity stabilization control
- Challenge: requires 300mA of photodetection

**Shot noise limited RIN**

$$\frac{i_{\text{shot}}}{i_{\text{DC}}} = \frac{\sqrt{2ei_{\text{DC}}}}{i_{\text{DC}}} = \sqrt{2e/i_{\text{DC}}}$$

- In-vacuum 8-branch Photodiode array



P. Kwee et al,  
Optics Express 20 10617-10634 (2012)

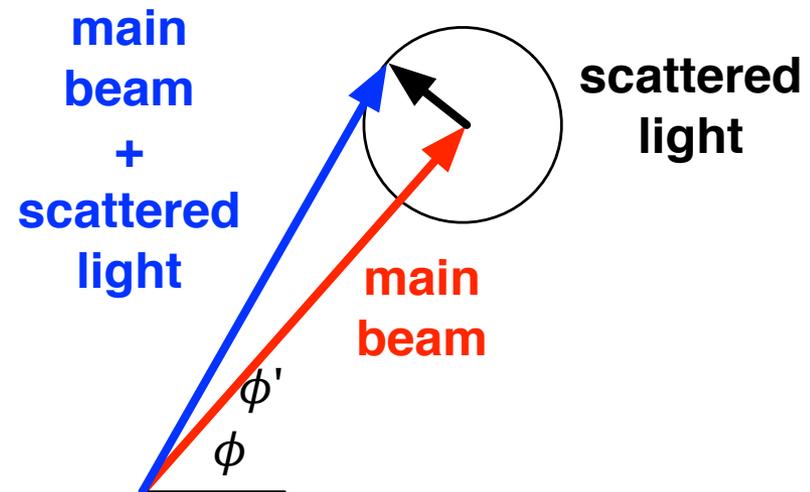
# Optical noises

## ■ Scattered light noise

- Scattered light recouples to the interferometer beam with an arbitrary phase  
=> causes amplitude and phase fluctuation
- Two effects:
  1. **Small motion regime:** linear coupling of the phase fluctuation
  2. **Large motion regime:** low freq large motion of the scattering object  
=> upconversion via fringe wrapping

## ■ Mitigation

- Reduce scattered light
- Vibration isolation of the scattering object



# Electrical noises

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# General rules for electrical noises

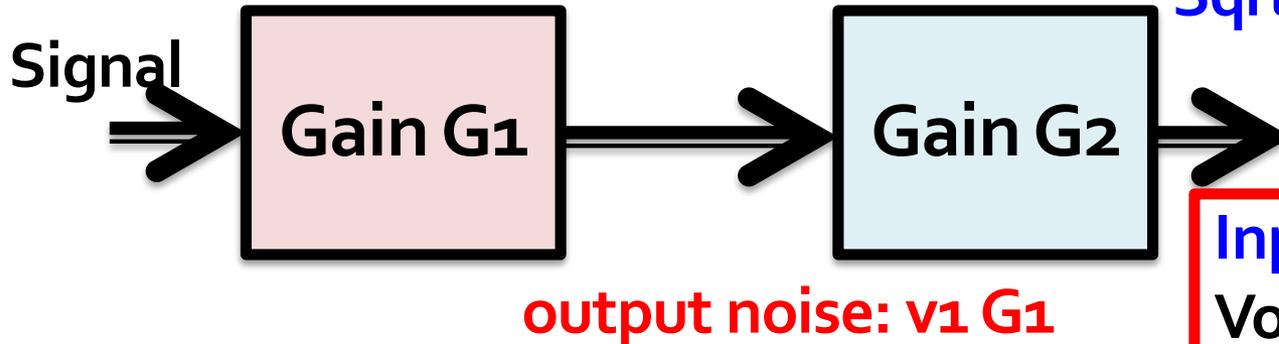
- Low noise amplification at the beginning
- Give necessary gain as early as possible
- Don't attenuate (and amplify again)

input noise:  $v_1$

input noise:  $v_2$

output noise  $V_{out}$ :

$$\text{Sqrt}[(v_1 G_1 G_2)^2 + (v_2 G_2)^2]$$



Input equivalent noise

$$V_{out} / (G_1 G_2)$$

$$= \text{Sqrt}[v_1^2 + (v_2/G_1)^2]$$

## ■ Lessons

- The input referred noise is determined by  $v_1$
- It won't become better by the later stages
- If  $G_1$  is big enough, we can ignore the noise of later stages

# Digitization (Quantization) noise

- Analog signals ( $\sim \pm 10\text{V}$ )  $\rightarrow$  Digital signal
  - Digitized to a discrete N bit integer number

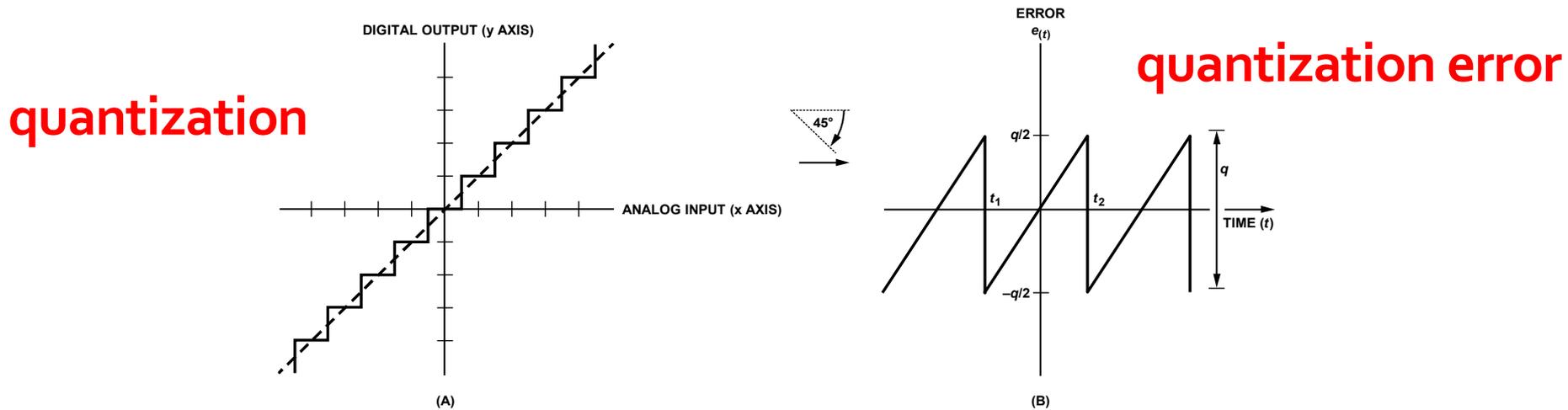


Figure 1. Ideal ADC Transfer Function (A) and Ideal N-Bit ADC Quantized Noise (B)

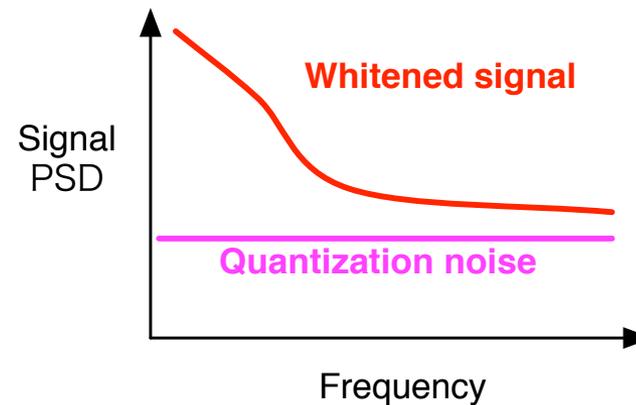
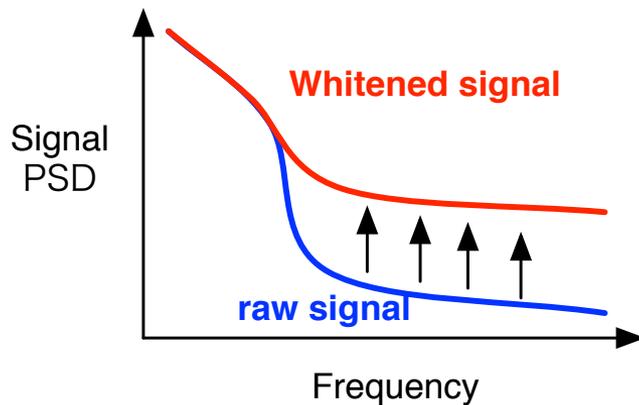
<http://www.analog.com/static/imported-files/tutorials/MT-229.pdf>

- Quantization causes a white noise  $V_n = \frac{\Delta}{\sqrt{12}} \text{ [V}/\sqrt{\text{Hz}}]$   
 e.g.  $\pm 10\text{V}$  16bit  $\Rightarrow \Delta = 0.3\text{mV} \Rightarrow V_n \sim 100 \mu\text{V}/\sqrt{\text{Hz}}$   
 cf. Input noise of a typical analog circuit  $10\text{nV}/\sqrt{\text{Hz}}$

# Digitization (Quantization) noise

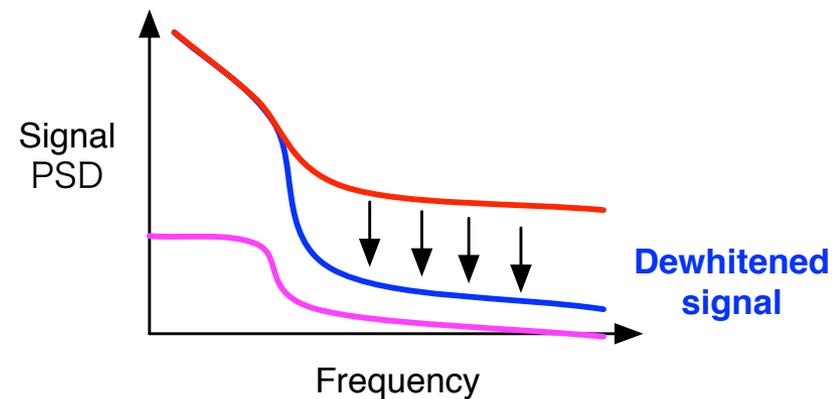
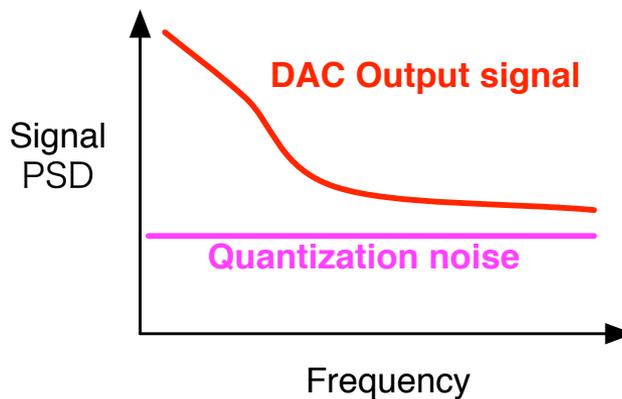
## Whitening

- Amplify a signal in the freq band where the signal is weak



## Dewhitening

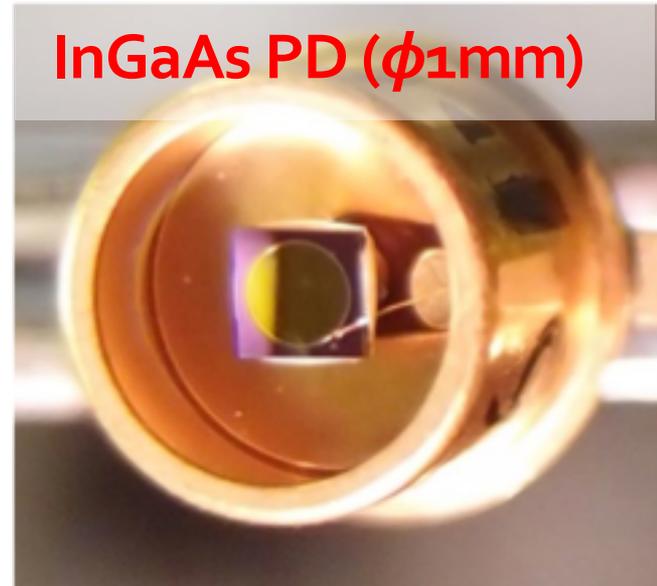
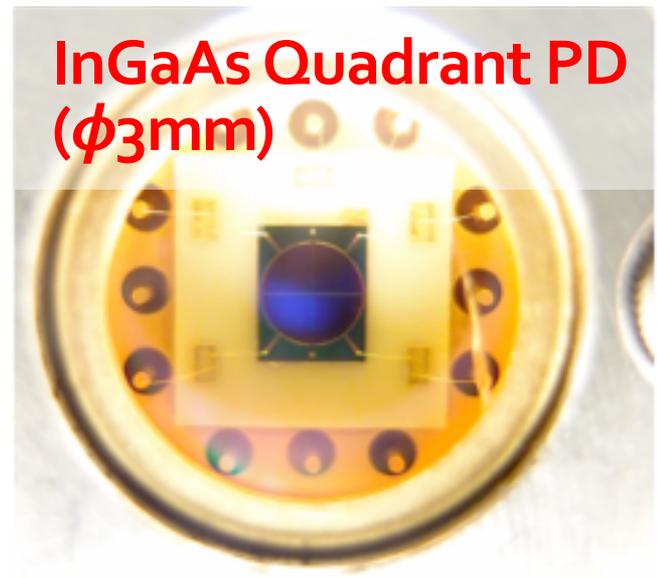
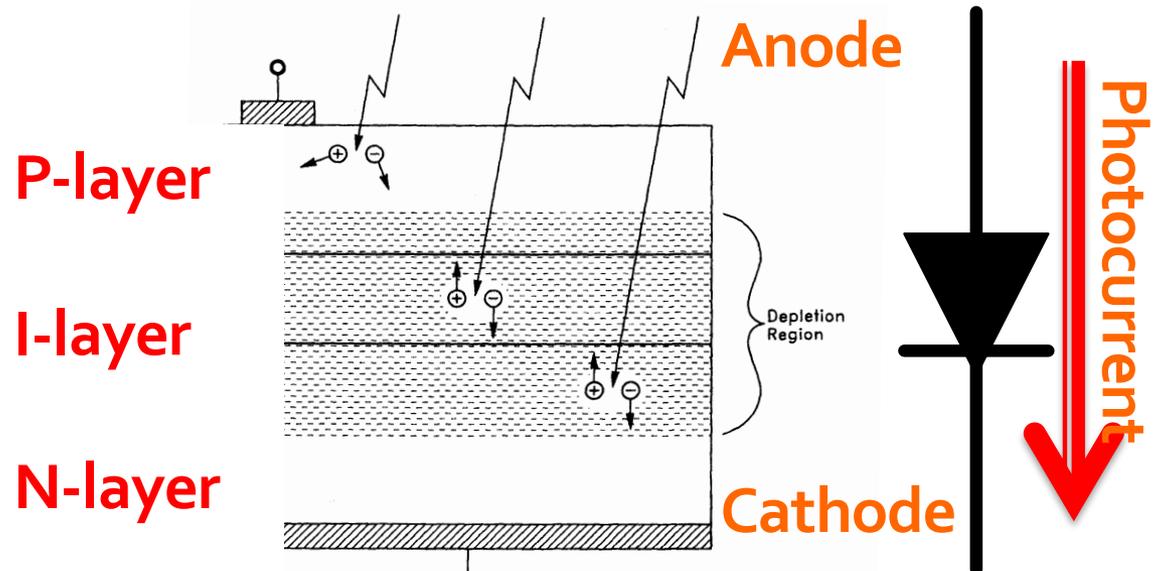
- Amplify a signal in the freq band where the signal is weak



# Noise in photodetectors

## ■ Photodiodes

- PIN photodiodes  
(InGaAs for near IR, Si for visible)
  - Good linearity
  - Low noise
  - High Quantum Efficiency (>90%)



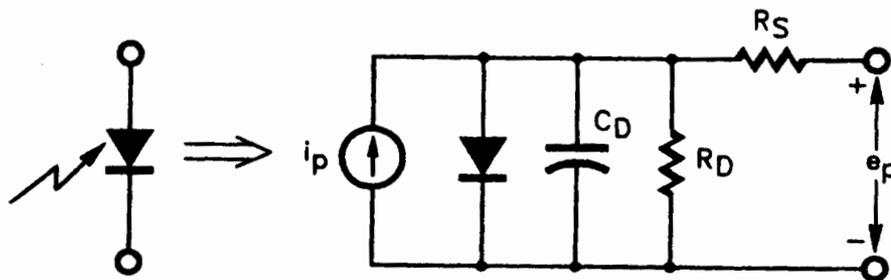
"Photodiode Amplifiers", J. Graeme (McGrawHill 1995)

# Noise in photodetectors

## ■ Noise in photodiodes

### ■ Photodiode equivalent circuit

- **Shunt Capacitance  $R_D$  ( $\sim 100\text{M}\Omega$ )** Usually not a problem
- **Junction Capacitance  $C_D$  (1pF $\sim$ 1nF)**
- **Series Resistance  $R_S$  (1 $\Omega$  $\sim$ 100 $\Omega$ )**



**Figure 1.3** The circuit model of a photodiode consists of a signal current, an ideal diode, a junction capacitance, and parasitic series and shunt resistances.

input referred noise current

$$i_{R_S} \sim \omega C_d \sqrt{4k_B T R_S}$$

The diode aperture size needs to be  $\sim$ mm  $\Rightarrow$   $C_d$  tends to be big.

2mm InGaAs PD:  $R_S \sim 10\Omega$ ,  $C_d \sim 100\text{pF}$

$\Rightarrow i_{R_S} = 20 \text{ pA}/\sqrt{\text{Hz}}$  @100MHz

(equivalent to the shot noise of 1mA light  $\sim 1.3\text{mW}$  @1064nm)

# Summary

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# Summary

- **Summary**
  - **There are such large number of noises**
  - **They are quite omnidisciplinary**
  - **Only one untamed noise can ruin GW detection**
  
  - **GW detection will be achieved by**
    - **Careful design / knowledge / experience**
    - **Logical, but inspirational trouble shooting**
  
  - **Noise “hunting”**
    - **Systematic approach: noise budget**