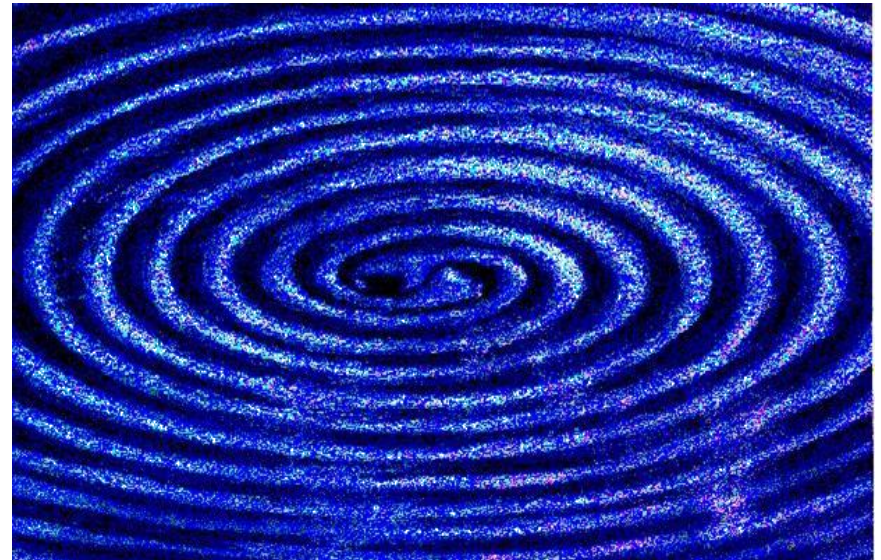


***New Low-Frequency GW Detector
with Superconducting Instrumentation***



Ho Jung Paik
Department of Physics, University of Maryland
GWADW, May 26, 2015

(delivered by Krishna Venkateswara, U Wash)

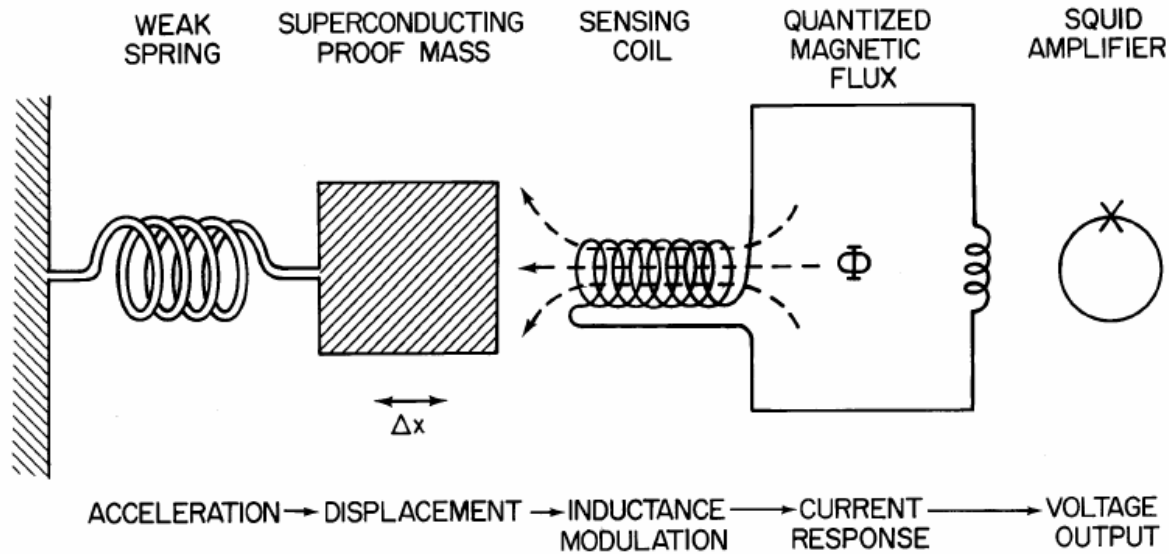
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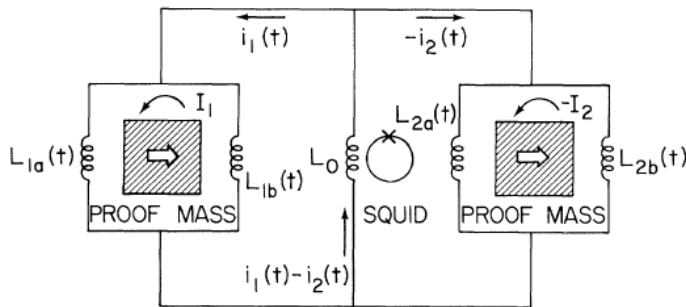
- ❑ **Introduction to Superconducting Gravity Gradiometers (SGG)**
- ❑ **SGG as a Gravitational-Wave Detector**
- ❑ **Current and Future SGGs and Sensitivity**
- ❑ **SOGRO: Design and Important Noise Sources**
- ❑ **Sensitivity Goal**

Introduction

Introduction to Superconducting Gravity Gradiometer (SGG)



SQUID output is proportional to the difference in displacement of the two masses



Gravity gradient tensor

$$\Gamma_{ij} = - \frac{\partial^2 \phi}{\partial x_i \partial x_j}$$

$i=j$: inline-component
 $i \neq j$: cross-component

FIG. 1. A schematic of the superconducting gravity gradiometer.

Test of Newton's Inverse-square law using SGG (1993)

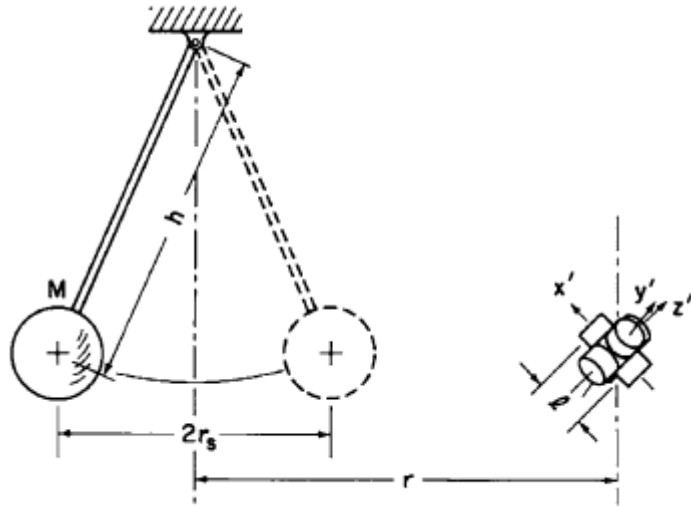
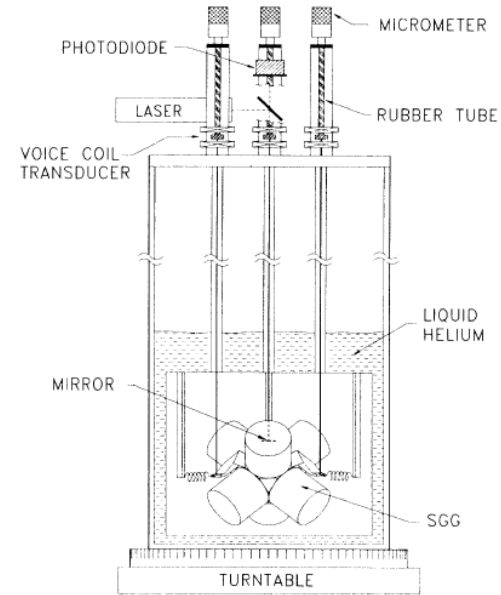
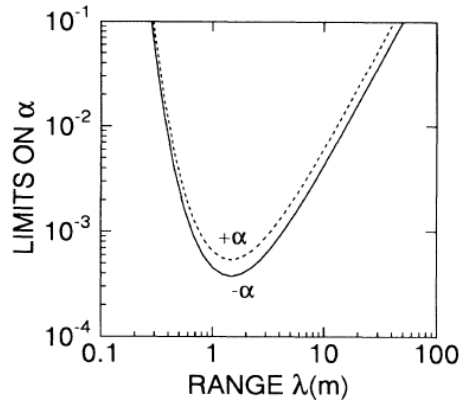


Illustration of Test



Cryostat Schematic



Poisson's equation
$$\sum_i \Gamma_{ii} = \nabla^2 \phi = 4\pi G \rho.$$

In free space,
$$\Gamma_{xx} + \Gamma_{yy} + \Gamma_{zz} = 0$$

2σ limits on new Yukawa interaction

Gravitational Wave Detection

SGG as a Gravitational Wave Detector?

- In the Newtonian limit, Riemann tensor becomes gravity gradient.
 ⇒ Gravity gradiometer measures Riemann tensor components.

$$R_{0i0j} \rightarrow \Gamma_{ij} = -\partial^2 \phi / \partial x_i \partial x_j$$

PHYSICAL REVIEW D

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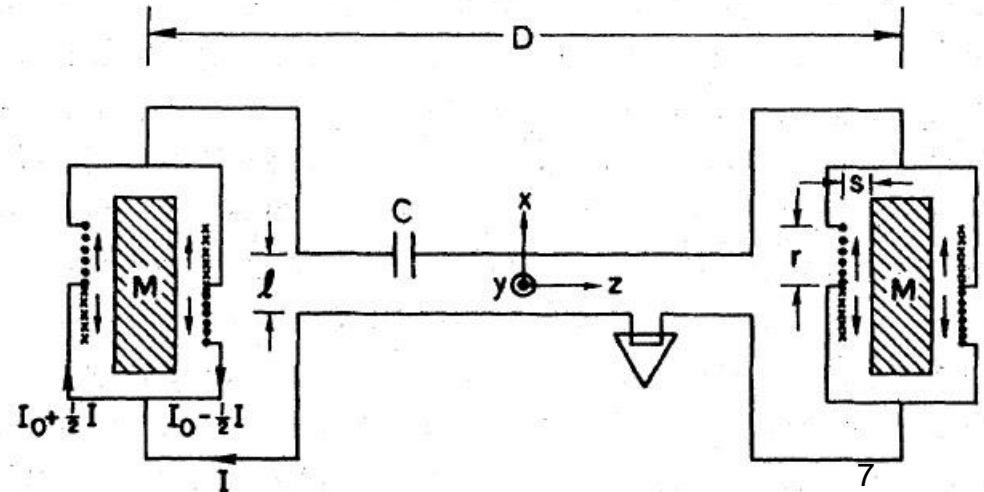
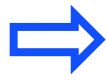
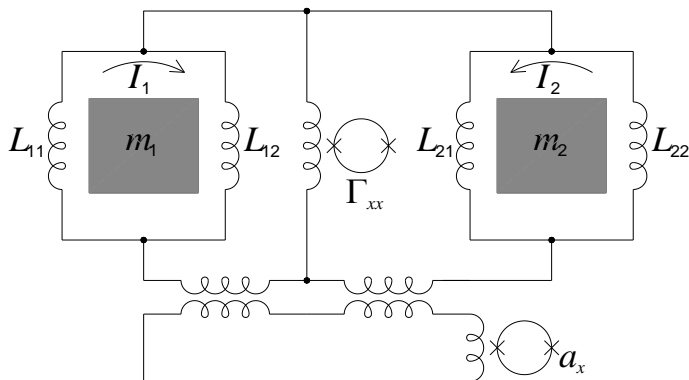
Tunable "free-mass" gravitational-wave detector

Robert V. Wagoner, Clifford M. Will, and Ho Jung Paik*

Institute of Theoretical Physics and Department of Physics, Stanford University, Stanford, California 94305

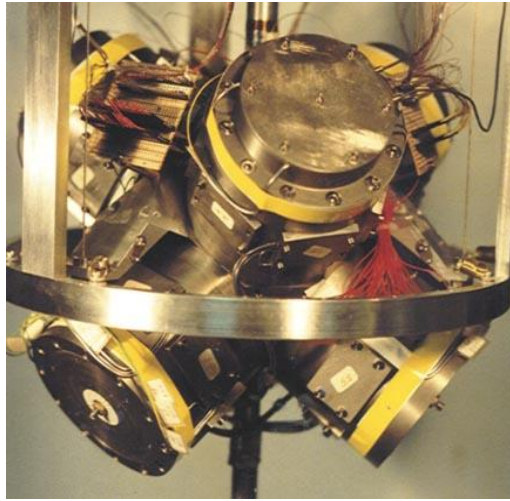
(Received 10 July 1978)

We propose a new type of detector for gravitational radiation. It consists essentially of two masses whose relative motion produces the driving emf of a resonant L - C circuit. The relative momentum of the masses induced by a gravitational wave is determined by the current in the circuit. A unique feature of this system is its ability to be tuned over a wide frequency range. If a quality factor $Q \sim 10^8$ can be achieved in the circuit, a laboratory-size detector cooled to 0.05 K in the absence of other noise could detect a continuous wave metric perturbation $h \gtrsim 3 \times 10^{-26}$ at the frequency of the Crab pulsar after integration for 100 days.

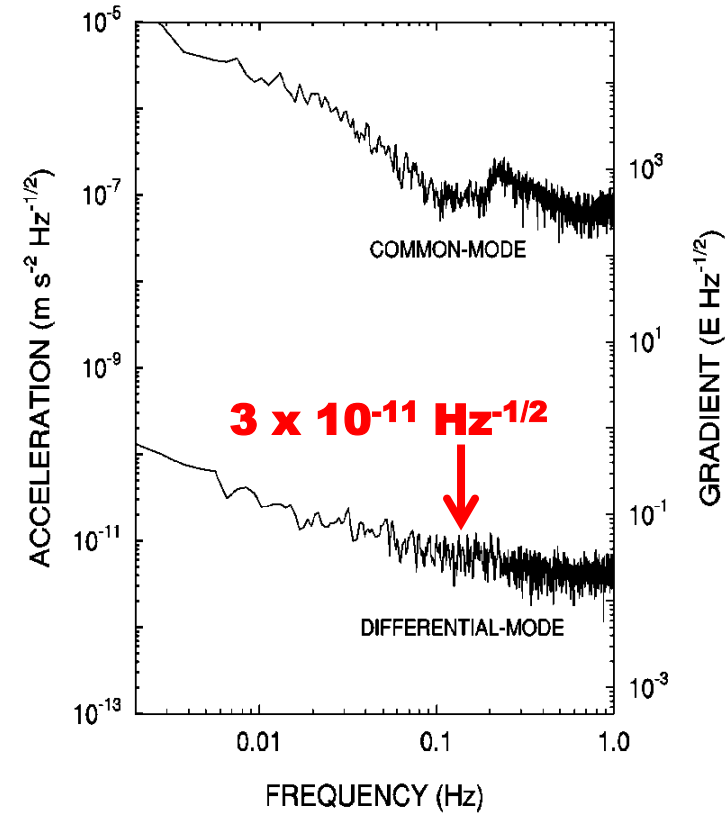
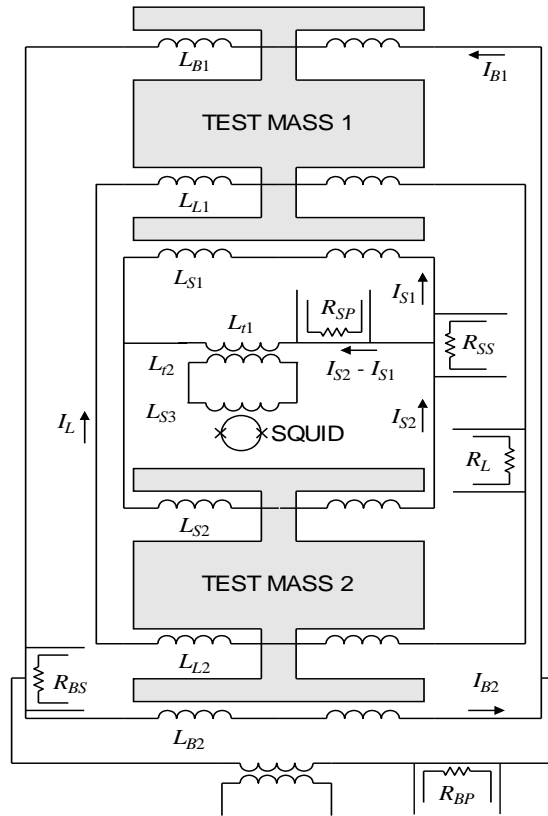


Current and Future SGGs

UM diagonal-component SGG (1993)

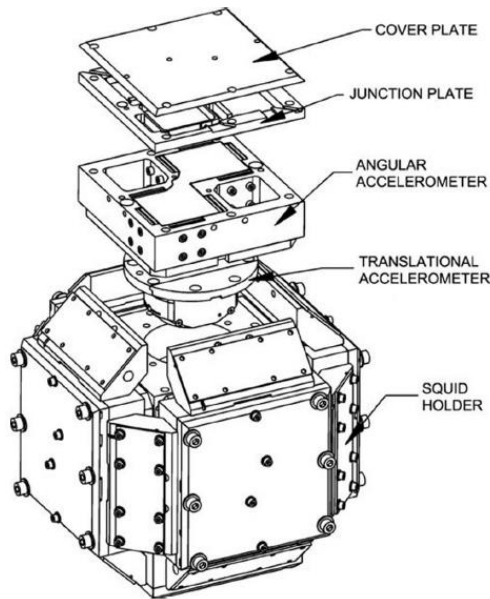


(Moody *et al.*, 2002)

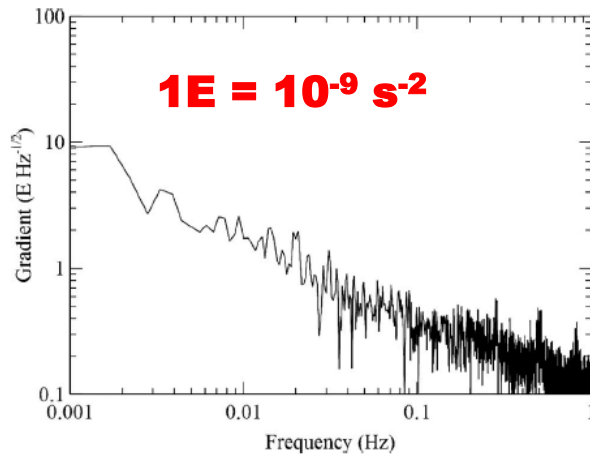


- Test masses are **mechanically suspended** ($f_{DM} \sim 10 \text{ Hz}$).
- Development was completed by early 1990s.
- 100 times better amplitude sensitivity than TOBA 20 years earlier.
 \Rightarrow **10^4 times better limit than TOBA in GW energy density.**

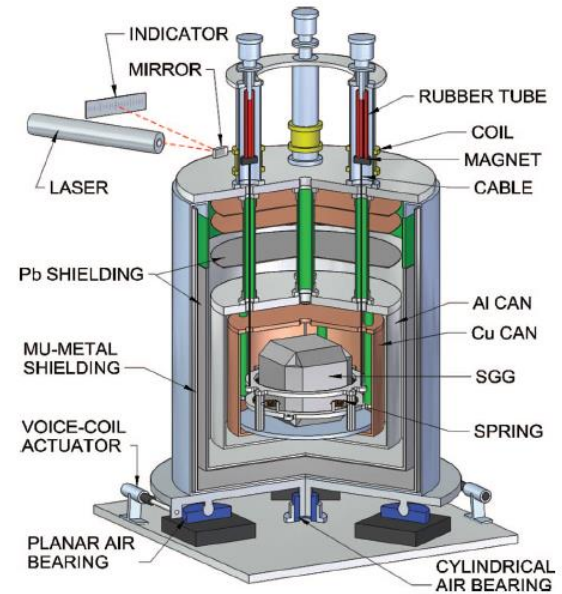
Cross-component SGG (2011)



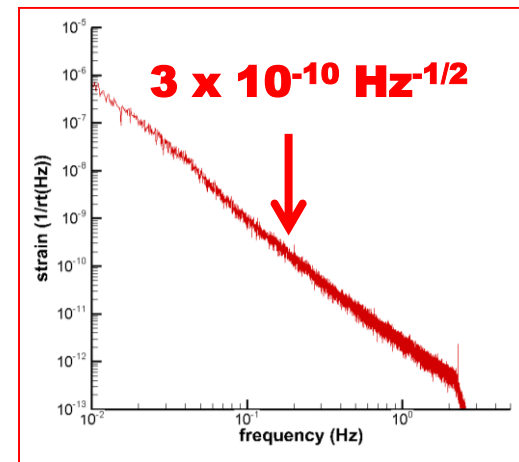
3-axis Cross-component SGG



Gravity gradient sensitivity



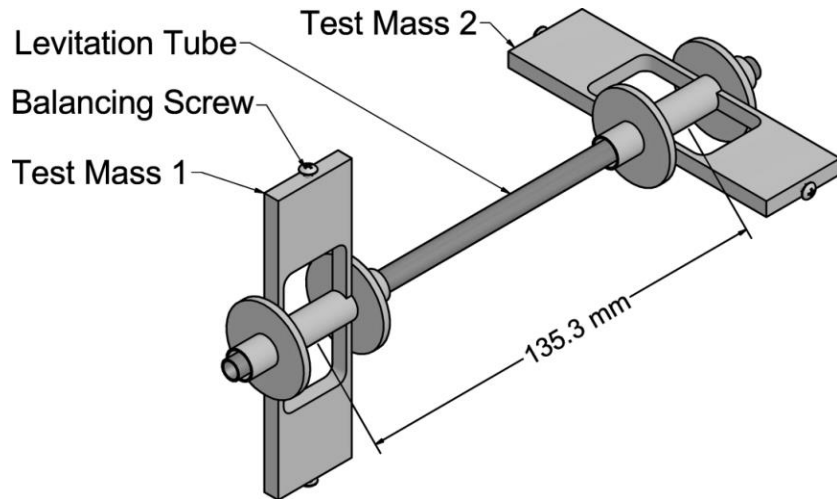
Cryostat Schematic



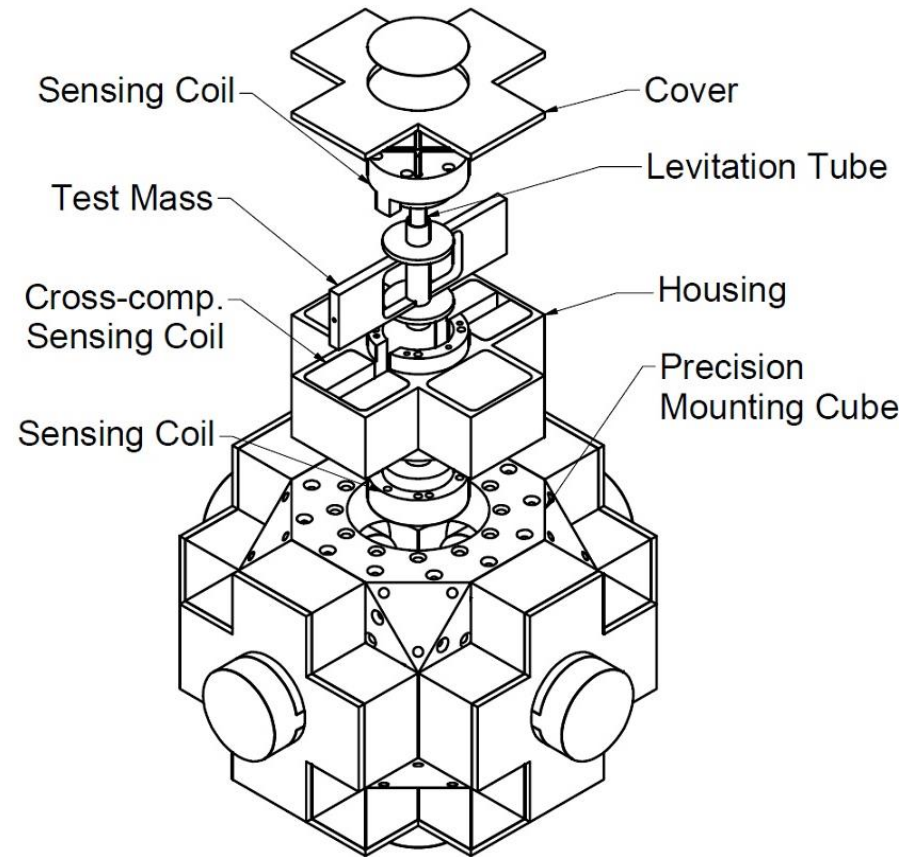
GW strain sensitivity

New Superconducting tensor gravity gradiometer

- More sensitive SGG is under development with NASA support.
- Test masses are **magnetically suspend** ($f_{DM} \sim 0.01$ Hz).
⇒ **Very high sensitivity**



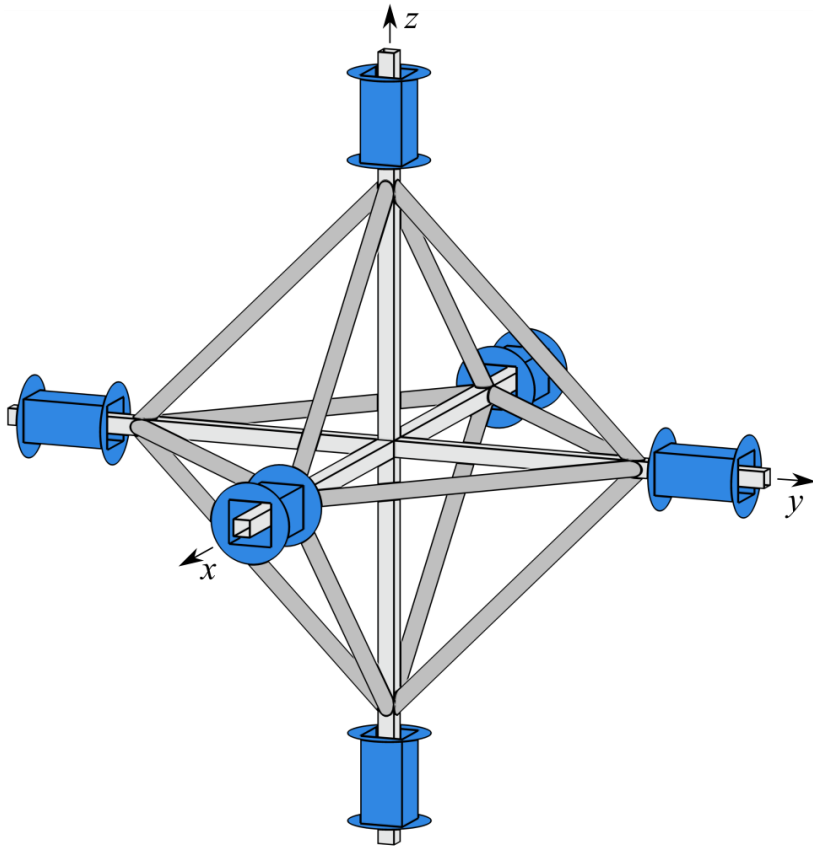
Test masses are levitated by a current induced along a tube.



Six test masses mounted a cube form a *tensor* gradiometer.

SOGRO:
Superconducting Omni-directional
Gravitational Radiation Observatory

Superconducting tensor gravitational wave detector



SOGRO

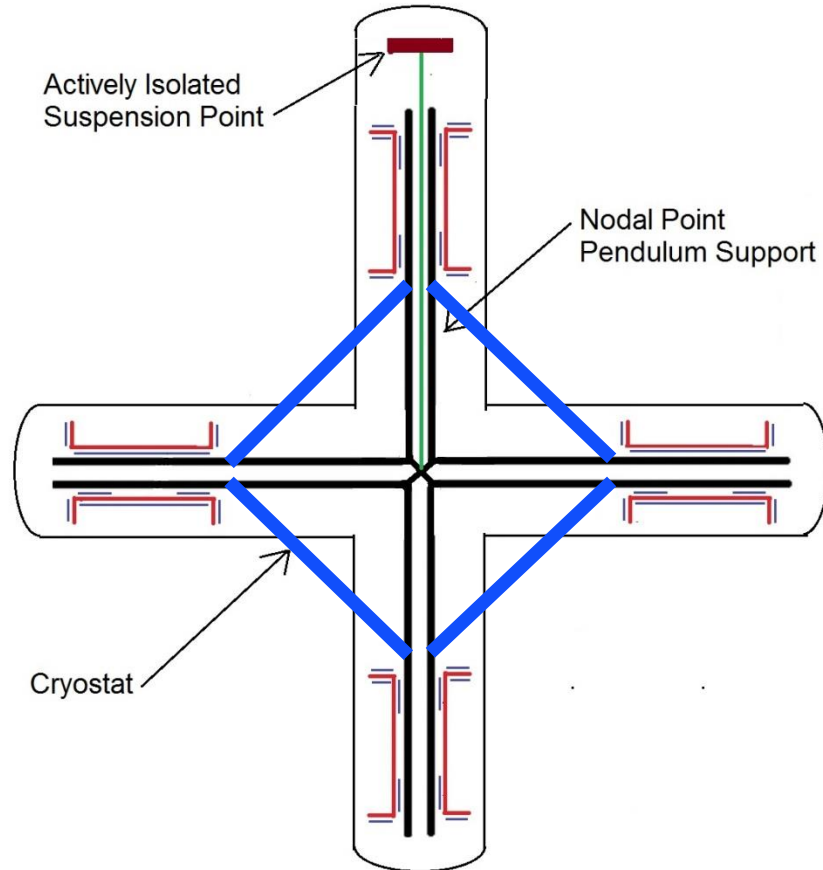
(Superconducting Omni-directional Gravitational Radiation Observatory)

- Each test mass has **three degrees of freedom**.
- Combining six test masses, a **tensor GW detector** is formed.

By detecting all six components of the Riemann tensor, the source direction (θ, ϕ) and wave polarization (h_+, h_\times) can be determined by a single detector. \Rightarrow **“Spherical” Antenna**

Suspension of SOGRO

- **Go underground** to reduce seismic and gravity gradient noise, as well as to be far away from moving objects.



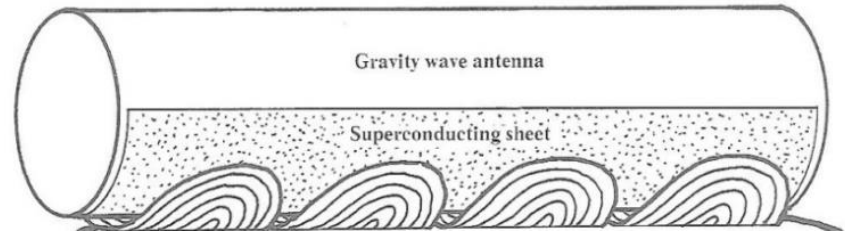
- **Nodal support** prevents odd harmonics from being excited.
- **Cable suspension** gives $f_r < 10^{-3}$ Hz for three angular modes.
⇒ **Active isolation unnecessary.**
- **25-m pendulum** gives $f_p = 0.1$ Hz for two horizontal modes.
⇒ **Provide passive isolation for high frequencies.**
- **Problem: Platform is not rigid enough.**
⇒ **Triangulate with struts.**

Alternative suspension: Optical rigid body
⇒ **Simpler cryogenics, larger baseline (up to 3 km?)**

Magnetic levitation

Stanford (1976)

1 ton Al antenna wrapped with Nb-Ti sheet was levitated magnetically.



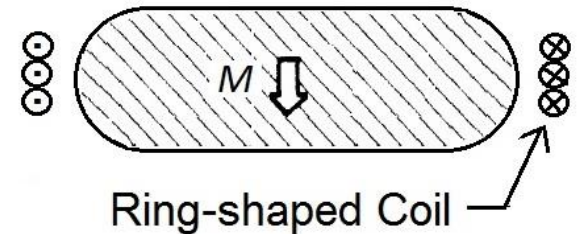
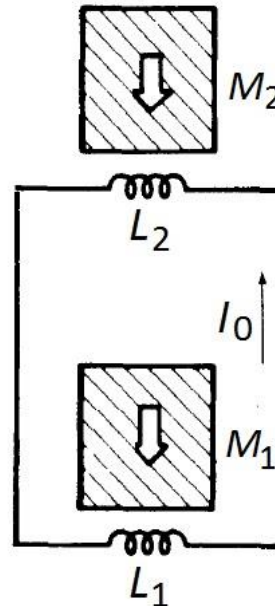
- Field required to levitate 5-ton mass:

$$\frac{B^2}{2\mu_0} A = Mg, B = \left(\frac{2\mu_0 Mg}{A} \right)^{1/2} = \left(\frac{2 \times 4\pi \times 10^{-7} \times 5 \times 10^3 \times 9.8}{10} \right)^{1/2}$$

$$= 0.11 \text{ T} < H_{c1} = 0.16 \text{ T (Nb at } T = 4 \text{ K).}$$

- The biggest challenge: To obtain symmetry, vertical DM resonance frequencies must also be reduced to 0.01 Hz.
- Employ “push-pull levitation” plus negative spring.

Push-pull levitation



Superconducting negative spring

Parametric transducer

- Near quantum-limited SQUIDs have $1/f$ noise below 100 kHz.
- Signal needs to be upconverted to $f_p = 100$ kHz by pumping at f_p by using a **parametric transducer**.
- An **inductance bridge** can be coupled naturally to a SQUID.

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Physical Review D

15 JANUARY 1986

Superconducting inductance-bridge transducer for resonant-mass gravitational-radiation detector

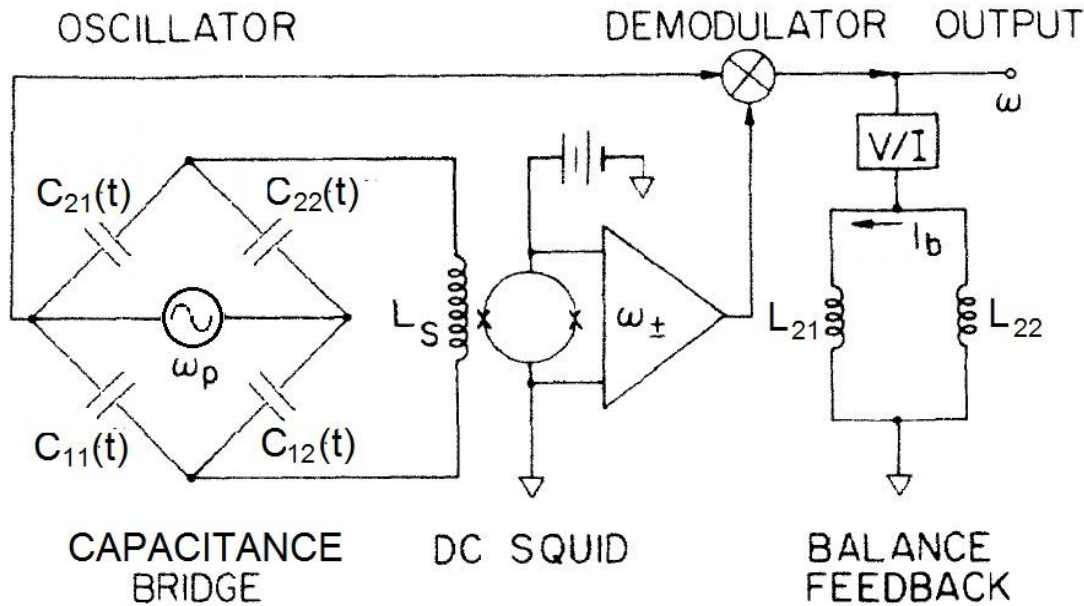
Ho Jung Paik*

Department of Physics, Stanford University, Stanford, California 94305

(Received 18 July 1985)

The sensitivity of cryogenic gravitational-radiation detectors is presently limited by the performance of the transducers. A superconducting ac-pumped inductance bridge is proposed as a new transducer for resonant-mass gravitational-radiation detectors. The impedance matrix of the transducer is computed to determine the input, output, and transfer characteristics of the electromechanical system. It is shown that the dissipative forces exerted on the proof mass by the bridge circuit through the two sidebands cancel each other almost exactly so that the Brownian-motion level is nearly unaffected by the electric sensing circuit. This implies that an effective energy coupling coefficient near unity could be used without being limited by the electrical Nyquist noise. With the parametric up conversion of the signal, the inductance bridge can be coupled to a nearly quantum-limited dc superconducting quantum interference device (SQUID). The sensitivity of the gravitational-radiation detector employing the new superconducting transducer is computed as a function of transducer parameters. It is shown that the proposed transducer, with modest values for its parameters, is capable of matching a high- Q gravitational-radiation antenna, cooled to 50 mK, to a nearly quantum-limited dc SQUID.

Tuned capacitor-bridge transducer



- Capacitor bridge coupled to a **near quantum-limited SQUID** thru S/C transformer.
- LC resonance increases **energy coupling β** by Q_p .
- **Oscillator noise is rejected by the bridge balance.**

⇒ **Maintain precise bridge balance by feedback.**

Parametric gain: $G_p^\pm = |Z_{\pm 1} / Z_{1\pm}| = \omega_\pm / \omega_m \approx \omega_p / \omega_m$

Energy coupling constant: $\beta_\pm \approx \frac{CE_p^2 Q_p}{M\omega^2}$

$$E_N(f) = \frac{k_B T \omega_D}{Q_D} + \frac{|\omega^2 - \omega_D^2|}{\omega_p} \left(1 + \frac{1}{\beta^2}\right)^{1/2} k_B T_N$$

Achievable detector noise

- For CW signal with impedance-matched bridge transducer,

$$S_h(f) = \frac{8}{ML^2\omega^4} \left\{ \frac{k_B T \omega_D}{Q_D} + \frac{|\omega^2 - \omega_D^2|}{\omega_p} \left(1 + \frac{1}{\beta^2} \right)^{1/2} k_B T_N \right\}, \quad k_B T_N = n\hbar\omega_p$$

Parameter	SOGRO 1	SOGRO 2	Method Employed (SOGRO 1 /2)
Each mass M	5 ton	5 ton	Nb square tube
Separation L	30 m	100 m	Over "rigid" mounting platform
Antenna temp T	1.5 K	0.1 K	Superfluid He / dilution refrigerator
DM frequency f_D	0.01 Hz	0.01 Hz	Magnetic levitation w/ negative spring
DM quality factor Q_D	10^8	10^9	Surface polished pure Nb
Signal frequency f	0.1-10 Hz	0.1-10 Hz	Detector noise computed at 1 Hz
Pump frequency f_p	50 kHz	50 kHz	Tuned capacitor bridge transducer
Amplifier noise no. n	200	10	Near-quantum-limited SQUID
Detector noise $S_h^{1/2}(f)$	$2 \times 10^{-20} \text{ Hz}^{-1/2}$	$2 \times 10^{-21} \text{ Hz}^{-1/2}$	Two phase development

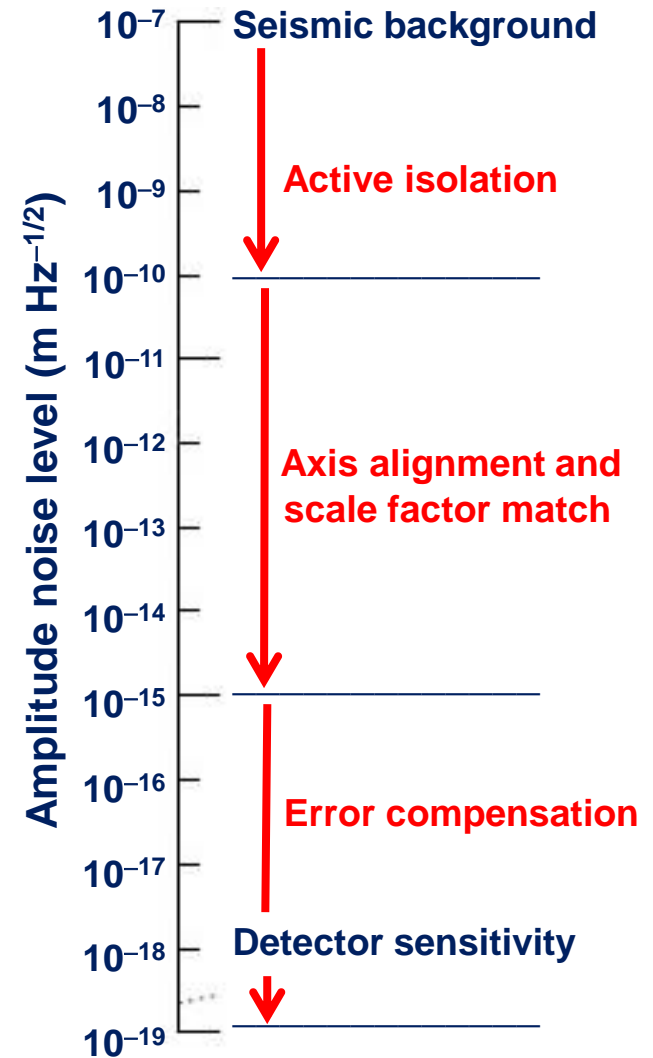
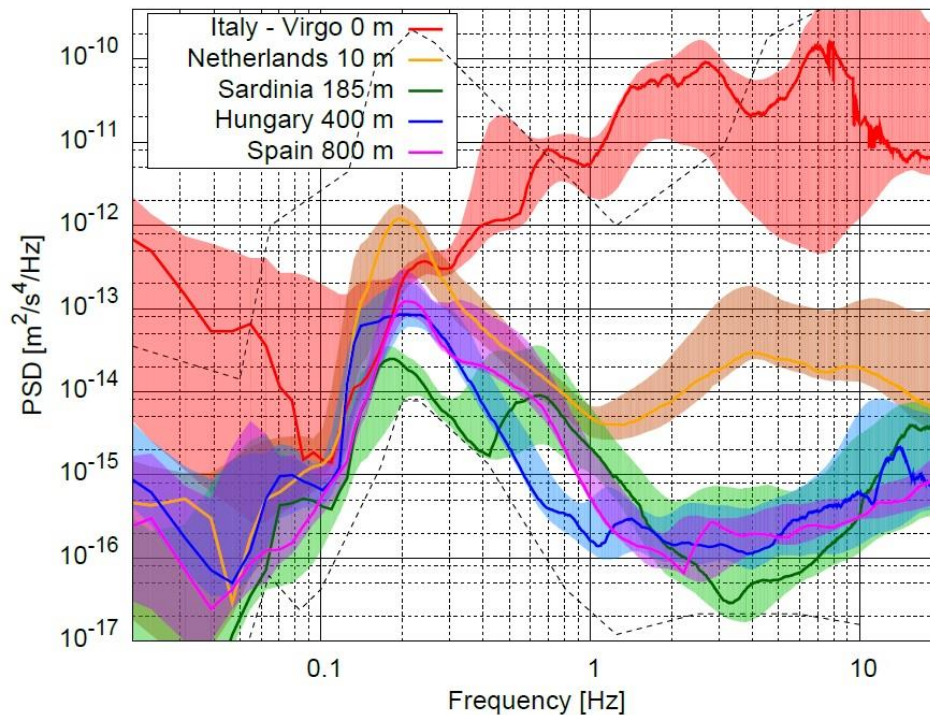
- SOGRO can also be operated as tunable resonant detector.**

$$S_h(f) = \frac{8}{ML^2\omega_D^3} \left[\frac{k_B T}{Q_D} + \frac{k_B T_N}{Q_D} \left(\frac{\omega_D}{\omega_p} \right) \right] \approx \frac{8}{ML^2\omega_D^3} \frac{k_B T}{Q_D}$$

Seismic noise

- **Seismic noise of bedrock: $10^{-7} \text{ m Hz}^{-1/2}$**
- **Could be reduced to the required level by combining **active isolation** with **CM rejection** of the detector.**
- **20-m pendulum with nodal support**
⇒ **Passive isolation for $f > 0.1 \text{ Hz}$.**

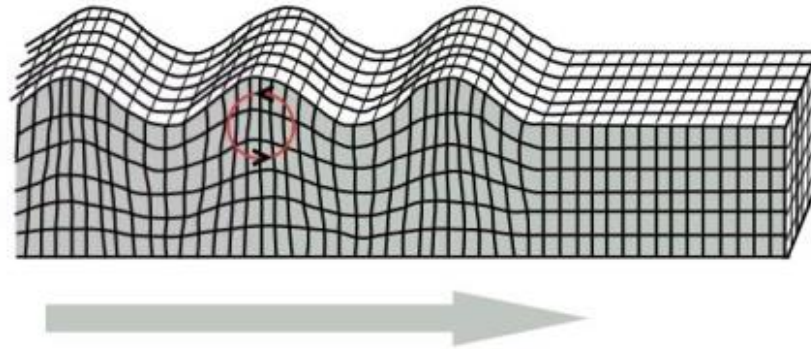
Seismic noise of underground sites



Newtonian gradient noise

- **The Newtonian noise from seismic and atmospheric density fluctuations cannot be shielded.**

Rayleigh Wave



- **GWs are transverse and cannot have longitudinal components whereas the Newtonian gradient does.**

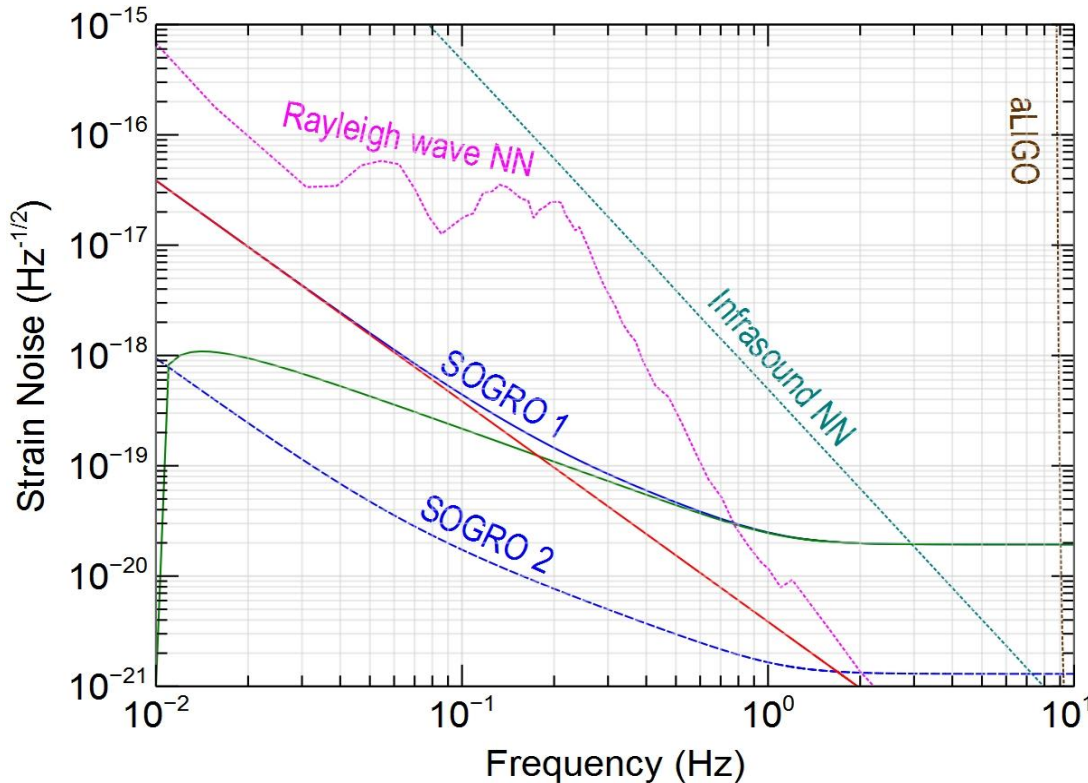
$$h_{GW}(t) = \begin{pmatrix} 0 & 0 & 0 \\ 0 & h_+(t) & h_\times(t) \\ 0 & h_\times(t) & -h_+(t) \end{pmatrix}, \quad h_{NG}(t) = \begin{pmatrix} h_{11}(t) & h_{12}(t) & h_{13}(t) \\ h_{12}(t) & h_{22}(t) & h_{23}(t) \\ h_{13}(t) & h_{23}(t) & h_{33}(t) \end{pmatrix}.$$

⇒ **Could be distinguished *in principle* from near-field Newtonian gradients.**

Sensitivity Goal

Sensitivity goals of SOGRO

SOGRO 1: $T = 1.5$ K, SOGRO 2: $T = 0.1$ K



- **Detector cooled to 1.5 K or 0.1 K and integrated with near-quantum-limited amplifiers.**
- **Seismic noise rejected by 10^{12} by combining CMRR, passive & active isolation.**
- **Newtonian noise rejected by $10^2 \sim 10^4$ using tensor nature of the detector and external sensors.**

Major challenges:

- **Large-scale cryogenics.**
- **Mitigation of Newtonian noise.**

Thank you!