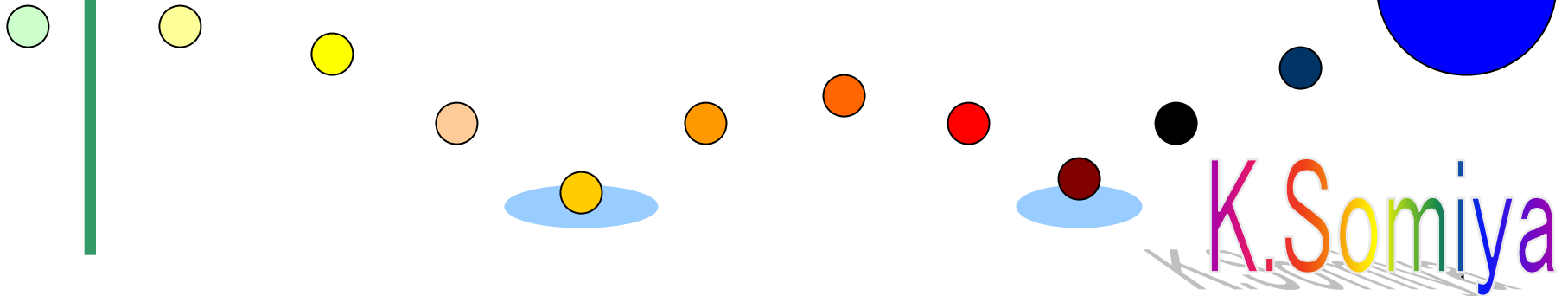


# Three different regimes in the verification stage according to the spring frequency

**K.Somiya, Y.Chen, S.Danilishin,  
H.Miao, H.Müller-Ebhardt, and H.Rehbein**

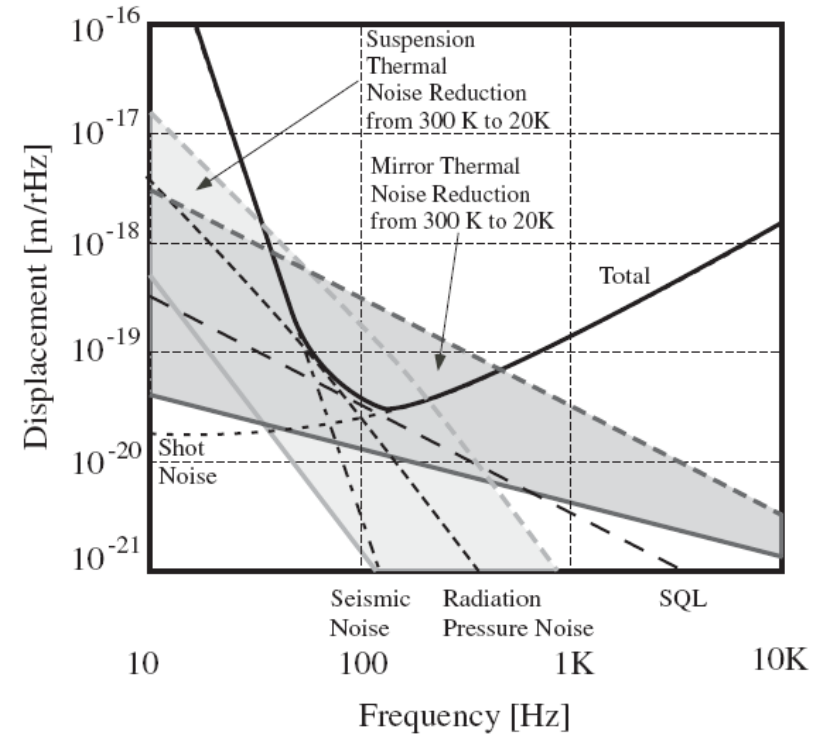
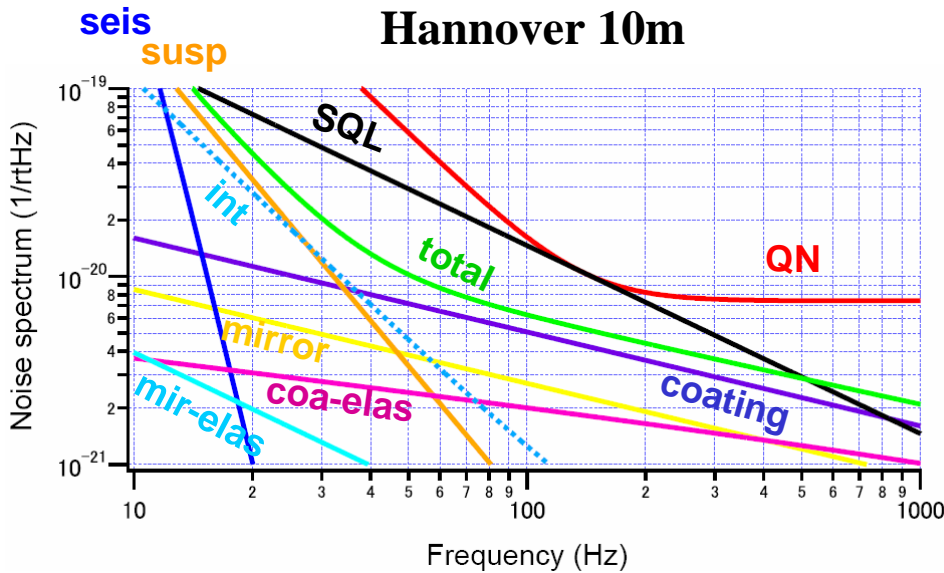
LSC Meeting @ Hannover  
2007.10.26



# MQM

CLIO (CQG 21 S1173)

Hannover 10m

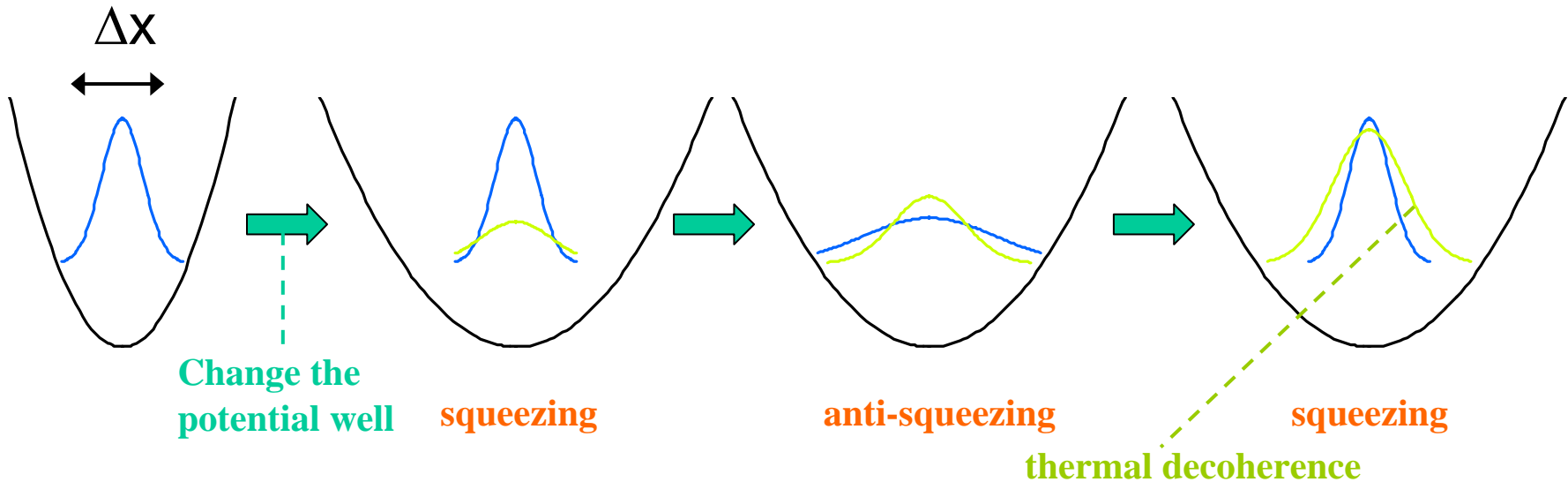


**Sub-SQL interferometer in a near future**



**We'll see the quantum behavior of a test mass  
(MQM=Macroscopic Quantum Measurement)**

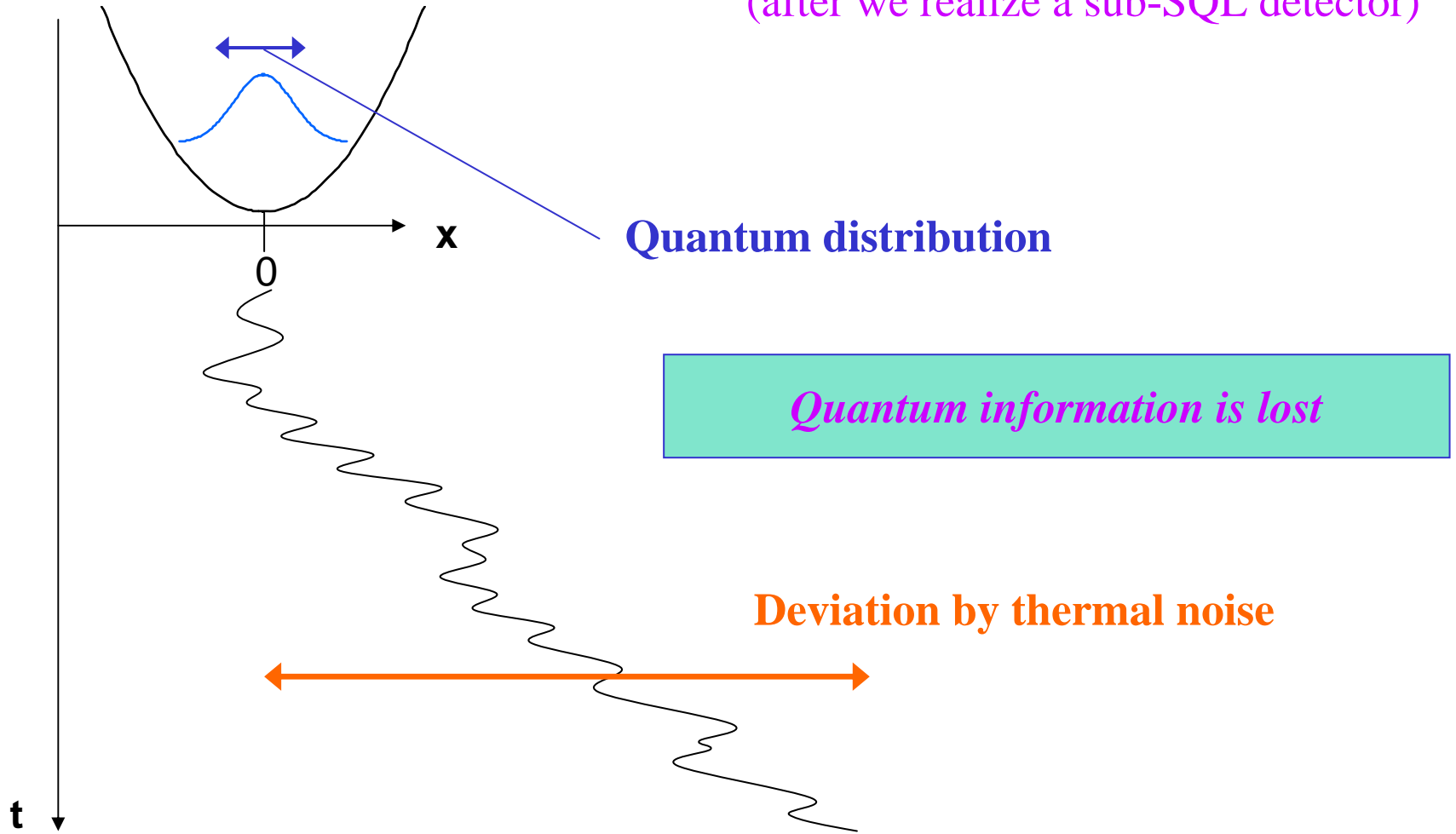
# Quantum behavior of a test mass



- Quantum breathing
- Thermal decoherence
- How can we see this?

# First we need to prepare a quantum state

(after we realize a sub-SQL detector)



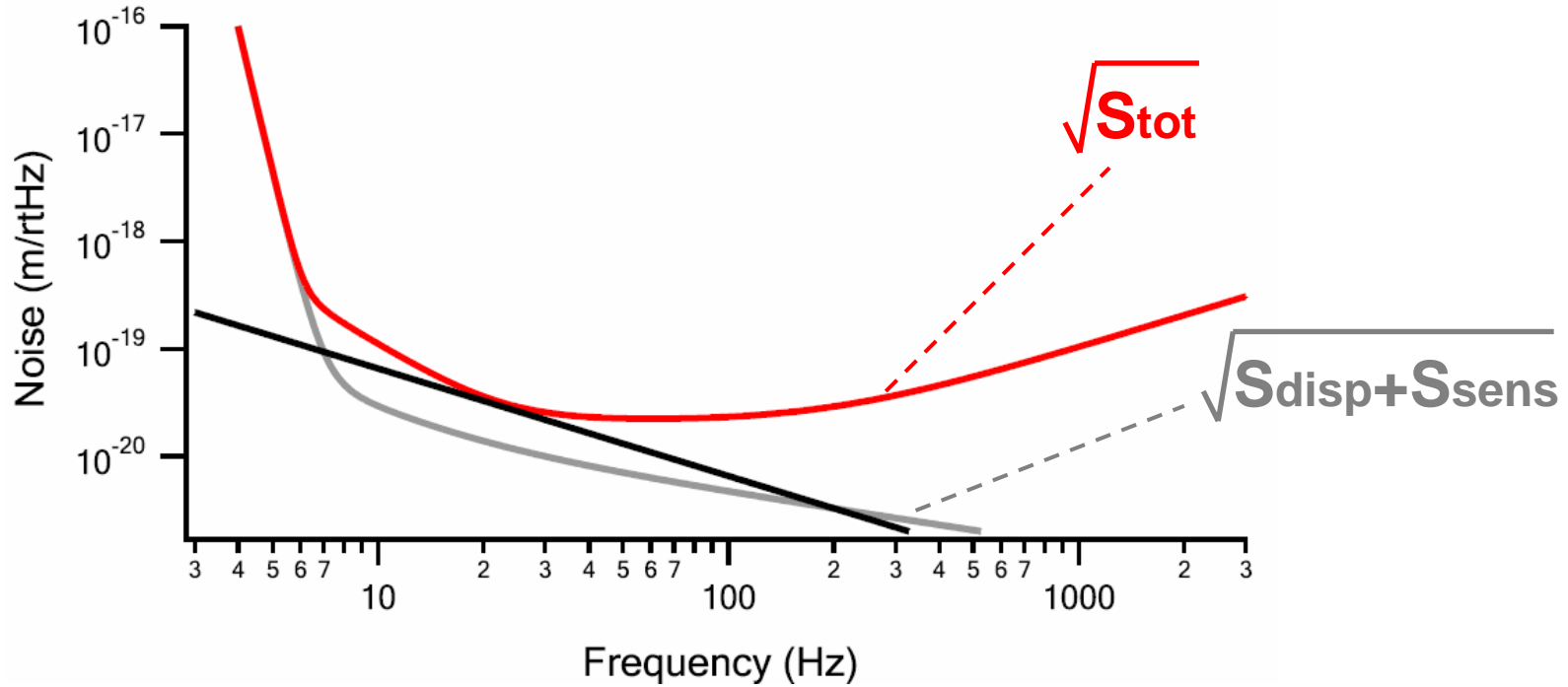
**We should identify the center of the wave function.**

# How to filter out the classical deviation



**Wiener Filter**

- designed from  $S_{disp}$  &  $S_{tot}$
- different filters for  $X$  &  $P$

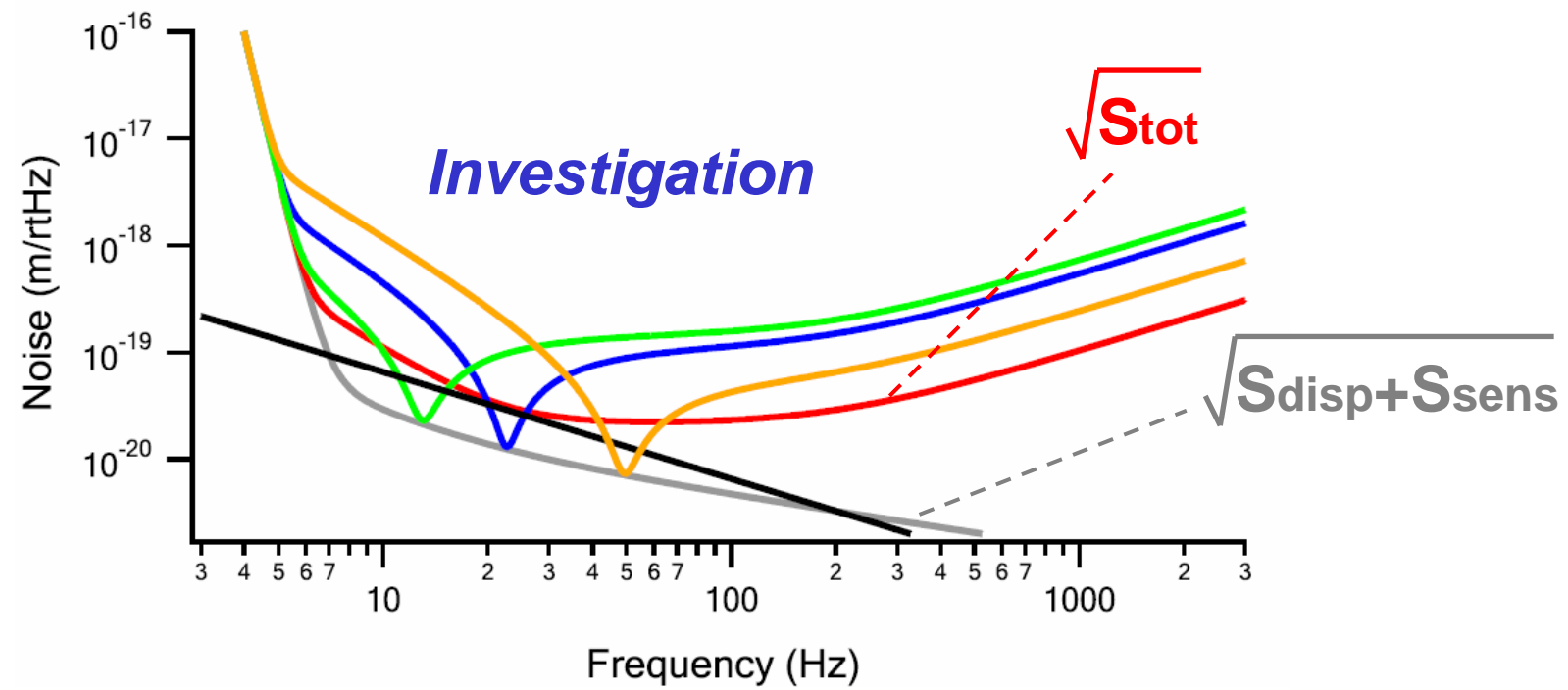


# How to filter out the classical deviation



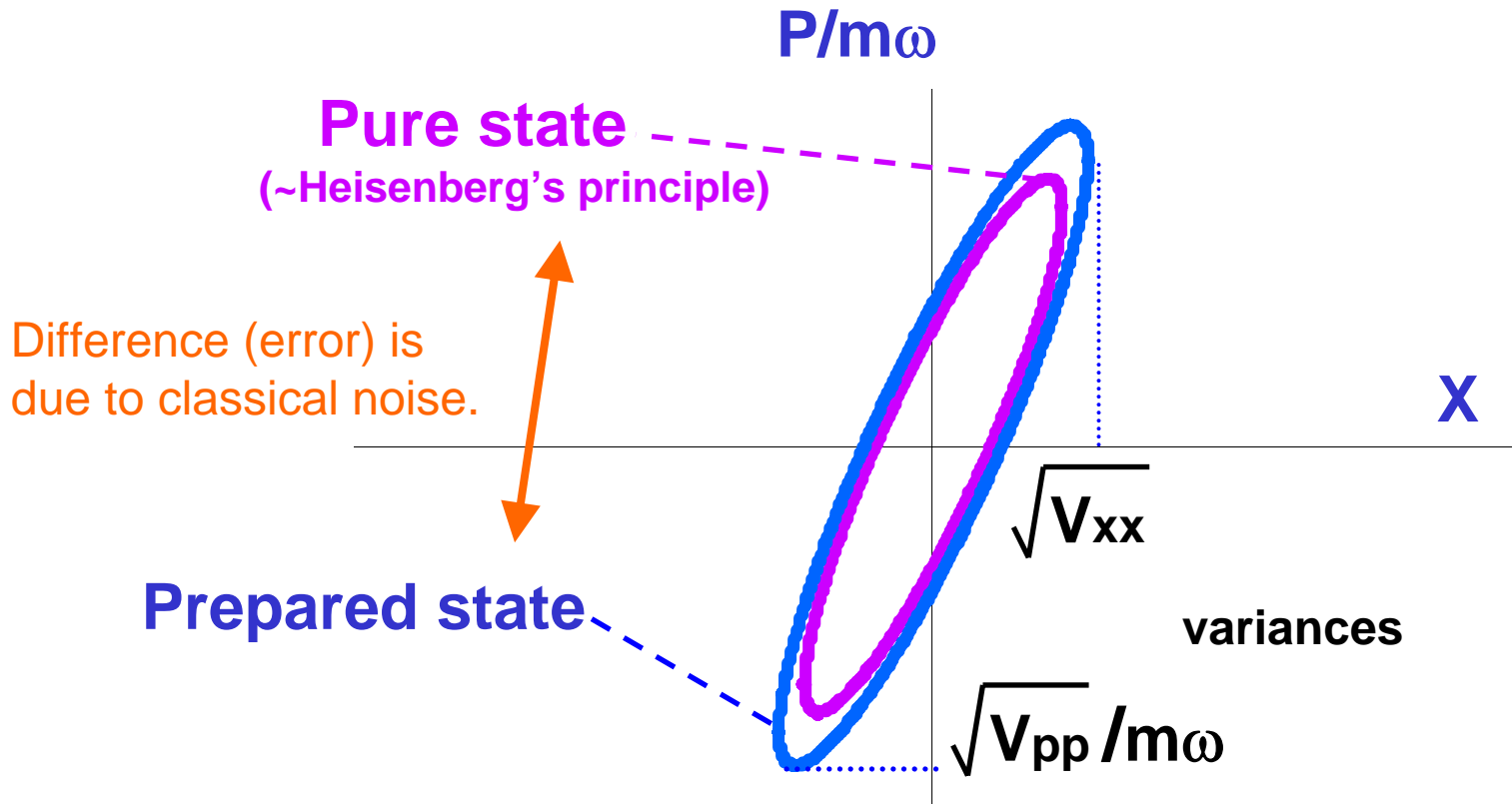
**Wiener Filter**

- designed from  $S_{disp}$  &  $S_{tot}$
- different filters for  $X$  &  $P$



The more the SQL is beyond classical noise, the closer the prepared state is to a pure quantum state.

# Prepared quantum state



Squeeze factor depends on the spring frequency  $\omega$

Now we should do measurement (verification)

## Then, the next step is...

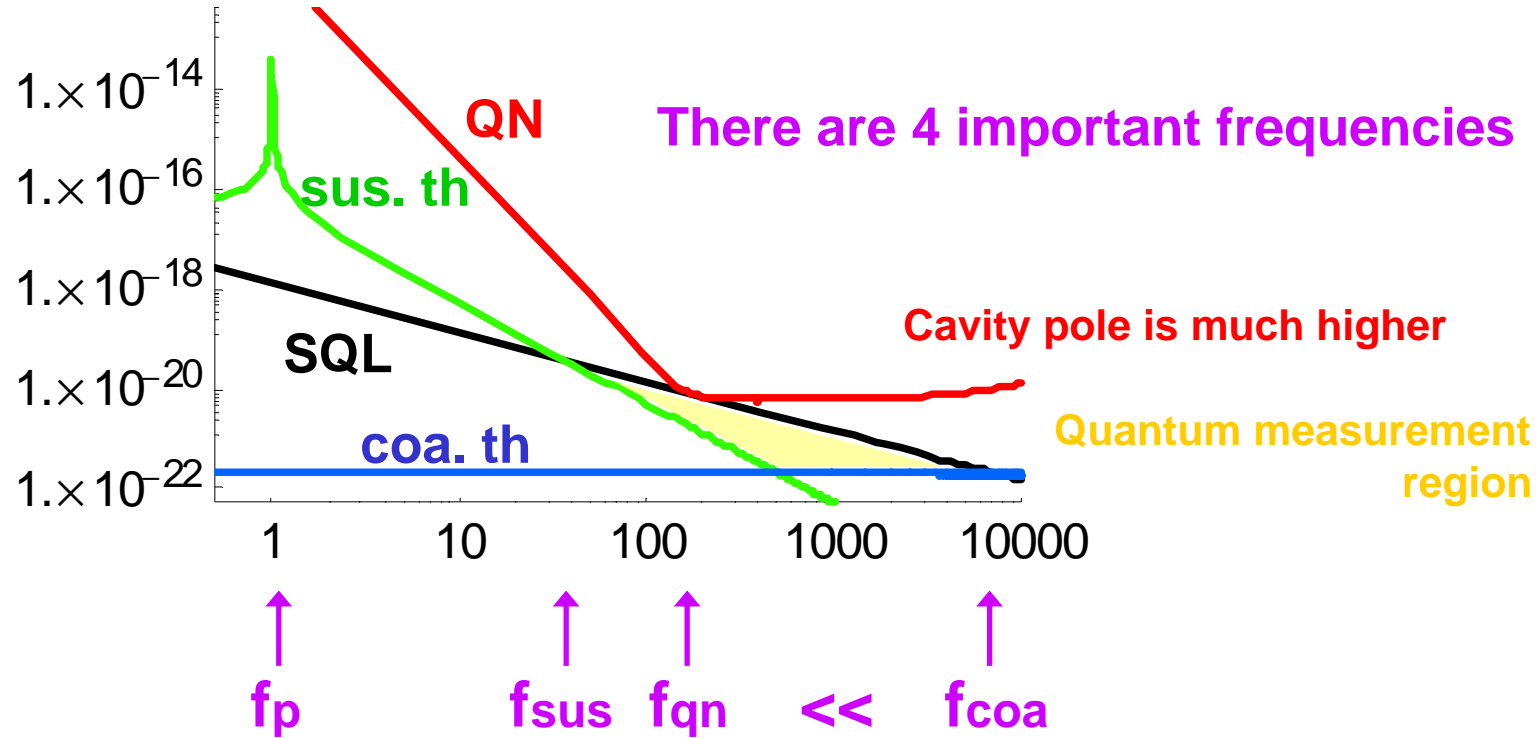
- (i) verify if it is almost a quantum state
- (ii) wait for one cycle ( $2\pi/\omega$ ) and see how different from (i)



**I'll show 3 different regimes  
depending on the spring frequency**

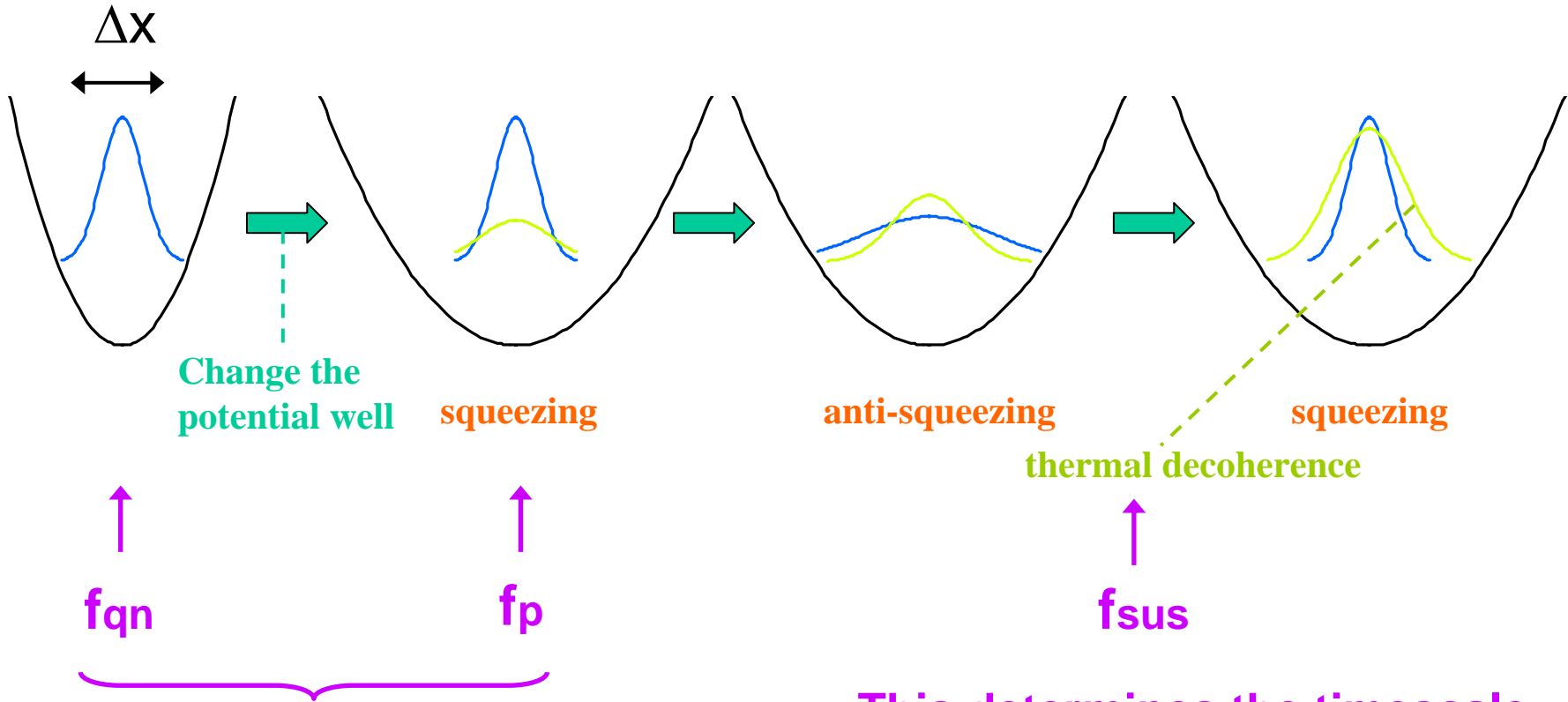


# Three situations in a simple model



- (1)  $f_p < f_{sus} < f_q$
- (2)  $f_{sus} < f_p < f_q$
- (3)  $f_{sus} < f_q < f_p$

# What do these frequencies mean?



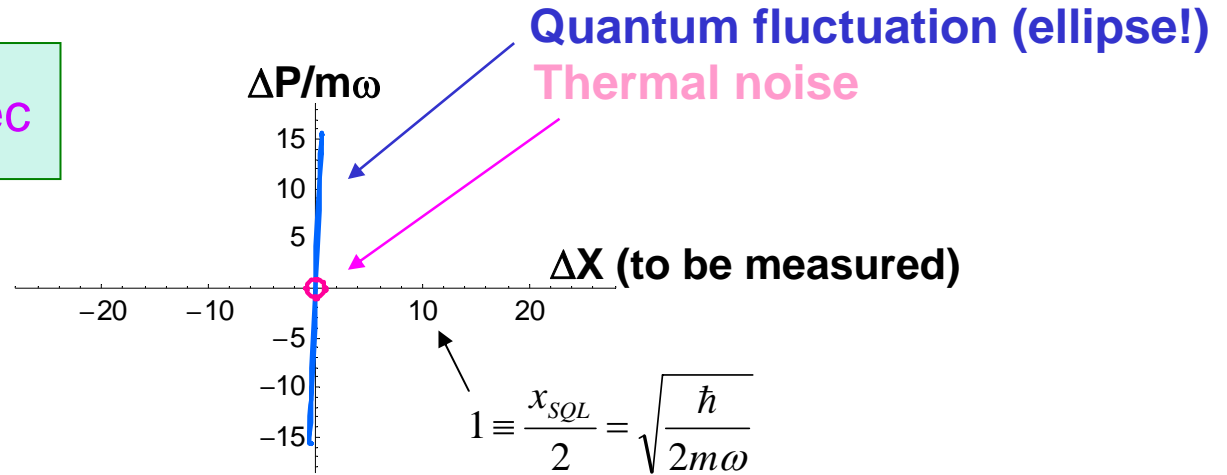
These determine the width of the potential wells.

This determines the timescale of thermal decoherence.

# (1) $f_p < f_{sus} < f_{qn}$ [pendulum]

(1Hz, 35Hz, 160Hz)

Snap shot at  $t=0.005\text{sec}$

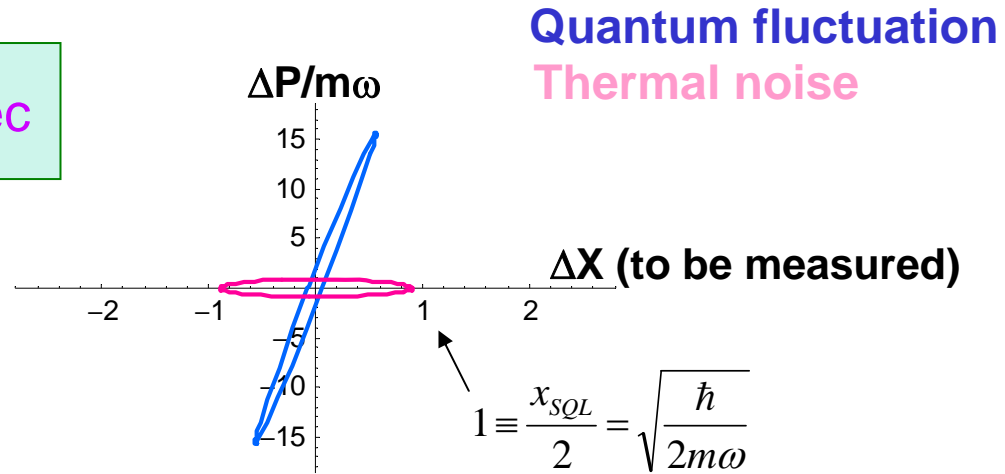


Highly squeezed ( $f_p \ll f_{qn}$ )

# (1) $f_p < f_{sus} < f_{qn}$ [pendulum]

(1Hz, 35Hz, 160Hz)

Snap shot at  $t=0.005\text{sec}$



Highly squeezed ( $f_p \ll f_{qn}$ )

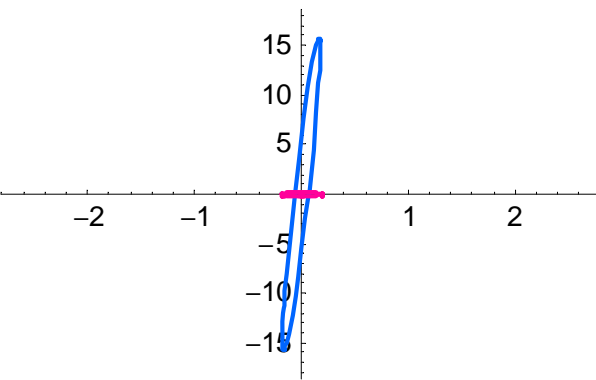
# (1) $f_p < f_{sus} < f_{qn}$ [pendulum]

(1Hz, 35Hz, 160Hz)

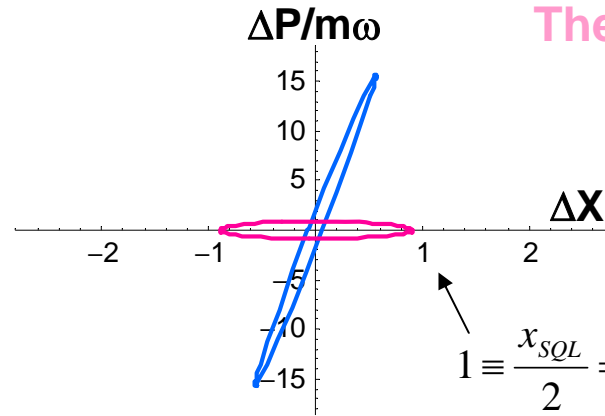
Evolution in time

Quantum fluctuation

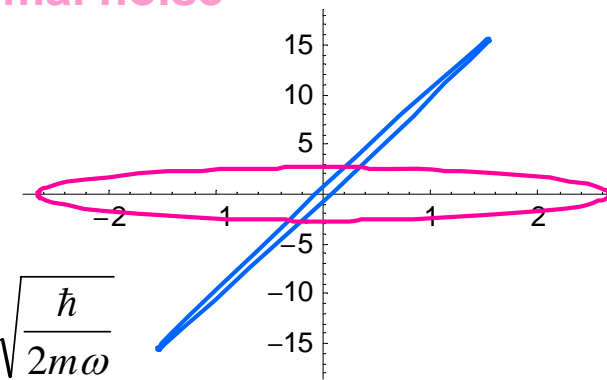
Thermal noise



t=0.001 sec



t=0.005 sec



t=0.015 sec

$$1 \equiv \frac{x_{SQL}}{2} = \sqrt{\frac{\hbar}{2m\omega}}$$

Thermal noise becomes bigger very soon.

But the information of P and X can be obtained quickly due to the strong squeezing (hope).

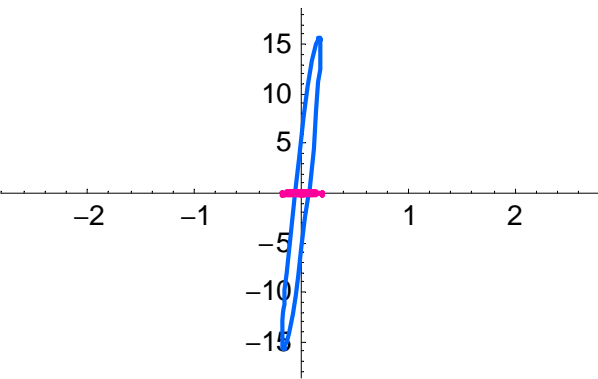
# (1) $f_p < f_{sus} < f_{qn}$ [pendulum]

(1Hz, 35Hz, 160Hz)

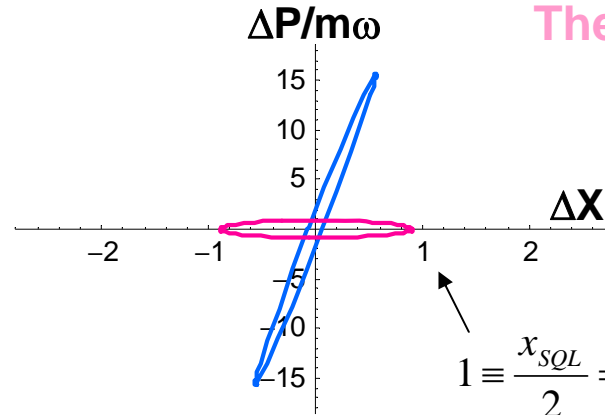
Evolution in time

Quantum fluctuation

Thermal noise

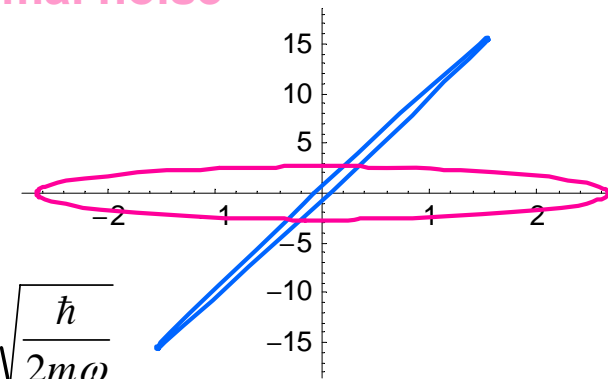


t=0.001 sec



t=0.005 sec

$$1 \equiv \frac{x_{SQL}}{2} = \sqrt{\frac{\hbar}{2m\omega}}$$



t=0.015 sec

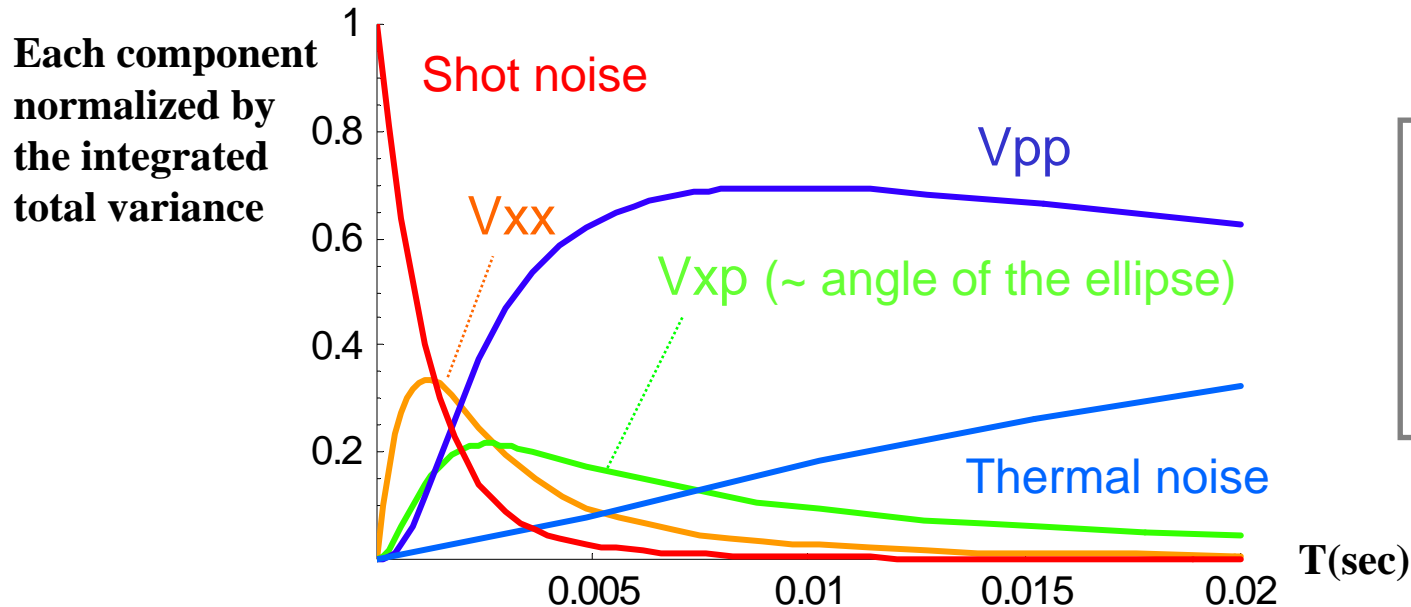
Note that these are snap shots.

We shall measure it ( $\Delta x$ ) for some amount of time.  
(integrate)

# (1) $f_p < f_{sus} < f_{qn}$ [pendulum]

(1Hz, 35Hz, 160Hz)

Let's see the integral of  $\langle X(t)X(t') \rangle$  over time T.



*notes*

$$\phi_{sus} = 2e-6$$

$$P_0 = 1W$$

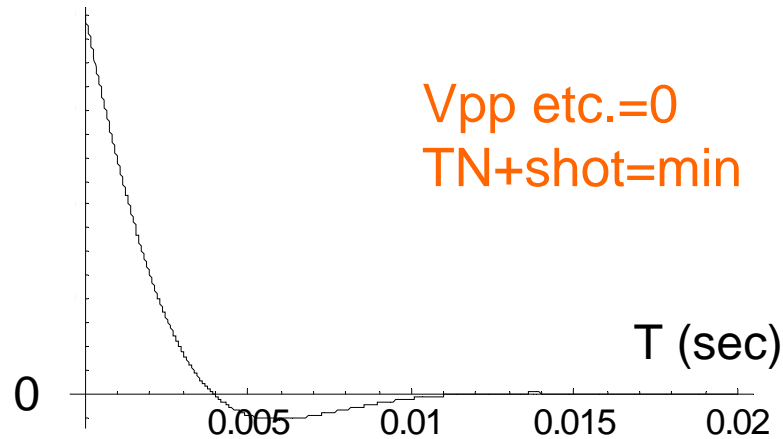
BAE (no RPN)

- At a peak of each,  $V_{xx}/V_{tot} = 34\%$  and  $V_{pp}/V_{tot} = 70\%$
- Optimal filter will make these number better

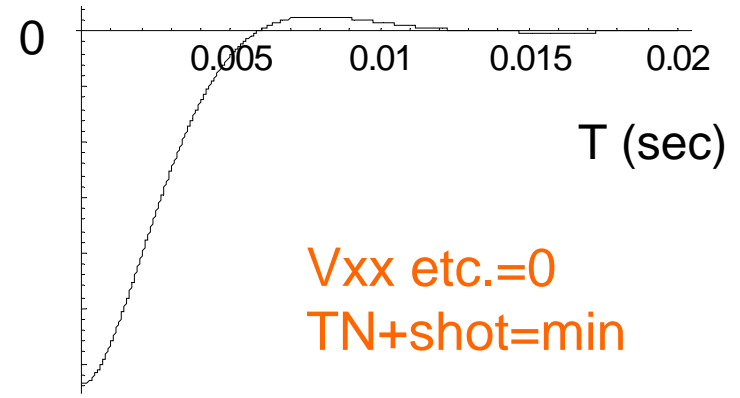
# (1) $f_p < f_{sus} < f_{qn}$ [pendulum]

## Optimal filters

for  $V_{xx}$



for  $V_{pp}$

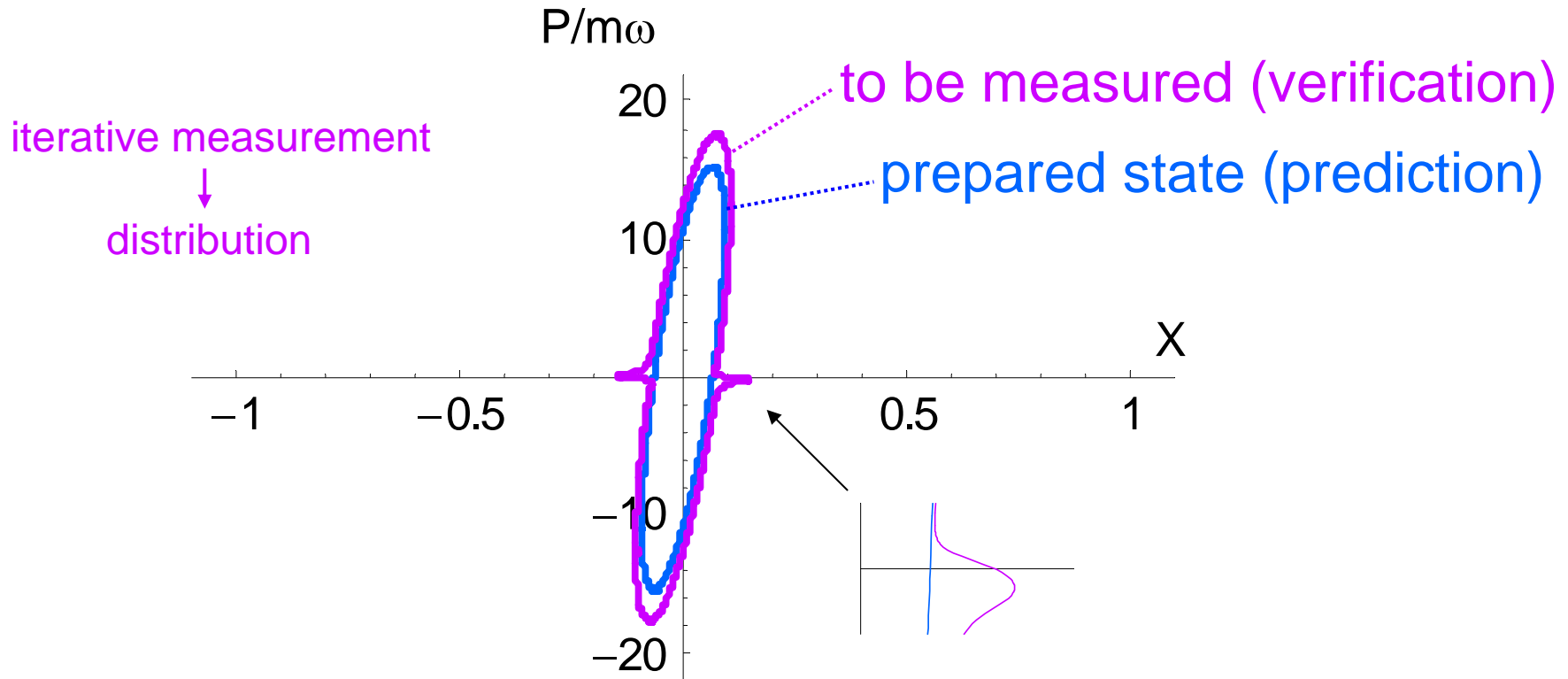


optimal function	<b>48%</b>	<b>76%</b>
(simple window)	34%	70%

We can use the filter for each quadrature between X and P.



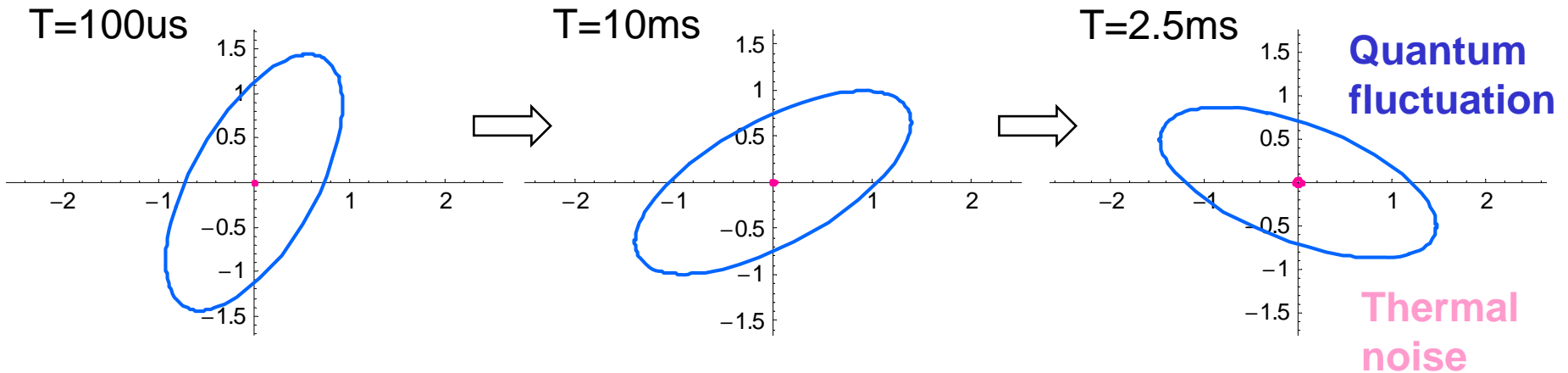
# (1) $f_p < f_{sus} < f_{qn}$ [pendulum]



If a quantum state is prepared correct,  
we'll be able to verify it in this way.

## (2) $f_{\text{sus}} < f_p < f_{\text{qn}}$ [spring]

(35Hz, 100Hz, 160Hz)



Modest squeezing ( $f_p \sim f_{\text{qn}}$ )

Thermal noise grows slowly ( $f_{\text{sus}} < f_p$ )

Let's see the integrated variances.

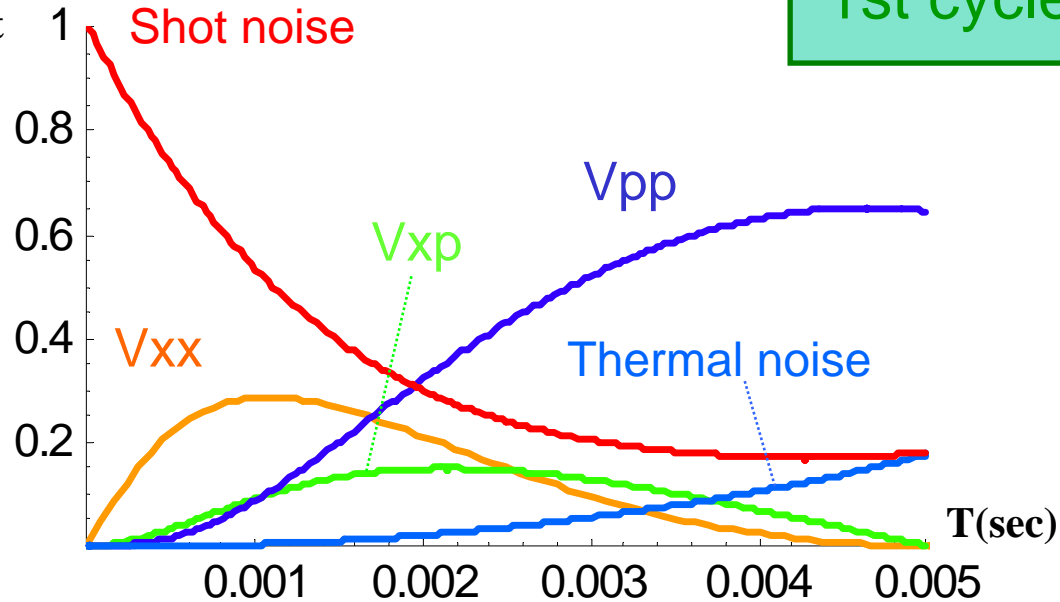
## (2) $f_{sus} < f_p < f_{qn}$ [spring]

(35Hz, 100Hz, 160Hz)

measurement starts at

1st cycle

Each component  
normalized by  
the integrated  
total variance



Reasonably fine.

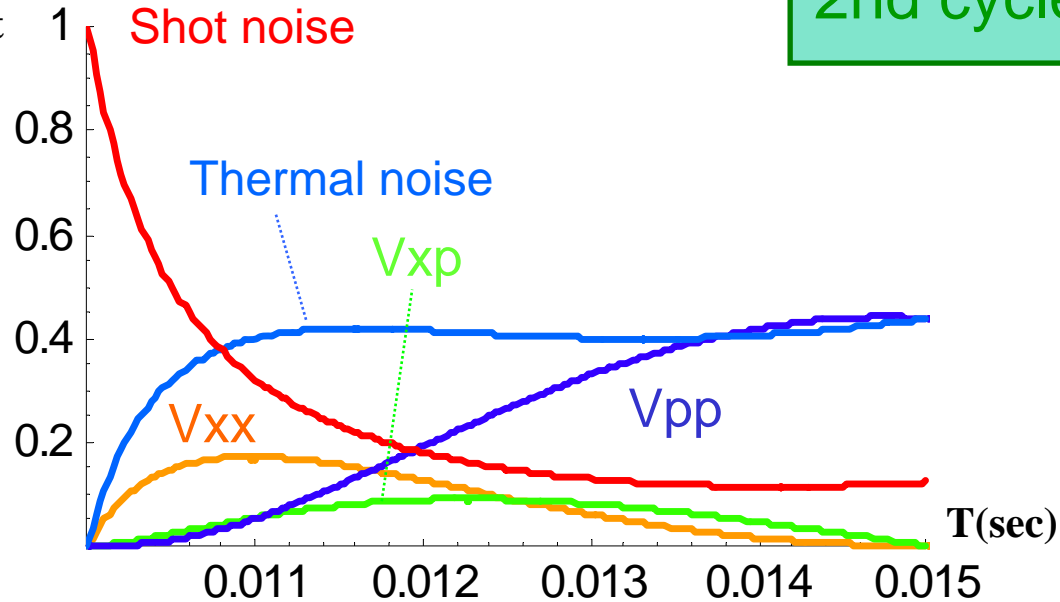
## (2) $f_{sus} < f_p < f_{qn}$ [spring]

(35Hz, 100Hz, 160Hz)

measurement starts at

2nd cycle

Each component  
normalized by  
the integrated  
total variance



Thermal decoherence is bigger than the 1<sup>st</sup> cycle.

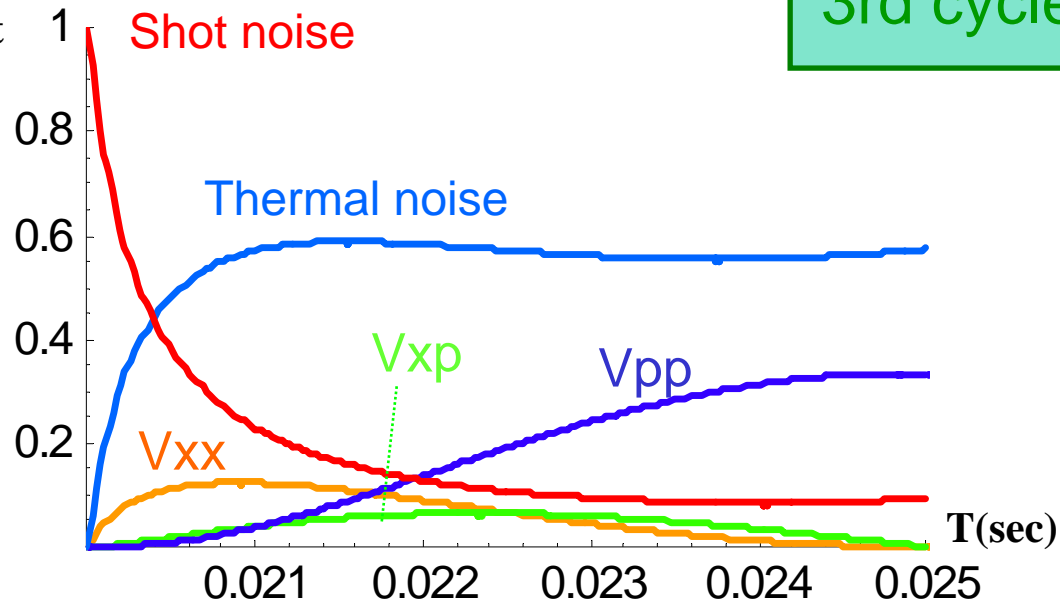
## (2) $f_{sus} < f_p < f_{qn}$ [spring]

(35Hz, 100Hz, 160Hz)

measurement starts at

3rd cycle

Each component  
normalized by  
the integrated  
total variance



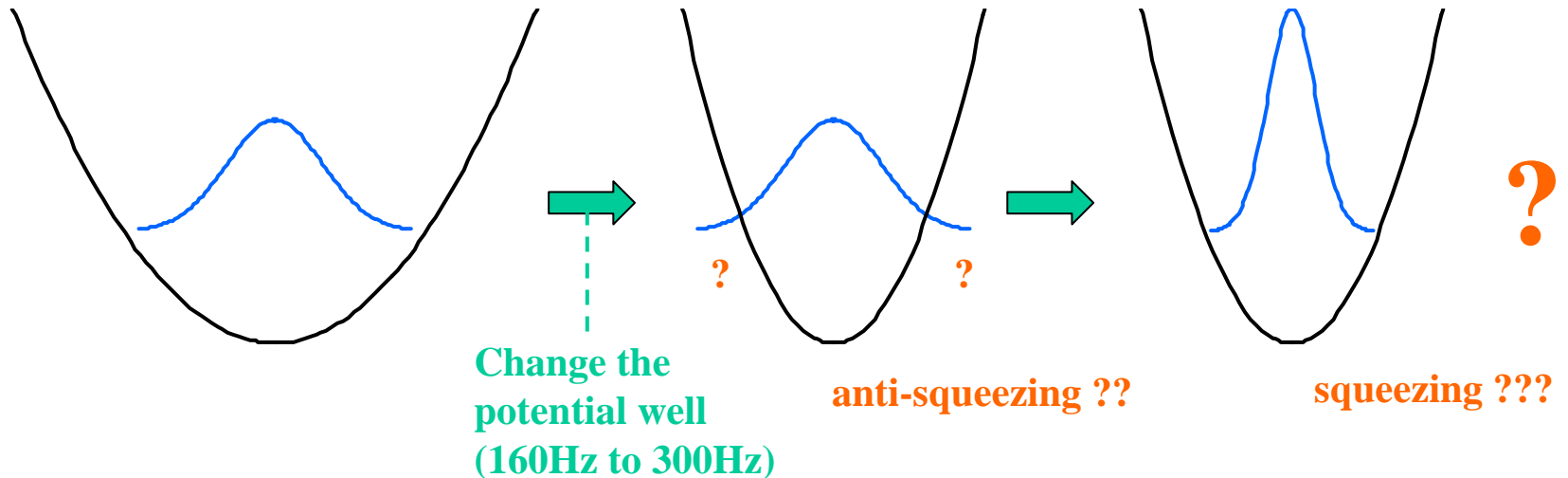
We'd like to see the results with the optimal filter.

➔ However it is not available yet... I'll do soon.

### (3) $f_{sus} < f_{qn} < f_p$ [spring]

(35Hz, 160Hz, 300Hz)

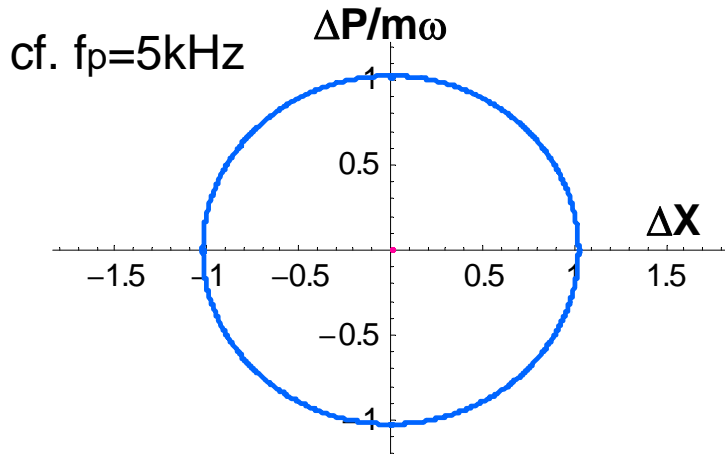
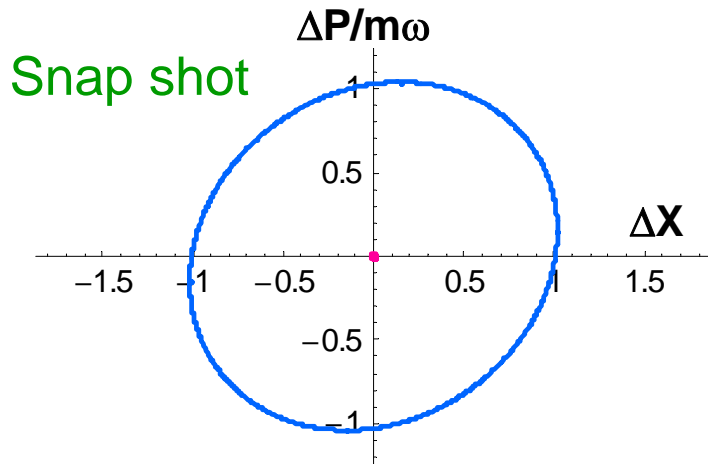
This is a subtle situation...



Let's see the calculated result.

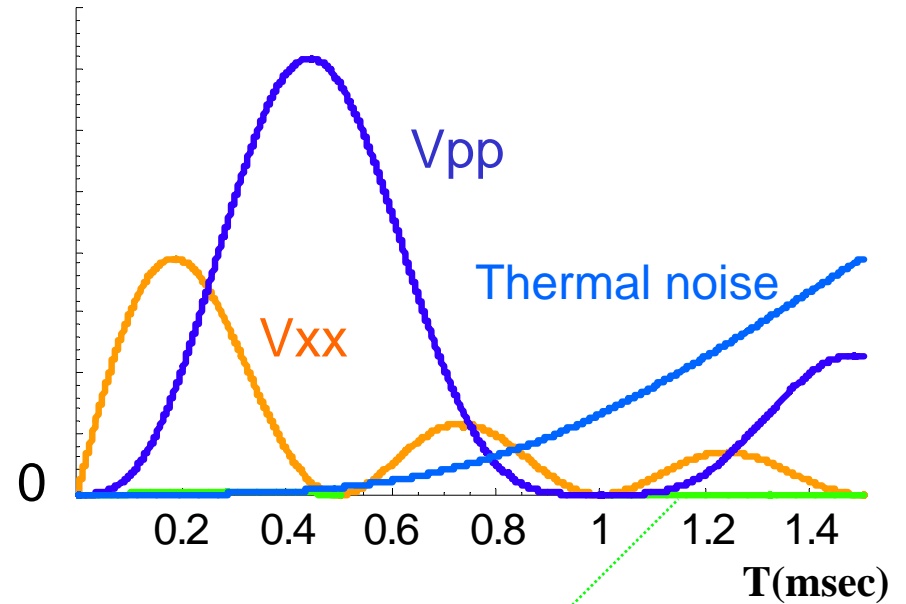
# (3) $f_{sus} < f_{qn} < f_p$ [spring]

(35Hz, 160Hz, 300Hz)



Integral (prop.)

Shot noise is much higher



$V_{xp}$  is quite small  
(as it's not squeezed)

So, this situation is not interesting...

# Conclusion

- The pendulum case seems *hopeful*:
  - $f_p \ll f_q \longrightarrow$  high squeezing
  - $f_p < f_{\text{sus}} \longrightarrow$  thermal decoherence grows fast
- The middle-freq spring case seems *interesting*:
  - $f_p \sim f_q \longrightarrow$  not so high squeezing
  - $f_{\text{sus}} < f_p \longrightarrow$  measurement can be done several times
- The high-freq spring case was *different* from what I thought:
  - $f_q \ll f_p \longrightarrow$  almost no squeezing (no anti-squeezing)

Some more calculation should be done.  
Note that this is with a very simple model.