

## Numerical optimization of AdvLIGO for simultaneous multiple source types detection

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LSC-VIRGO Meeting, QND workshop  
Hannover, October 26, 2006



Max-Planck-Institut  
für Gravitationsphysik  
(Albert-Einstein-Institut)



# Outline

1 Advanced LIGO interferometers conventional optimization

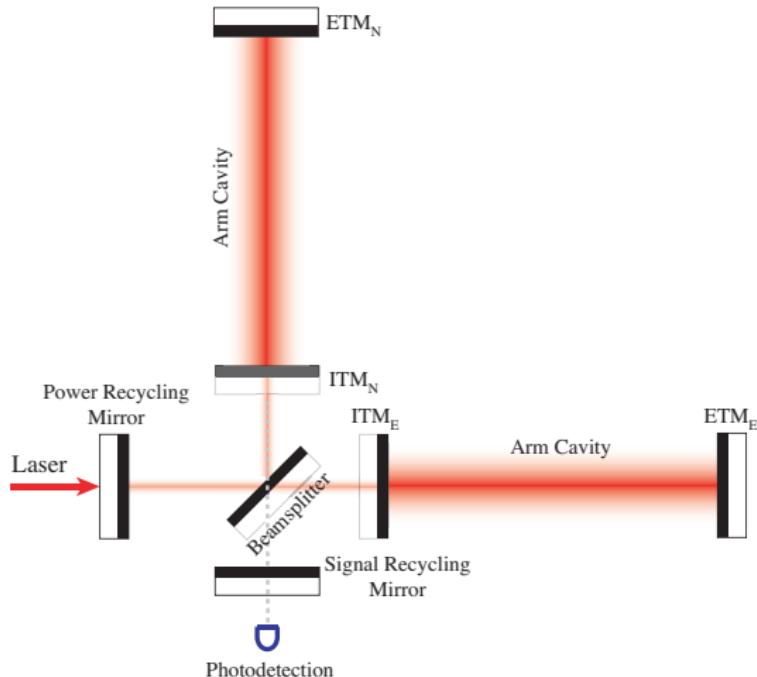
2 Simultaneous optimization for different sources

- Optimization for HF pulsars
- Optimization for GW bursts

3 Conclusion



# Advanced LIGO signal recycled interferometers (SRI)



Characteristic parameters of SRI:

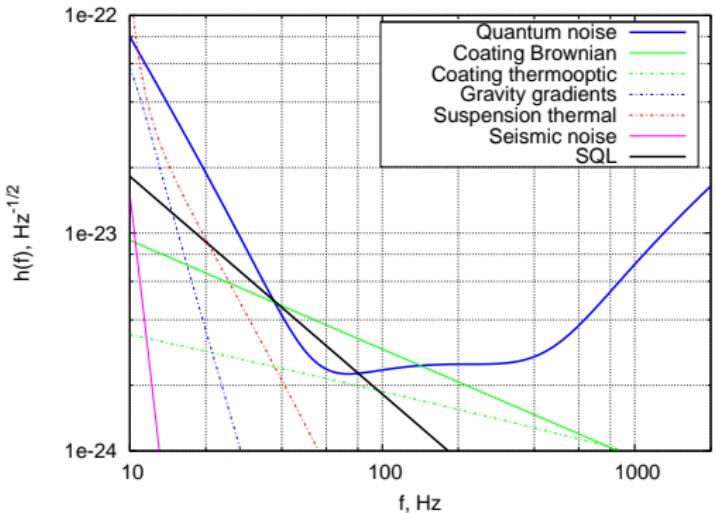
Quantum noise of SRI  $\xrightarrow{}$  3 parameters

- Bandwidth  
 $\gamma \iff (\rho_{SRM}, \phi_{SRM})$ ;
- Detuning  
 $\delta \iff (\rho_{SRM}, \phi_{SRM})$
- Homodyne quadrature angle  $\varphi$

A. Buonanno, Y. Chen, Phys.  
Rev. D 67, 062002 (2003)



# Advanced LIGO signal recycled interferometers (SRI)



Optical configuration:

Certain set of 3 numbers ( $\gamma, \delta, \varphi$ )

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*A. Buonanno, Y. Chen, Phys. Rev. D 67, 062002 (2003)*

# How to evaluate which configuration of SRI is better?

## Detection range for NS-NS binaries

Evident candidate  $\Rightarrow$  NS-NS binaries detection range:

$$r_{\text{NS-NS}} = \left( \frac{2}{15} \frac{G^{5/3}}{\pi^{4/3} c^3} \frac{\mathcal{M}^{5/3}}{\bar{\rho}_0^2} \int_{f_{\min}}^{f_{\max}} \frac{df}{f^{7/3} S_h(f)} \right)^{1/2}.$$

Conventional way  $\Rightarrow$  optimal configuration for NS-NS binaries

Quantum part of  $S_h(f)$  depends on  $(\gamma, \delta, \varphi)$

$$S_h(f) = S_h^q(f, \gamma, \delta, \varphi) + S_h^{cl}(f)$$



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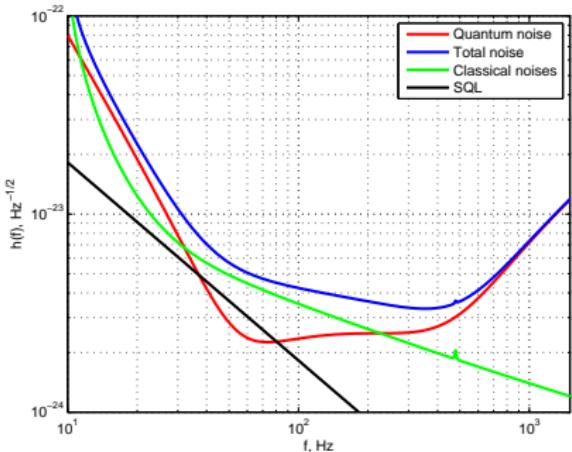
Conventional way  $\Rightarrow$  optimal configuration for NS-NS binaries

Optimizing  $r_{\text{NS-NS}}$  in  $(\gamma, \delta, \varphi)$   $\Rightarrow$  optimal configuration

$$\min[r_{\text{NS-NS}}(\gamma, \delta, \varphi)] \longrightarrow (\gamma_{\text{opt}}, \delta_{\text{opt}}, \varphi_{\text{opt}})_{\text{NS-NS}}$$



# Why one might need another optimization?



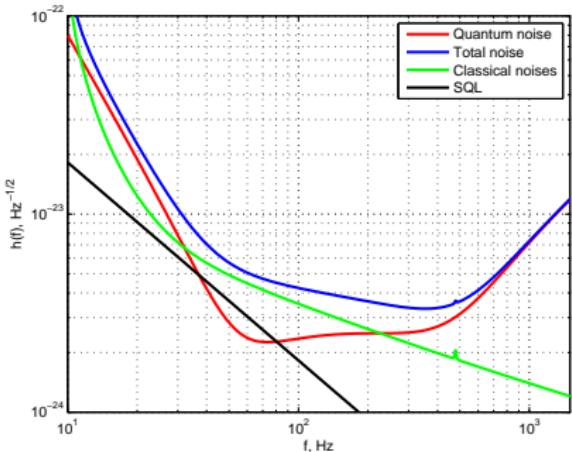
$$\gamma_{\text{opt}}^{NS} \simeq 960 \text{ sec}^{-1}, \delta_{\text{opt}}^{NS} \simeq 2090 \text{ sec}^{-1},$$
$$\varphi_{\text{opt}}^{NS} \simeq -1.02 \text{ rad}$$

## Disadvantages of conventional way

- ➊ Lock-in to specific GW sources type;
- ➋ High classical noises at medium frequencies  $\Rightarrow$  weak dependence on optical parameters;
- ➌  $|h_{\text{NS}}(f)|^2 \sim \frac{1}{f^{7/3}}$   $\Rightarrow$  no optimization at high frequencies.



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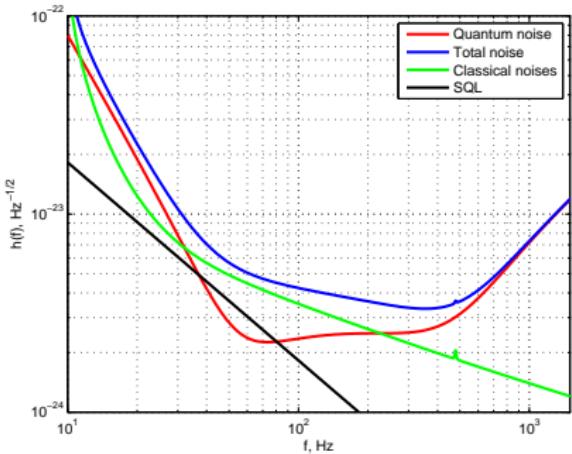
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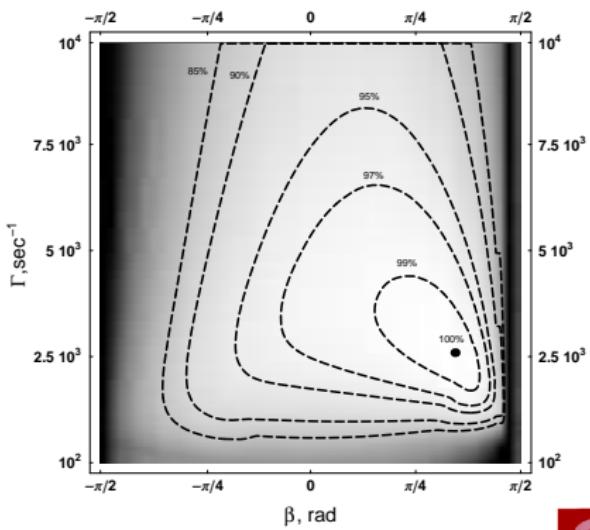
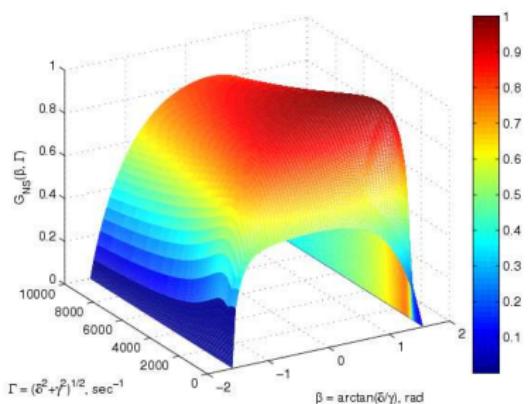
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# Why one might need another optimization?

Dependance of DR for NS-NS binaries on optical configuration is rather weak!



## Redefinition of optical parameters

$$\gamma, \delta \implies \Gamma = \sqrt{\gamma^2 + \delta^2}, \beta = \arctan \frac{\delta}{\gamma}$$

Why don't we sacrifice a little bit  
in sensitivity to NS-NS  
and improve sensitivity to other sources?



# What sources can be included into consideration?

- ➊ High frequency pulsars with  $f_{\text{rot}} \gtrsim 300$  Hz
- ➋ Bursts of GWs from, say, supernovae explosions

## Criteria for optimization

Relative gain for specific source:

$$G_{\text{source}} = \frac{\text{SNR}_{\text{source}}(\Gamma, \beta, \varphi)}{\text{SNR}_{\text{source}}(\Gamma_{\text{opt}}^{\text{NS}}, \beta_{\text{opt}}^{\text{NS}}, \varphi_{\text{opt}}^{\text{NS}})}$$



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## SNR for HF pulsars

$$\text{SNR}_{\text{puls}} \propto \left[ \frac{1}{S_h(f_{\text{puls}}, \Gamma, \beta, \varphi)} \right]^{1/2}$$



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## SNR for GW bursts

$$\text{SNR}_{\text{burst}} \propto \left[ \int_{f_{\min}}^{f_{\max}} \frac{df}{S_h(f, \Gamma, \beta, \varphi)} \right]^{1/2}$$



# How do we optimize?

## Our optimization procedure:

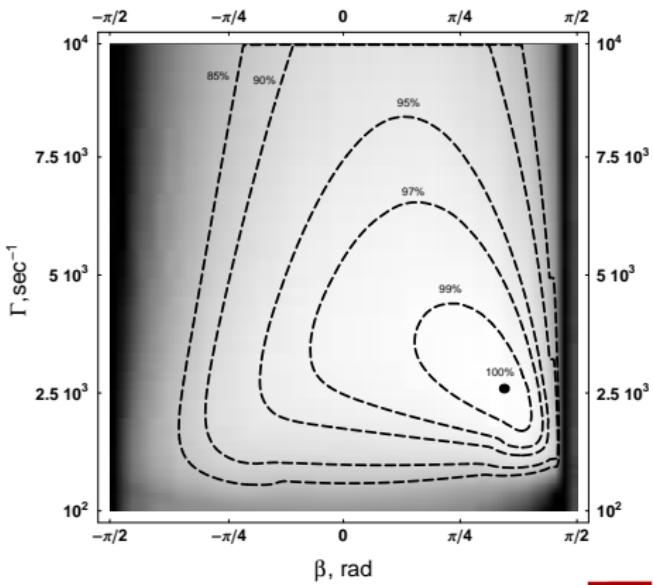
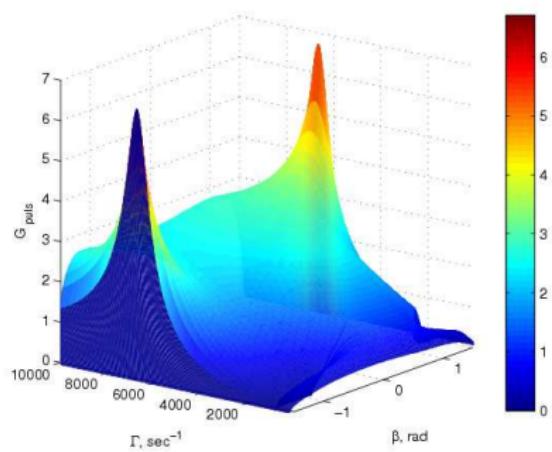
- ① Calculate and optimize over all possible configurations  $\Rightarrow G_{\text{NS}}(\gamma, \beta, \phi)$
- ② Optimize  $G_{\text{NS}}(\gamma, \beta, \varphi)$  over all  $\varphi \Rightarrow G_{\text{NS}}(\gamma, \beta)$
- ③ Fix price to pay (% of loss in NS sensitivity) and maximize  $G_{\text{puls}}$  and  $G_{\text{burst}}$  taking this price into account

*BENCH Software*, <http://ilog.ligo-wa.caltech.edu:7285/advligo/Bench/>



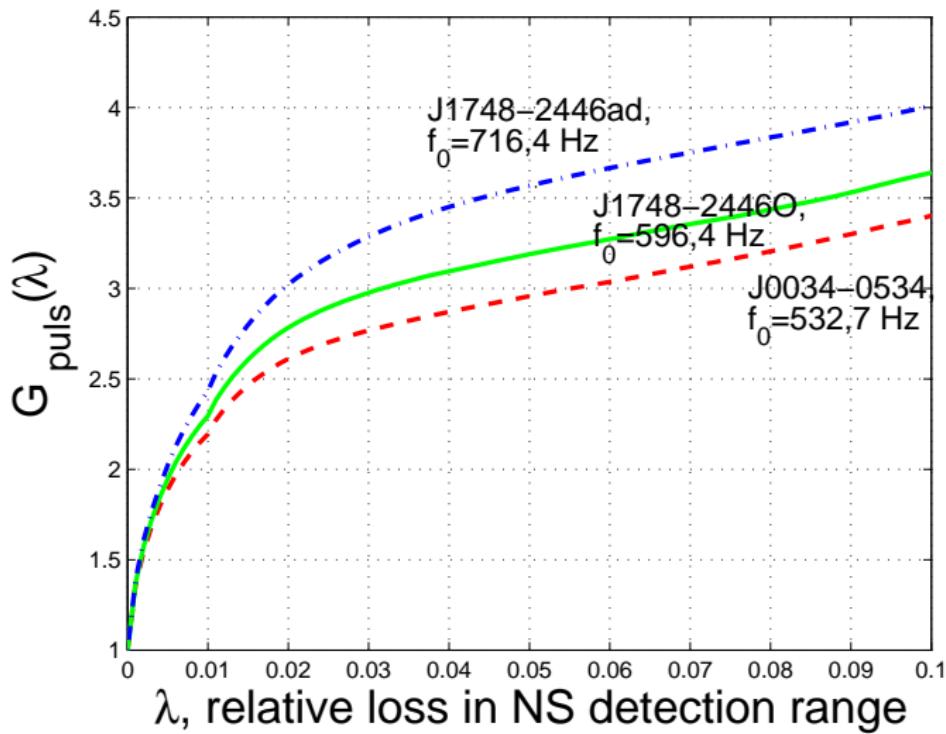
# Optimization for HF pulsars

$G_{puls}(\Gamma, \beta)$  for pulsar J0034-0534 with rotation frequency  $f_{\text{rot}} = 532.7$  Hz



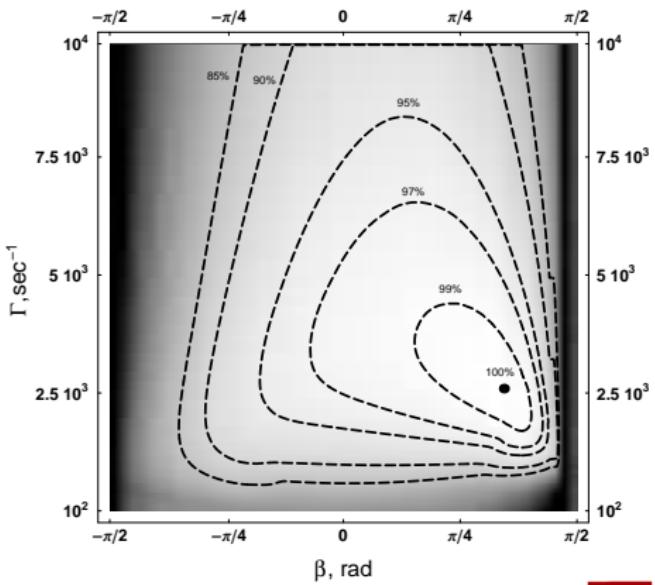
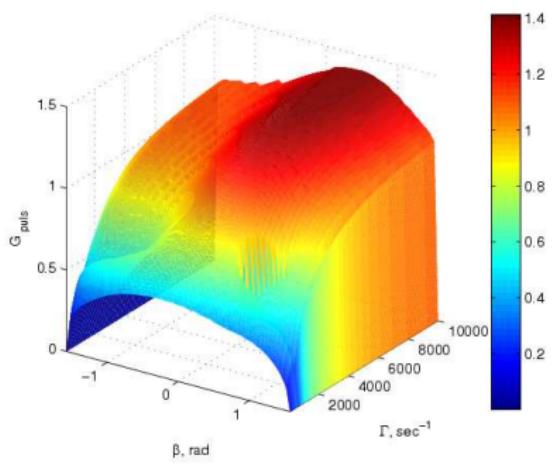
# Optimization for HF pulsars

Measure of sacrifice in NS sensitivity:  $\lambda = 1 - G_{\text{NS}}(\Gamma_{\text{puls}}, \beta_{\text{puls}}, \varphi_{\text{puls}})$



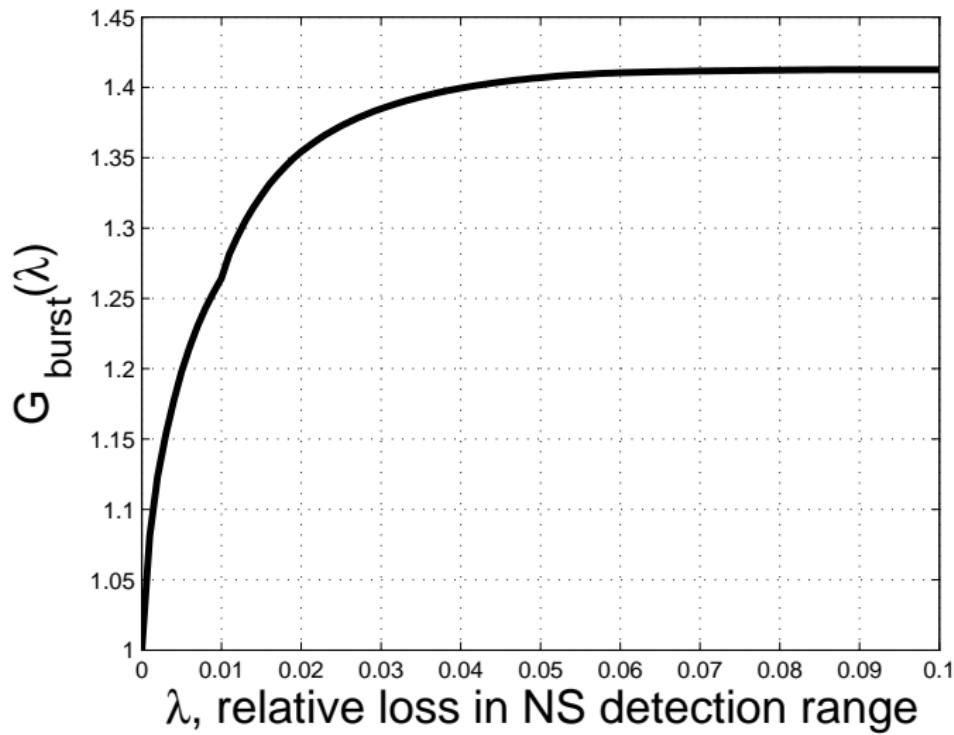
# Optimization for GW bursts

Relative loss in sensitivity for GW bursts  $G_{burst}(\Gamma, \beta)$



# Optimization for GW bursts

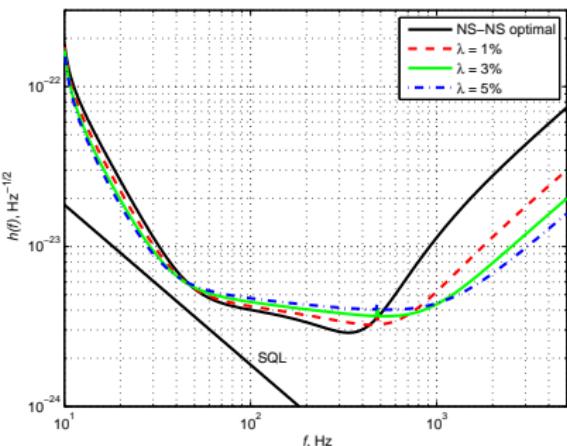
Measure of sacrifice in NS sensitivity:  $\lambda = 1 - G_{\text{NS}}(\Gamma_{\text{burst}}, \beta_{\text{burst}}, \varphi_{\text{burst}})$



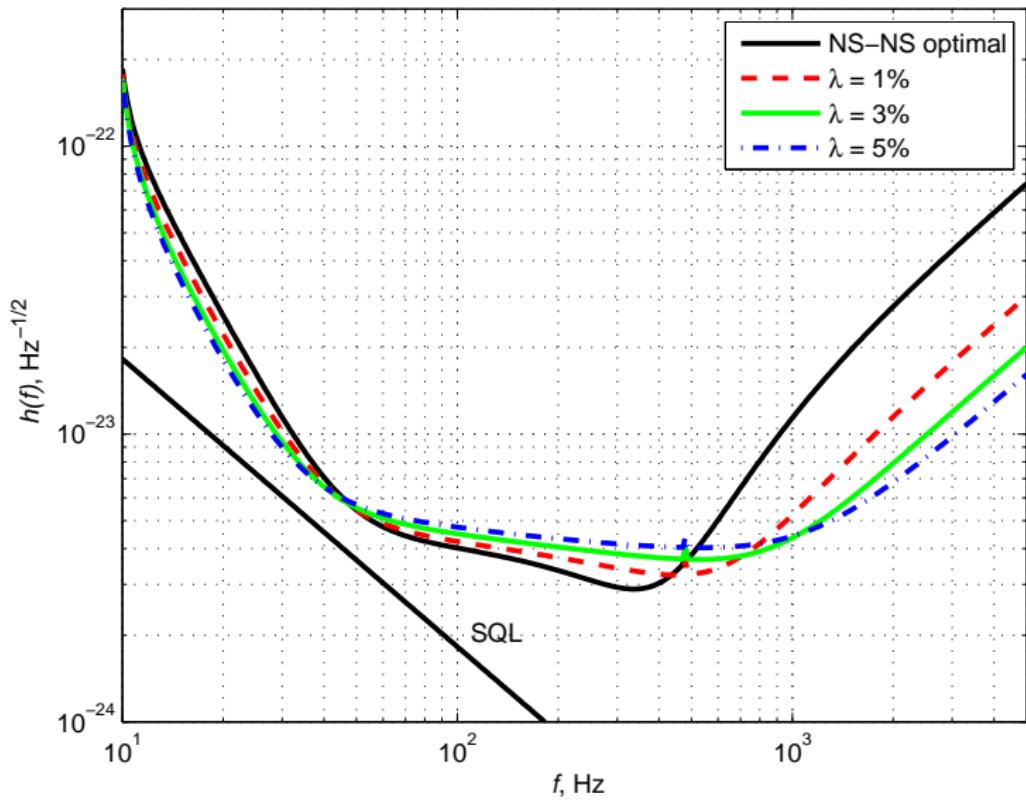
# Optimized total noise curves and corresponding configurations

For given  $\lambda$  one can choose optimal configuration of SRI

Rel. NS loss $\lambda$ ,	0.01	0.02	0.03
$\gamma$ , sec $^{-1}$	2294	3160	4076
$\delta$ , sec $^{-1}$	2775	3063	3229
$\varphi$ , rad	-0.59	-0.44	-0.34
$G_{\text{burst}}$	1.3	1.36	1.39
$G_{\text{puls}}(f_1)$	2.35	2.64	2.77
$G_{\text{puls}}(f_2)$	2.46	2.81	2.99
$G_{\text{puls}}(f_3)$	2.6	3.06	3.31



# Optimized total noise curves and corresponding configurations



- ➊ Optimization for NS-NS binaries seems to be non-optimal for HF sources of GWs
- ➋ Optimization of AdvLIGO SRI sensitivity aiming at several different sources is demonstrated
- ➌ Significant improvement in sensitivity for HF pulsars and GW bursts along with small loss in sensitivity to NS is shown



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**THANK YOU  
FOR YOUR ATTENTION!!!**