

**Latest news on
continuous-wave sources**

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LIGO-G000038-00-D

Summary

Lots of papers on r-modes, one monster paper on electron capture.

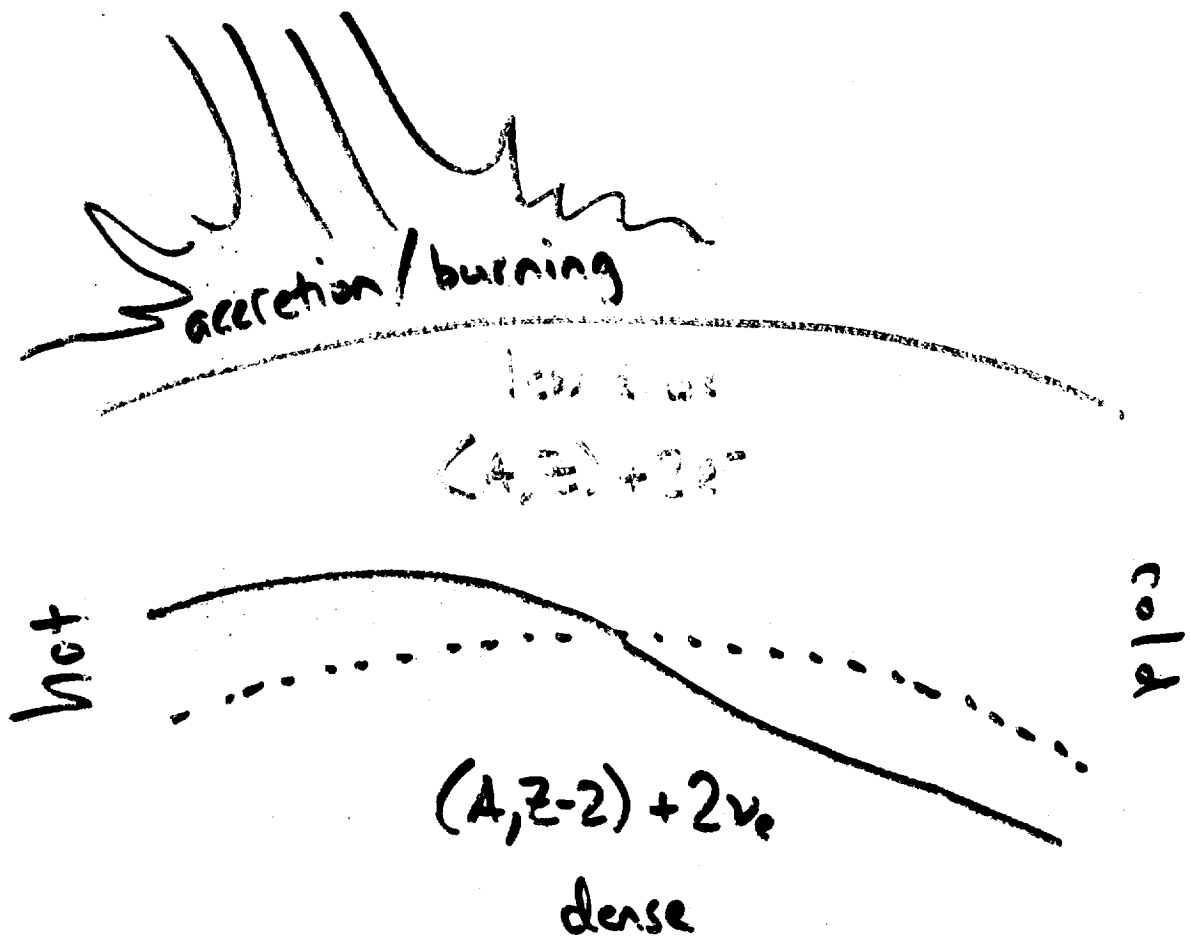
LMXBs are sure to have large electron-capture quadrupoles (LIGO-II).

The r-modes are dead, at least in LMXBs.

Long live the r-modes in newborn neutron stars, though not as long as we thought (still good for LIGO-II).

Electron capture in crust

Bildsten (1998): Strong temperature dependence of e-capture lets accreting neutron stars have quadrupoles detectable by LIGO-II.

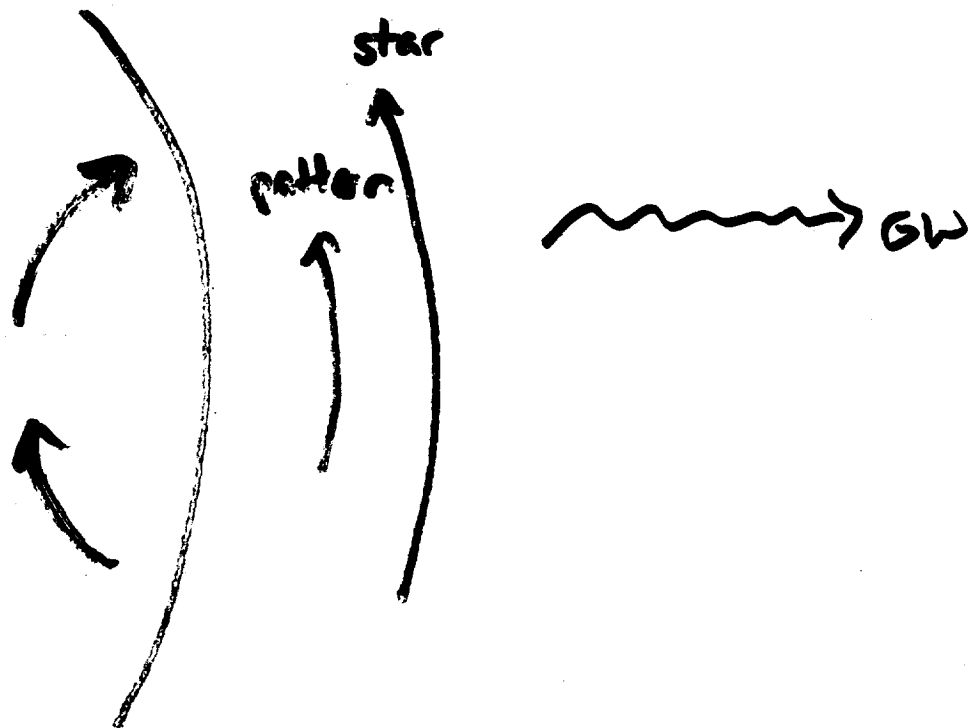


... but neutron-star μ/p is basically Jell-O, mountains tend to sink.

R-mode instability

Owen et al. (1998): Fluid r-modes are CFS unstable in newborn neutron stars, detectable by LIGO-II at several per year.

Bildsten (1998), Andersson, Kokkotas, & Stergioulas (1999): LMXBs too?



Driving beats damping for about a year then superfluid transition.

G. Ushomirsky, C. Cutler & L. Bildsten, astro-ph/0001136

Quadrupole down by 20-50 due to sinkage, but comes back up thanks to multiple layers deep in crust.

Tooth fairy only needs few % lateral temperature gradients, fraction of % (Z/A) gradients (reasonable from x-rays).

Needs breaking strain 10^{-2} , higher than known materials.

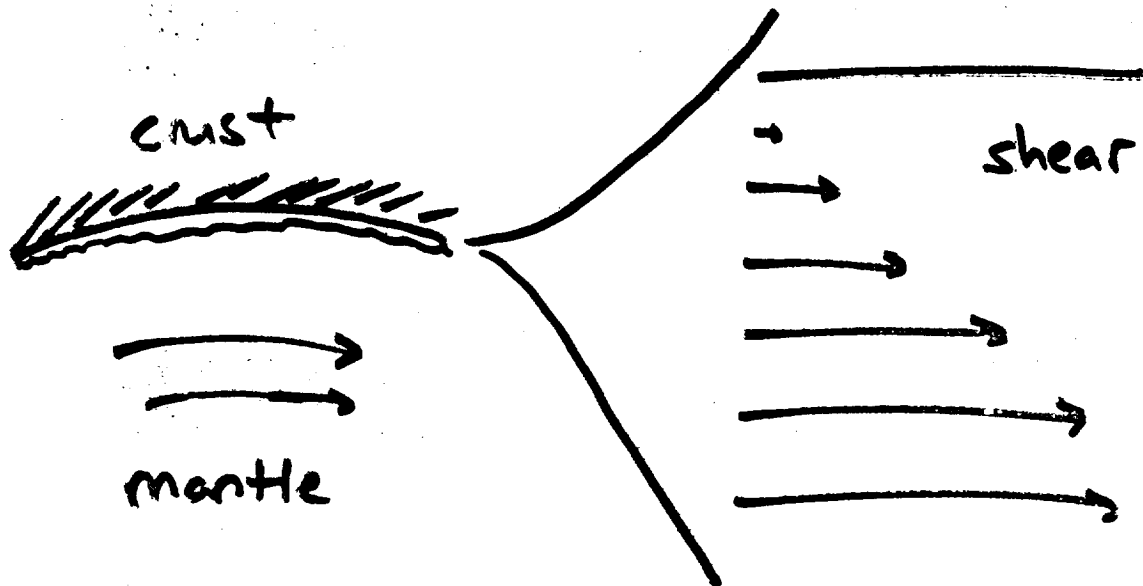
Turned around, maybe Eddington-rate accretion pushes crusts to the breaking point, explaining the 300Hz hangup.

Some (?) quadrupole left over long after accretion stops, hope for old neutron stars.

Boundary layer

L. Bildsten & G. Ushomirsky, ApJ 529, L33

Shear in boundary layer under crust damps 10^5 – 10^7 times faster than we thought, r-modes are no longer unstable with crust, kills LMXBs.



N. Andersson et al., astro-ph/0002114. dispute this, M. Rieutord, astro-ph/0003171 confirms it.

B.J. Owen, unpublished

Crust forms around 10^{10} K, when r-modes are already growing in newborn neutron stars.

Boundary layer → extreme localized heating, neutrino emission and thermal conduction cannot compensate → crust melts (stays melted) if r-mode is big enough.

'Big enough' means $\delta v/v \sim 10^{-3}$, reasonable.

But very sensitive to crust melting temperature (among other things).

Magnetic fields

W.C.G. Ho & D. Lai, astro-ph/9912296: Magnetic braking of spin substantially reduces SNR for $B > 10^{14}$ G (reliable).

L. Rezzolla, F.K. Lamb, & S.L. Shapiro, ApJ 531, L139 and *Y. Levin & G. Ushomirsky, astro-ph/9911295*: r-modes \rightarrow differential rotation \rightarrow field windup (in toy models), substantial loss of SNR for $B > 10^{10}$ G.

But it's a nonlinear hydro effect they're trying to get with linear hydro. *P.M. Sa, unpublished* has found severe problems with the toy models.

Parameter-space grid for Hough hierarchical pulsar search

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What we've got

Parameter space metric: power loss = $g_{\mu\nu}d\lambda^\mu d\lambda^\nu$
for parameters λ^μ guessed wrong by $d\lambda^\mu$.

Define a canonical time so signal $\sim e^{i2\pi f_0 t_c}$,

$$t_c = t_b + \sum_{k=1}^s \frac{f_k}{k+1} t_b^{k+1}. \quad (1)$$

Ptolemaic astronomy is good enough for the metric (CPU-friendlier than JPL ephemeris):

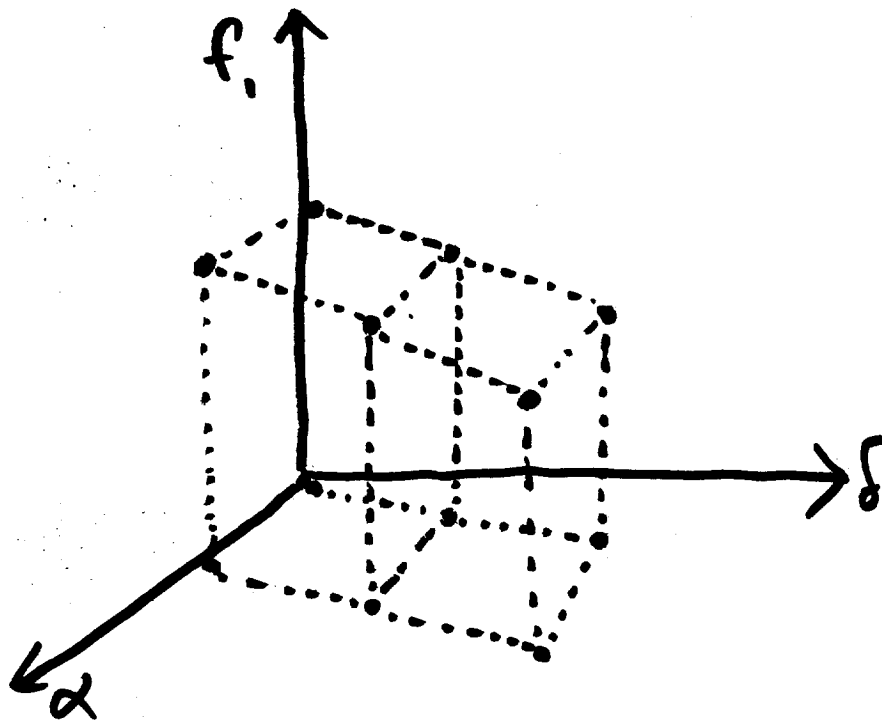
$$t_b = t + r \cos \delta \cos(\omega t + \psi_0) + R \cos \beta \cos(\Omega t + \phi_0) \quad (2)$$

Metric is obtained from derivatives

$$g_{\mu\nu} = 4\pi^2 [\langle (f_0 t_c)_{,\mu} (f_0 t_c)_{,\nu} \rangle - \langle (f_0 t_c)_{,\mu} \rangle \langle (f_0 t_c)_{,\nu} \rangle]. \quad (3)$$

Right now just ~ 1 day coherent integration (coarse stage of hierarchical search), include only daily modulation and spindown (fake yearly as spindown). Harder than expected.

Input: maximum f_0 , minimum spindown age, coherent integration time, sky area to search.



Output: grid points (neglect spindown-position correlations for now, converting to LAL specs)

What's next

LAL specs instead of quick and dirty

Spindown-position correlations change, can't be neglected on longer timescales (fiddle with spindown definitions?)

How to (efficiently) do Hough histogram when the bins are changing? How much are they changing?

Note 1, Linda Turner, 04/27/00 02:37:44 PM
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